Achieving High Efficiency in Solution-Processed Perovskite Solar Cells using C₆₀/C₇₀ Mixed Fullerenes

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ABSTRACT

Fullerenes have attracted considerable interest as an electron transporting layer in perovskite solar cells. Fullerene-based perovskite solar cells produce no hysteresis and do not require high-temperature annealing. However, high power conversion efficiency can only be achieved when the fullerene layer is thermally evaporated, which is an expensive process. In this work, the limitations of a solution-processed fullerene layer have been identified as high crystallinity and the presence of remnant solvents, in contrast to a thermally deposited C₆₀ film, which has low crystallinity and no remaining solvents. As a solution to these problems, a mixed C₆₀ and C₇₀ solution-processed film, which exhibits low crystallinity, is proposed as an electron transporting layer. The mixed-fullerene-based devices produce power conversion efficiencies as high as that of the thermally evaporated C₆₀-based device (16.7%), owing to improved fill factor and open-circuit voltage. In addition, by vacuum-drying the mixed fullerene film, the power conversion efficiency of the solution-processed perovskite solar cells is further improved to 18.0%. This improvement originates from the enhanced transmittance and charge transport by removing the solvent effect. This simple and low-cost method can be easily used in any type of solar cells with fullerene as the electron transporting layer.

1. INTRODUCTION

Organometal halide perovskite solar cells (PSCs) have emerged as next-generation thinfilm solar energy harvesters with high efficiency, owing to the long exciton diffusion length, high absorption coefficient, and high carrier mobility of the perovskite.^{1–5} Fullerenes, which have been used in organic solar cells (OSCs) for more than two decades, have attracted considerable interest as an electron transporting layer (ETL) in PSCs.⁶ Unlike in OSCs, fullerenes are not used as an electron acceptor in the photoactive layer of PSCs; instead, fullerenes function as a chargeselective layer in PSCs and high concentration of fullerene solution is not necessary for PSCs. Furthermore, pristine fullerenes, such as [60]fullerene (C₆₀) and [70]fullerene (C₇₀), which possess intrinsically higher mobility and conductivity than fullerene derivatives, can be directly applied to PSCs without difficult and expensive chemical modification.⁷⁻¹⁰ Currently, metal oxides ETLs, such as TiO₂ and SnO₂, are widely used in PSCs. However, the use of metal oxide ETLs results in large hysteresis of devices,¹¹ originating from charge accumulation¹² and charge trapping^{13,14}. Moreover, their high sintering temperature limits the flexible and large-size application of PSCs^{15,16}. Fullerenes can be used as an ETL to avoid these problems.^{17,18} Fullerenes have low chemical capacitance and high electron affinity, which helps avoid hysteresis.^{19,20} In addition, fullerenes are reported to function as a passivation layer for perovskite grain boundaries, reducing the surface trap density.^{21,22} Nevertheless, high power conversion efficiencies (PCEs) were reported only when fullerenes were thermally deposited or over-coated on metal oxides; the thermal deposition is a slow and expensive process, and the over-coating process requires a metal oxide layer, which renders the use of fullerenes futile.^{23–25} Spin coating of fullerene solution, which is a simple and low-cost process that enables large-size

application, is an ideal alternative. However, solution-processed (spin-coated) fullerene-based PSCs cannot achieve PCEs as high as those of thermally deposited fullerene-based PSCs.

In this work, we studied the reasons for the lower performance of spin-coated fullerenebased PSCs compared with that of thermally deposited fullerene PSCs. Accordingly, we provide a method for achieving higher efficiency using spin-coated fullerenes. With regard to pristine fullerenes, we compared thermally deposited and spin-coated pristine fullerenes in PSCs, and observed that C₆₀ is a better ETL than C₇₀, owing to its lower absorption and better morphology. In addition, thermally deposited fullerenes are a better ETL than spin-coated fullerenes, owing to the amorphous state of the thermally deposited C_{60} films in contrast to the crystalline state of the spin-coated C₆₀ films. Considering that the crystallinity of the fullerene films is linked to the performance of ETLs, we showed that mixing C_{70} into C_{60} reduces the crystallinity of the film. Consequently, the spin-coated mixed-fullerene-based PSC devices produced a PCE of 16.7%, which is similar to that of thermally deposited C_{60} -based PSCs. To enhance the performance of the spin-coated mixed-fullerene-based PSCs further, the fullerene films were subjected to vacuum drying to remove the remaining solvents. The removal of the intercalated solvent from the fullerenes improved the optical transparency and charge transport of the film. The vacuumdried spin-coated mixed-fullerene-based PSCs produced a PCE of 18.0%, which was significantly higher than that of the thermally evaporated fullerene-based PSCs. This study not only achieved a high PCE of the spin-coated fullerene-based PSCs, but also provided an insight into the fullerene engineering of PSCs through various analyses.

