

Tailoring carbon nanostructures: Unraveling the electronic properties of low- dimensional quantum solids

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Outline

- Introduction/Motivation/Experimental
- Part 1: Angle resolved dielectric response of SWCNT
 - a) SWCNT bundles vs. individual SWCNT
 - b) Implications on excitation spectrum in 1D,2D
 - c) 1D „Drude“ plasmon in intercalated SWCNT?
- Part 2: Electronic structure of metallicity selected SWCNT
 - a) Textbook example for intrinsic XPS line shapes
 - b) Unravelling the 1D DOS in PES and XAS
 - c) Nature of metallic ground state in metallic SWCNT
- Part 3: Functionalized SWCNT:
 - a) Examples for intercalation, filling and substitution
 - b) Nanochemical reactions inside SWCNT
- Part 4: Electronic and phononic structure of graphene systems
 - a) Electron dispersion, phonon dispersion and EPC
 - b) GIC KC_8 revisited: A key to graphene
- Summary and outlook

Low dimensional quantum systems?

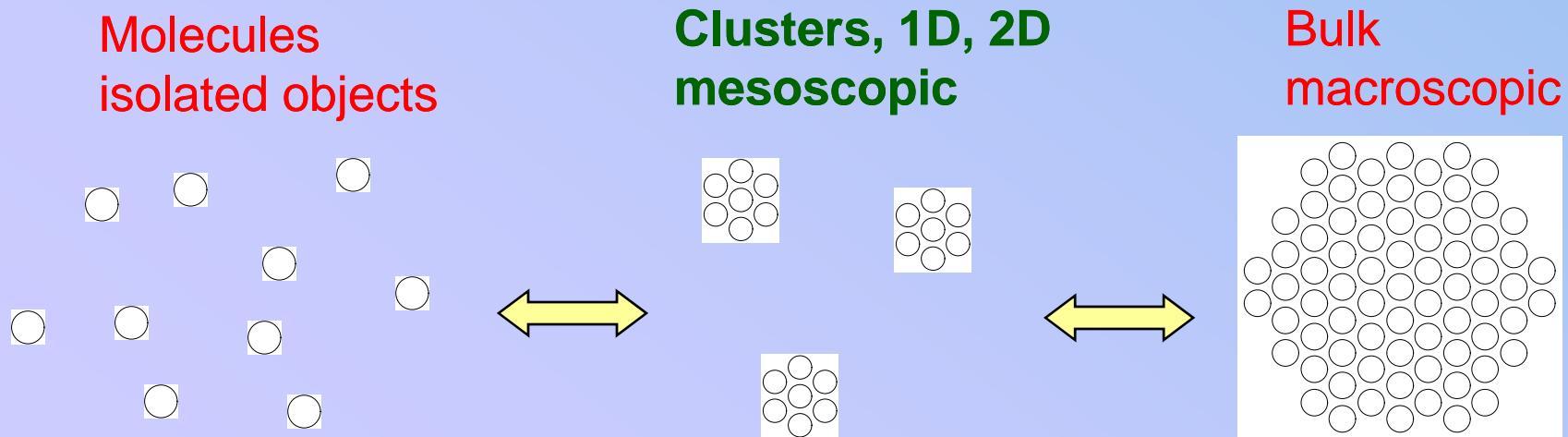
- All solids are of course quantum systems
- 1D and 2D quantum solids:
 - Mesoscopic systems with size quantization
 - Quantum effects already at room temperature
 - Correlation effects crucial:
electron-hole, electron-electron, electron-phonon
 - Instabilities in 1D and 2D systems!



Molecular carbon nanostructures are
archetypical examples of perfect 1D and 2D solids