2. RESULTS AND DISCUSSION

2.1. Investigation of Pristine Fullerene ETLs in PSCs.

There have been a few reports on the comparison between C₆₀ and C₇₀ as the ETL in PSCs, but these comparisons were limited in terms of device architectures and deposition methods.^{23,26,27} Some of these reports contradicted each other in terms of PCEs and analyses. Therefore, we first comprehensively compare C_{60} and C_{70} in normal and inverted architectures using thermal and solution deposition methods (Figure 1a and S1). Table 1 shows the photovoltaic performance of the thermally deposited C₆₀- and C₇₀-based PSCs in the normal and inverted architectures (Figure S2). In both device types, the C60-based PSCs produced higher PCEs than the C70-based devices, which is consistent with the previous reports.^{23,26,27} The same open-circuit voltage (V_{OC}) values of the C₆₀- and C₇₀-based devices indicate that C₆₀ and C₇₀ possess similar energy disorders^{28–30} and similar energy levels, which are supported by the ab initio calculations (Figure 1b, Figure S3) and photoelectron yield spectroscopy (PYS) measurements (Figure S4).^{31,32} The density functional theory (DFT) calculations of C₆₀ and C₇₀ single molecules reveal that they have similar lowest unoccupied molecular orbital (LUMO) levels whereas their highest occupied molecular orbital levels are within 0.1 eV of each other (Figure S3). This result is corroborated by PYS data, which show that the valence levels of C_{60} and C_{70} are -6.4 eV and -6.3 eV, respectively, for both thermally deposited and spin-coated films (Figure S4a, b, c, and d). The short-circuit current density (Jsc) was lower for the C70-based PSCs than for the C60-based PSCs for the normal architecture, but was similar for both PSCs for the inverted architecture. This indicates that the lower J_{SC} is due to the higher absorption of C_{70} (Figure S5 and S6). The fill factor (FF) of the C₇₀-based PSCs was lower than that of the C₆₀-based PSCs for both the normal- and inverted-type devices. The lower FF of the C70-based PSCs is due to the higher series resistance (Rs) and lower shunt resistance (RsH). Previous studies reported that C70 thinfilms have higher conductivity^{33,34} and electron affinity³⁵ than C₆₀ thin-films. Moreover, C₇₀ is

observed to be more effective in preventing the recombination of excitons than C_{60} , because C_{60} possesses many circulating circuits, resulting in inefficient quantum current distributions inside the molecules.³⁶ However, based on photoluminescence (PL) measurement (**Figure 1c**), it was observed that the C_{60} films quenched the excitons from MAPbI₃ better than the C_{70} films. Notably, the PL peak of MAPbI₃ on C_{60} films is blue-shifted slightly. This blue shift indicates the passivation of the trap states between C_{60} and MAPbI₃ as the trap states lead to spontaneous radiative recombination at the surface.²¹ Furthermore, the full width at half maximum (FWHM) was narrower for the MAPbI₃ on C_{60} compared with that on C_{70} . The narrower FWHM indicates a reduction in the shallow trap density at the interface.³⁷ It has been reported that the low molecular symmetry of C_{70} leads smoother morphology of C_{60} leads to better electron transfer from MAPbI₃.^{38,39} Therefore, we can conclude that the higher FF is due to the better interface between C_{60} and MAPbI₃, thus facilitating the use of C_{60} as the ETL in PSCs.



Figure 1. a) Illustration of normal-type and inverted-type PSC structures used in this work with C₆₀ and C₇₀ molecules. b) Calculated energy band diagrams of C₆₀, C₇₀, and mixed fullerene

analyses in a normal-type PSC. c) Photoluminescence of MAPbI₃ on glass (blue line), MAPbI₃ on C₇₀ on glass (green line), MAPbI₃ on C₆₀ on glass (black line), and MAPbI₃ on C₆₀/C₇₀ on glass (red line).

Table 1. Photovoltaic parameters of the normal-type and inverted-type PSCs using thermally evaporated C₆₀ and C₇₀ as the ETLs under one sun (AM 1.5 G, 100 mW cm⁻²).

Device Type	Fullerene	J _{sc} [mA cm ⁻²]	Voc [V]	FF	$R_{\rm S}$ [Ω cm ²]	$R_{\rm SH}$ [Ω cm ²]	PCE _{Best}	PCE _{Average}	Hysteresis Index ^{11b}
Normal	C ₆₀	22.7	1.01	0.72	37	3.7x10 ⁴	16.7%	16.6% ±0.4	0.01
Normai	C ₇₀	19.3	1.01	0.63	41	2.3x10 ⁴	12.3%	12.1% ±0.2	0.03
Inverted	C ₆₀	22.9	0.90	0.77	9	1.1x10 ⁴	15.8%	15.3% ±0.4	0.01
	C ₇₀	22.9	0.91	0.55	27	7.2x10 ³	11.4%	11.3% ±0.3	0.02

Spin-coated C₆₀ and C₇₀ films were compared in the normal-type PSCs. As the normaltype PSCs exhibit higher efficiency than the inverted-type PSCs, we only consider the normal architecture henceforth. The difference between the spin-coated C₆₀- and C₇₀-based PSCs was similar to that between the thermally deposited C₆₀- and C₇₀-based PSCs (**Table 2, Figure S7a**); the C₇₀-based devices showed a lower J_{SC} and FF than the C₆₀-based devices. Spin-coated fullerene-based PSCs showed lower PCEs than the thermally deposited fullerene-based PSCs (based on a comparison between **Table 1** and **Table 2**). Although the spin-coated fullerene-based PSCs resulted in lower efficiency, the solution process is preferred over the thermal deposition process owing to its lower fabrication cost, higher speed, and simplicity. Thus, it is important to determine why the spin-coated fullerene-based devices resulted in lower performance than the thermal-processed PSCs and to provide a solution to this limitation.

As the lower V_{OC} and FF resulted in the lower PCE of the spin-coated fullerene-based PSCs (thermally deposited C₆₀-based PSCs: $J_{SC} = 22.7$ mA cm⁻², $V_{OC} = 1.01$ V, FF = 0.72 from Table 1; spin-coated C₆₀-based PSCs: $J_{SC} = 23.0 \text{ mA cm}^2$, $V_{OC} = 0.99 \text{ V}$, FF = 0.65 from Table 2), we attribute the lower performance mainly to the crystallinity of fullerene films. In general, thermally deposited fullerene films are reported to be nearly amorphous,⁴⁰⁻⁴² whereas spincoated fullerene films are reported to possess high crystallinity.⁴²⁻⁴⁵ We used grazing-incidence X-ray diffraction spectroscopy (GIXRD) to analyze the thermally evaporated C₆₀ films and the spin-coated C₆₀ films (Figure S8). Figure 2a shows that the spectra of the thermally deposited C_{60} films do not display any peaks, but those of the spin-coated C_{60} films display the (111) peak at 10° and the (220) peak at 18°.46 This was also confirmed by the transmission electron microscope (TEM) images. The thermally deposited C60 films exhibited an amorphous state (Figure 2b, Figure S9a) with no crystal domains (Figure 2c), whereas the spin-coated C60 films exhibited high crystallinity (Figure 2d, Figure S9b) with crystal domains of sizes 20 nm to 100 nm (Figure 2e, Figure S9c). It is known that the boundaries between crystal domains hinder the charge flow whereas charge mobility may be high within a crystal domain. Therefore, unless a film is in a predominately single crystalline phase, amorphous films are expected to have better charge transport. This explains the higher series resistance (Rs) and lower FF of the spin-coated C₆₀-based PSCs as compared with those of the thermally deposited C₆₀-based PSCs. The lower charge transport indicates more charge recombination; thus, lower Voc was obtained for the spincoated fullerene-based PSCs.



Figure 2. a) GIXRD 2 θ scan (incident angle $\omega = 1$ degree) of thermally deposited C₆₀ film (blue line), spin-coated C₆₀ film (black line), and spin-coated mixed fullerene film (red line). TEM images of b), c) thermally deposited C₆₀ film; d), e) spin-coated C₆₀ film with an SAED as an inset; and f), g) spin-coated mixed fullerene film with an SAED as an inset. Optimized structures using Grimme dispersion parameters of h) C₆₀ in fcc crystal structure, and i) C₇₀ in hcp crystal structure.

Table 2. Photovoltaic parameters of the normal-type PSCs using spin-coated C60, (C70, and mixed
fullerene with different ratio as the ETLs under one sun (AM 1.5 G, 100 mW cm ⁻²)).

Fullerene	J _{sc} [mA cm ⁻²]	V _{oc} [V]	FF	$R_{ m S}$ [Ω cm ²]	$R_{ m SH}$ [Ω cm ²]	PCE _{Best}	PCE _{Average}	Hysteresis Index
C ₆₀	23.0	0.99	0.65	42	1.7x10 ⁴	14.8%	14.6% ±0.7	0.03
C ₇₀	19.5	0.97	0.62	36	6.8x10 ⁴	11.7%	11.4% ±0.6	0.07
C ₆₀ /C ₇₀ (9:1)	22.9	1.02	0.71	30	1.5x10⁵	16.7%	16.6% ±0.5	0.04
C ₆₀ /C ₇₀ (5:5)	21.2	1.02	0.55	65	5.7x10 ³	11.8%	11.4% ±0.5	0.05
C ₆₀ /C ₇₀ (1:9)	20.4	1.00	0.55	60	6.8x10 ³	11.1%	11.0% ±0.5	0.05

2.2. Improvement of Spin-coated Fullerene ETLs in PSCs

Considering that the crystallinity of the C_{60} films leads to lower device performance, we added a small amount of C_{70} into the C_{60} solution to suppress the crystallinity of the spin-coated C_{60} film. C_{60} - and C_{70} - mixed fullerene films (referred to as "mixed fullerene films" hereafter) are reported to possess lower crystallinity than pure C_{60} .^{47,48} Furthermore, the native C_{60} and C_{70} mixtures obtained during fullerene synthesis are considerably less expensive than the pure compounds.⁴⁹ GIXRD spectra in **Figure 2a** show that the mixed fullerene film exhibits much weaker (111) and (220) peaks. TEM analysis also reveals that the mixed fullerene film (Figure 2f) displays suppressed crystallinity in contrast to the spin-coated C_{60} film (Figure 2d). This is evident from the selected-area electron diffraction ring patterns (SAED); the SAED pattern of the spin-coated C₆₀ film (Figure 2d inset, Figure S9e) shows stronger crystallinity than that of the mixed fullerene film (Figure 2f inset, Figure S9f). Moreover, crystal domains were not observed in the mixed fullerene films (Figure 2g, Figure S9d). To the best of our knowledge, the low crystallinity of mixed fullerene has not been discussed in detail in the literature so far. We can explain this phenomenon in terms of their packing. Theoretically, C₆₀ packing is more stable in a face-centered cubic (fcc) form (Figure 2h) than in a hexagonal close-packed (hcp) form, whereas C₇₀ is more stable in the hcp form (Figure 2i).⁵⁰ However, in reality, C₆₀ is present in both hcp and fcc forms, whereas C70 only exists in the hcp form.⁵¹ Therefore, in mixed fullerenes, as C₆₀ molecules attempt to form an fcc structure, C₇₀ molecules existing between C₆₀ molecules prevent the formation of fcc packing. For C_{60} in the hcp form, which is less stable, the larger and asymmetric C₇₀ molecules distort the packing, inhibiting the crystallization. The morphologies of the films were investigated using the topography and adhesion modes of atomic force microscopy (AFM). The mixed fullerene film showed lower surface roughness with a roughness average (Ra) of 2.43 nm compared with the spin-coated C₆₀ films, which showed the Ra of 7.19 nm (Figure 3a and b). We could also detect the crystal domain boundaries of the fullerene films by probing the adhesion of the film surface. From the high adhesion force of the crystal boundaries, we could conclude that the C_{60} films possessed many domains of *ca*. 50 nm (Figure 3c). In the case of mixed fullerene, the boundaries were much less defined (Figure 3d). This is evident from the 3D images in which the mixed fullerene films have much flatter topography.



Figure 3. AFM topographic images of spin-coated a) C_{60} film and b) mixed fullerene film with roughness average values. AFM adhesion images of spin-coated c) C_{60} film and d) mixed fullerene film with the 3D images. e) Transmittance of fullerene films with different C_{60} to C_{70} ratios along with their pictures showing different colors.

PSCs with normal architecture were fabricated using the spin-coated mixed fullerene ETLs with different C₆₀ to C₇₀ ratios (**Table 2, Figure S7b**). The best PCE of 16.7% was achieved with the native C₆₀ and C₇₀ mixture ratio (C₆₀:C₇₀ = 9:1 in weight). This PCE is as high as that of the thermally deposited C₆₀-based PSCs (**Table 1**). The devices with greater C₇₀ content exhibited lower J_{SC} and lower FF (**Table 2**). UV–vis spectra with different C₆₀/C₇₀ ratios showed that the increase in the amount of C₇₀ reduces the transmittance of the film (**Figure 3e**), thus reducing the J_{SC} . The C₆₀/C₇₀ films with the weight ratio 9:1 showed almost the same absorption as that of the pure C₆₀, as indicated by the similar J_{SC} values. The lowering trend of FF is observed because the charge transfer between C₇₀ and the perovskite layer is not favored as discussed in section 2.1.1. We used solid-state ab initio calculations to assess the energy level change when C₆₀ and C₇₀ are mixed and the contributions from each component to the density of

states (DOS) in the conduction band. Solid-state calculations were performed at the dispersioncorrected density functional tight binding level (Figure S3). The DOS of C₆₀ and C₇₀ solids (Figure S10a) show the same energy levels to C₆₀ and C₇₀, similar to the comparison at the single-molecule level in **Figure S3**. Mixing C_{60} and C_{70} yielded a similar result in that the energy contributions from C₇₀ and C₆₀ to the conduction band are similar (Figure 1b, Figure S10b), indicating that there is no LUMO offset within the mixed film, which could potentially inhibit charge transport. This is corroborated experimentally by the PYS measurement (Figure S4e). We also conducted simulations of selected low-index surfaces and confirmed that there is no change in DOS owing to the presence of the mixed fullerene interface (Figure S11). In addition, the ab initio calculations suggest that the mixed system is only approximately 0.03 eV/molecule above the convex hull (Figure S12), which is approximately kT at room temperature; there is, therefore, no significant driving force for segregation or desegregation. These results indicate that there is no significant change in the energy level when C₇₀ is introduced, and the reduction in Voc can be ascribed to the increase in charge recombination. Notably, all devices that used the mixed fullerenes showed a higher Voc than pure C₆₀ or C₇₀ was applied. According to our computational results, it vividly demonstrated that the conjugation of the mixed C₆₀ and C₇₀ is much intact, providing a better interface for forming the perovskite at the interface with few voids, which thusly slightly increased the V_{OC} of devices. Consequently, we can conclude that the native mixed fullerene weight ratio of 9:1 is the optimized ratio, and mixing too much C₇₀ undermines the device performance (Figure 4a).



Figure 4. Photovoltaic parameters and statistical analysis of thermally deposited C_{60} -based PSCs (black circles), spin-coated C_{60} -based PSCs (red triangles), spin-coated mixed-fullerene-based PSCs (blue diamonds), and vacuum-dried mixed-fullerene-based PSCs (purple squares), showing a) *J*sc, b) *V*oc, c) FF, and d) PCE. e) XPS chlorine spectra of a mixed fullerene film thermally annealed at 100 °C (black), a mixed fullerene film thermally annealed at 200 °C (green), and a mixed fullerene film thermally annealed at 100 °C (black), a mixed fullerene film thermally annealed at 100 °C (black) and a vacuum-dried mixed fullerene film thermally annealed at 100 °C (black) and a vacuum-dried mixed fullerene film thermally annealed at 100 °C (black) and a vacuum-dried mixed fullerene film (purple). Cross-sectional SEM pictures of spin-coated g) mixed fullerene film and h) the same film after vacuum-drying it on ITO substrates with magnified images as insets.

2.3. Further Enhancement of Mixed Fullerene ETL.

Considering that the mixed fullerene ETLs can achieve PCE as high as that of the thermally deposited C₆₀ ETLs, we report a method to improve the PCE of these mixed fullerene ETLs even further. In addition to the amorphous state, a favorable aspect of the thermally evaporated C₆₀ is that it is solvent-free. By removing the trapped solvents from the spin-coated mixed fullerene films, we postulate that the fullerenes can be packed in a more compact manner, improving the charge transport and eliminating the solvent effects.^{52,53} To remove the orthodichlorobenzene (ODCB) effectively, we applied three different treatments to the spin-coated mixed fullerene films: thermal annealing (TA) at 100 °C, TA at 200 °C, and TA at 100 °C under a light vacuum (0.01 MPa). Table 3 shows the photovoltaic performances of the PSCs with different post-treatments applied to the spin-coated mixed fullerene ETLs. The PSCs treated with TA at 100 °C did not exhibit a significant change in the device performance (Figure S13). We suspect that this temperature is not sufficiently high to remove ODCB (boiling point of 180 °C). In contrast, the PSCs treated with TA at 200 °C exhibited reduced device performance (Table 3, Figure S13). According to the AFM morphology analysis in Figure S14a, TA at 200 °C resulted in aggregation of the fullerene films. While such a high temperature is effective for removing ODCB, it leads to strong crystallization, which severely roughens the morphology as indicated by the low FF.54,55 Finally, we vacuum-dried the film at 100 °C for 10 min. The device performance increased to 18.0% with the increase in J_{SC} and V_{OC} whereas FF remained the same (Figure 4a, b, c, and d; Table 3; Figure S13). The AFM roughness image and the corresponding Ra value reveal that the morphology of the fullerene film improved with the vacuum-drying treatment (Figure S14b). However, the improvement in morphology does not explain the improvements in J_{SC} and V_{OC}. X-ray photoelectron spectroscopy (XPS) was used to

confirm the removal of solvents via vacuum-drying. The data indicate that the chlorine from ODCB solvent, which was detectable in films treated with TA at 100 °C, was not detectable in the films treated with TA at 200 °C and vacuum-dried films (Figure 4e). UV-vis absorption spectra in Figure 4f indicate that the vacuum-drying treatment reduced the absorption of the film significantly. We conjecture that the removal of ODCB increased the transmittance of the film by removing solvatochromism.^{54,55} Cross-sectional scanning electron microscope (SEM) image of the spin-coated mixed fullerene film, with a magnified image as an inset, shows that the fullerene film has a thickness of approximately 32 nm (Figure 4g), which is similar to the optimized thickness i.e., 35 nm of the thermally deposited C₆₀ film (Figure S15). After the vacuum-drying treatment, we can observe that the thickness of the same mixed fullerene film decreased to *ca.* 12 nm (Figure 4h). GIXRD spectra showed a decrease in the crystallinity of the spin-coated mixed fullerene films upon vacuum-drying treatment, thus resembling the thermally evaporated C₆₀ film (Figure S8). This confirms that the removal of solvent affected the packing and crystallinity of the fullerenes. The combination of enhanced transmittance and charge transport led to the increase in J_{SC} and V_{OC} . Ultimately, the PCE of the vacuum-dried solutionprocessed mixed-fullerene-based PSC (18.0%) exceeded that of the thermally evaporated C₆₀based PSC (16.7%). Moreover, this vacuum-dry method can be applied to other solutionprocessed fullerene ETLs such as C₆₀ and C₇₀, which gave a better performance than none treated devices (Table S2, Figure S16). For instance, the vacuum-dried solution-processed C60 achieved a PCE in 16.0%, exceeding non-treated solution-processed C_{60} with a PCE in 14.8%. The vacuum-dried solution-processed C70 (12.1%) slightly exceed the non-treated solution-processed C70 (11.7%).

Table 3. Photovoltaic parameters of the normal-type PSCs using mixed fullerene as the ETL with different film treatments under one sun (AM 1.5 G, 100 mW cm^{-2}).

Fullerene	Treatment	J _{SC} [mA cm ⁻²]	V _{oc} [V]	FF	$R_{\rm S}$ [Ω cm ²]	$R_{ m SH}$ [Ω cm ²]	PCE _{Best}	PCE _{Average}	Hysteresis Index
Spin-coated C ₆₀ /C ₇₀ (9:1)	none	22.9	1.02	0.71	42	1.7x10 ⁴	16.7%	16.6% ±0.5	0.04
	100 °C TA	23.2	1.04	0.71	33	1.7x10 ³	16.9%	16.0% ±0.4	0.04
	200 °C TA	22.5	1.05	0.66	20	4.0x10 ³	15.7%	14.8% ±1.0	0.13
	Vacuum-dry at 100 °C	24.3	1.03	0.73	28	1.5x10 ⁴	18.0%	17.7% ±0.6	0.01

3. CONCLUSIONS

We compared C_{60} and C_{70} in PSCs. The comparison was extended to thermally evaporated and solution-processed deposition methods of fullerenes in normal and inverted architectures of PSCs. C_{60} as ETL in PSC exhibited higher photovoltaic performance than C_{70} . The thermally deposited C_{60} -based normal-type PSCs showed a PCE of 16.7%. By investigating the mechanism behind the efficiency enhancement, we discovered that the high performance of the thermally deposited C_{60} -based PSCs is attributed to the crystallinity of the fullerene and the presence of remnant solvents. By mixing a small amount of C_{70} into a solution of C_{60} , the crystallinity of the spin-coated fullerene films decreased substantially, resulting in a PCE close to that of the thermally deposited C_{60} -based PSCs. Furthermore, by vacuum-drying the mixed fullerene film, a PCE of 18% without hysteresis was achieved, which exceeds that of the thermally deposited C_{60} -based PSCs.

4. EXPERIMENTAL

Mixed-fullerene solution preparation. C_{60} and C_{70} were dissolved in a solution of ODCB in different ratios with a mass concentration of 20.0 mg mL⁻¹. The solution was filtered through a 0.45 µm polytetrafluoroethylene filter prior to use.

Spiro-MeOTAD solution preparation. A solution was prepared by mixing 85.8 mg spiro-MeOTAD (Lumtech), 19.3 μ L of a stock solution of 520 mg mL⁻¹ lithium bis(trifluoromethylsulphonyl)imide (Aldrich) in anhydrous acetonitrile, and 33.8 μ L of 4-*tert*-butylpyridine (Aldrich) in 1 mL anhydrous chlorobenzene.

PEDOT:PSS solution preparation. Poly(3,4-ethylenedioxythiophene) polystyrene sulfonate (PEDOT:PSS) solution was produced by adding 0.5 wt% polyoxyethylene(6) tridecyl ether (Sigma-Aldrich Chemical Co., Inc.) to a PEDOT:PSS dispersion in water (Clevios P VP, Heraeus Precious Metals GmbH & Co.). The solution was filtered through a 0.45 μm polytetrafluoroethylene filter prior to use.

Perovskite (MAPbI₃) precursor solution preparation. CH₃NH₃I (TCI), PbI₂ (TCI), and anhydrous dimethyl sulfoxide (TCI) (molar ratio 1:1:1) were mixed in anhydrous N,N-dimethylformamide (TCI) with a concentration of 50 wt%. The solution was filtered through a 0.45 µm polytetrafluoroethylene filter prior to use.

Fabrication of inverted-type planar heterojunction perovskite solar cells. Indium tin oxide (ITO) patterned glass substrates with a size of $15 \times 15 \text{ mm}^2$ and sheet resistance of 6 Ω sq⁻¹ (Techno Print Co., Ltd.) were cleaned and sonicated with a detergent, distilled water, acetone, and isopropanol in an ultrasonic bath for 15 min. Subsequently, the substrates were treated with ultraviolet/ozone (UV/O₃) for 15 min prior to use. Thereafter, 25 µL of PEDOT:PSS solution was spin-coated on UV/O₃-treated ITO substrates at 3000 rpm for 30 s. The coated ITO substrates were subsequently annealed at 105 °C for 10 min. Thereafter, 25 µL of the perovskite precursor solution was spin-coated on the PEDOT:PSS layer at 3000 rpm for 30 s, and 0.5 mL of anhydrous diethyl ether was slowly dripped onto the substrate 10 s after the start of the spin-coating process. Subsequently, the film was annealed at 100 °C for 10 min to obtain a dense brown MAPbI3 film. Subsequently, C₆₀ or C₇₀ was thermally deposited on the perovskite film to obtain a thickness of 30 nm at a constant evaporation rate of 0.02 nm s⁻¹. Similarly, bathocuproine was thermally deposited to obtain a thickness of 10 nm at the same evaporation rate of 0.05 nm s⁻¹.

Fabrication of normal type planar heterojunction perovskite solar cells. ITO patterned glass substrates were cleaned and sonicated with detergent, distilled water, acetone and isopropanol in an ultrasonic bath for 15 min, respectively. Subsequently, 25 μ L of mix-fullerene solution was spin-coated on ITO substrates at 3000 rpm for 30 s. The coated ITO substrates were used without any treatment or were followed by different processes, such as: 1) annealing at 100°C for 1 h; 2) annealing at 200 °C for 1 h; 3) drying up at 100°C under vacuum condition (~1.0×10⁻³ MPa) for 1 h. If thermally deposited fullerene film was used, C₆₀ or C₇₀ was thermally deposited on ITO substrates in a thickness of 30 nm at a constant evaporation rate of 0.02 nm s⁻¹.

Then, 25 μ L of perovskite precursor solution was spin-coated on the fullerene layer at 3000 rpm for 30 s, with a dropping of 0.5 mL of anhydrous diethyl ether was slowly dripped onto the substrate 10 s after the start of the spin-coating process. Next, the film was annealed at 100 °C for 10 min to obtain a dense brown MAPbI₃ film. The hole transporting layer was spin-coated from the spiro-MeOTAD solution at 4000 rpm for 30 s. Finally, a 70-nm-thick of Au anode was fabricated by thermal deposition at a constant evaporation rate of 0.05 nm s⁻¹.

Characterizations. The J-V characteristics were measured using a software-controlled source meter (Keithley 2400 SourceMeter) under dark conditions and the simulated sunlight irradiation of 1 sun (AM 1.5G; 100 mW cm⁻²) using a solar simulator (EMS-35AAA, Ushio Spax Inc.) with a Ushio Xe short arc lamp 500. The source meter was calibrated using a silicon diode (BS-520BK, Bunkokeiki). SEM analysis of the perovskite films was performed using an S-4800 (Hitachi). The TEM images are taken by JEM-2010F (JEOL Ltd.) with a thermal field emission gun operated at 200 keV. The Selected area electron diffraction (SAED) patterns are recorded by a Charge-coupled Device (CCD) at a camera length of 60 cm.

Shimadzu UV-3150 was used for the UV–vis–NIR measurement. The PL measurements were performed using JASCO Spectrofluorometer (FP-8300). Topography images were recorded by using an atomic force microscope (AFM) operating in tapping mode (SPI3800N, SII). The grazing-incidence XRD 2 θ scans were performed on a Jordan Valley D1 diffractometer with copper K α_1 radiation and a parallel beam source. In the 2 θ scans, the scattering angle 2 θ between incident beam and diffracted beam changes while the incident angle ω between incident beam and the sample surface is fixed at 1 degree. The valence band and Fermi levels measurements were performed using Riken Keiki PYS-A AC-2 and Kelvin probe spectroscopy in air (ESA), respectively. The photoemission measurements were performed using XPS (PHI5000, Versa

Probe) with monochromatic Al K α radiation. The water contact angle measurements were performed using a contact angle meter (DMo-501, Kyowa Interface Science Co., Ltd.).

ASSOCIATED CONTENT

Supporting Information.

Cross sectional SEM, device fabrication, various computational calculations, PYS, UV-vis absorption, GIXRD, TEM, AFM.

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Author Contributions

H.L and I.J conceived the experiment. H.L and S.S conducted device fabrication. R.X carried out TEM analyses. J.L and Y.Y carried out AFM analyses. C.L and M.G carried out GIXRD analyses. A.P and S.M conducted computational analyses. H.L, I.J, S.M, and Y.M did the rest of the analyses and the related discussion. H.L and I.J wrote the manuscript.

ACKNOWLEDGMENT

We gratefully acknowledge the Research and Education Consortium for Innovation of Advanced Integrated Science by Japan Science and Technology (JST) and Japan Society for the Promotion of Science (JSPS) KAKENHI Grant Numbers JP15H05760, JP16H02285, 17K04970, and 17H06609.

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Table of Contents



Supporting Information

Achieving High Efficiency in Solution-Processed Perovskite Solar Cells using C₆₀/C₇₀ Mixed Fullerenes

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Figure S1. Cross-sectional SEM images of a) normal-type PSCs and b) inverted-type PSCs.



Figure S2. *J*–*V* curves of the thermally deposited C₆₀- (black squares and line) and C₇₀- (red circles and line) based PSC for a) the normal-type structure and b) the inverted-type structure.

a) Frontier orbital comparison of single molecules at the DFT level: [6-31+G(d,p)] The calculations were performed in Gaussian 09^[S1] with the 6-31+g(d,p) basis, in vacuum. HOMO-LUMO distributions:



b) Energy level comparison of single molecules at the DFTB level:

The calculations were performed in DFTB+^[S4] using the 3ob-3-1^[S5,S6] parameter set The HOMO-LUMO distributions:



The HOMO-LUMO values are as below:

C ₆₀ : HOMO: -5.67 eV	C ₇₀ : HOMO: -5.60 eV
LUMO: -3.88 eV	LUMO: -3.97 eV

Figure S3. Frontier orbital calculations of C₆₀ and C₇₀ a) at the DFT level, and b) at the DFTB level.



Figure S4. PYS spectra of a) thermally deposited C₆₀, b) thermally deposited C₇₀, c) spin-coated C₆₀, d) spin-coated C₇₀, e) spin-coated C₆₀/C₇₀ (9:1), and f) vacuum-dried spin-coated C₆₀/C₇₀ (9:1).



Coated on Silicon Coated on glass Coated on ITO

Figure S5. Various C₆₀ and C₇₀ films on glass and silicon substrates. The difference in transparency between C₆₀ and C₇₀ of the same thickness is visible even for the naked eye.



Figure S6. a) UV-vis absorption of spin-coated fullerene films with different C₆₀ to C₇₀ ratios. b) UV-vis absorption spectroscopy of thermally deposited C₆₀ and C₇₀.



Figure S7. *J*–*V* curves of a) the normal-type PSCs using spin-coated C₆₀- (black) and C₇₀- (red); the normal-type PSCs using b) mixed fullerenes with different C₆₀ to C₇₀ ratios; c) the inverted-type PSCs using thermally deposited C₆₀- (black) and C₇₀- (red), and mixed fullerenes (blue).

Table S1. Photovoltaic parameters of the inverted-type PSCs using thermally deposited C_{60} , thermally deposited C_{70} , and spin-coated mix-fullerene as the ETLs under one sun (AM 1.5 G, 100 mW cm⁻²).

Fullerene	J _{sc} [mA cm ⁻²]	Voc [V]	FF	$R_{\rm S}$ [Ω cm ²]	$R_{ m SH}$ [Ω cm ²]	PCE _{Best}	PCE _{Average}	Hysteresis Index
C ₆₀	22.9	0.90	0.77	9	1.1x10 ⁴	15.8%	15.3% ±0.4	0.01
C ₆₀ /C ₇₀ =9:1	22.1	0.90	0.76	10	1.0x10 ⁴	15.2%	15.0% ±0.5	0.01
C ₇₀	22.9	0.91	0.55	27	7.2x10 ³	11.4%	11.3% ±0.3	0.02



Figure S8. GIXRD scan of the thermally deposited C₆₀ film (black), the spin-coated C₆₀ film (red), the spin-coated mixed fullerene film (blue), the vacuum-dried mixed fullerene film (purple), and 200 °C TA-treated mixed fullerene film.



Figure S9. Supplementary TEM images of a) a thermally deposited C₆₀ film, b) a boundary of crystal domains in a spin-coated C₆₀ film, c) many crystal domains in a spin-coated C₆₀ film, d) one large uniform domain in a mixed fullerene film, and supplementary SAED of d) a spin-coated C₆₀ film and e) a mixed fullerene film.



Comparison of solids at the DFTB level (dispersion correction with the Grimme scheme):^[S7]

a) Density of states (DOS) of C₆₀ and C₇₀ crystals (left) with the LUMO magnified (right), showing the C₆₀ and C₇₀ LUMO position separately.



b) DOS of C_{60} and C_{70} parts in the mixed C_{60}/C_{70} structure.

Figure S10. Densities of states of a) C₆₀ and C₇₀, and b) mixed fullerene (C₆₀:C₇₀ 29:3) in solid state at the DFTB level.

Effect on surface on DOS at the DFTB level:

111 surface of C_{60} in fcc structure:





 $001 \ surface \ of \ C_{70} \ in \ hcp \ structure:$ Similar to $C_{60}, \ we \ cut \ the \ 001 \ surface \ of \ C_{70} \ hcp \ crystal.$



DOS of C_{70} and 001 surface of C_{70} :



The lowest energy surfaces of both the C₆₀ and C₇₀ crystal behave similar with the bulk.

Figure S11. Computational energy level calculations of C₆₀, C₇₀, and mixed fullerene in which crystallinity and surface have been taken into account in solid at DFTB level.

Convex hull: It seems from the convex hull that there is very small possibility of segregation in mixed system, as the formation energy change is negligible ($\sim 0.03 \text{ eV}$, i.e on the order of kT). Thus, the mixed C₆₀/C₇₀ system likely does not segregate into C₆₀/C₇₀ layers.



The electron transfer rate of C_{60}/C_{70} mixed system.

Figure S12. Segregation mechanism calculation of C₆₀ and C₇₀ mixture and electron transfer rate calculation.



Figure S13. *J–V* curves of PSCs after 100 °C TA (green line), 200 °C TA (blue line), and 100 °C TA under light vacuum (purple line) on spin-coated mixed fullerene ETLs.



Figure S14. AFM images of a) 200 °C TA-treated spin-coated mixed fullerene film, and b) 100 °C TA-treated under light vacuum of spin-coated mixed fullerene film.



Figure S15. Cross sectional SEM of thermally deposited C60 on ITO substrate.



Figure S16. J–V curves of the normal-type PSCs using spin-coated C₆₀ (black), spin-coated C₇₀ (red), vacuum-dry C₆₀ (blue), and vacuum-dry C₇₀ (green).

Table S2. Photovoltaic parameters of the normal-type PSCs using spin-coated C₆₀, C₇₀ as the ETLs with vacuum-dry treatment under one sun (AM 1.5 G, 100 mW cm⁻²).

Fullerene	J _{SC} [mA cm ⁻²]	V _{oc} [V]	FF	$R_{ m S}$ [Ω cm ²]	$R_{ m SH}$ [Ω cm ²]	PCE _{Best}	PCE _{Average}	Hysteresis Index
C ₆₀	23.0	0.99	0.65	42	1.7x10⁴	14.8%	14.6% ±0.7	0.03
Vacuum-dry C ₆₀	24.1	0.99	0.67	37	1.0x10 ⁴	16.0%	15.7% ±0.5	0.03
C ₇₀	19.5	0.97	0.62	36	6.8x10 ⁴	11.7%	11.4% ±0.6	0.07

Vacuum-dry C ₇₀	21.7	0.98	0.60	38	1.2x10⁴	12.1%	11.9% ±0.4	0.05
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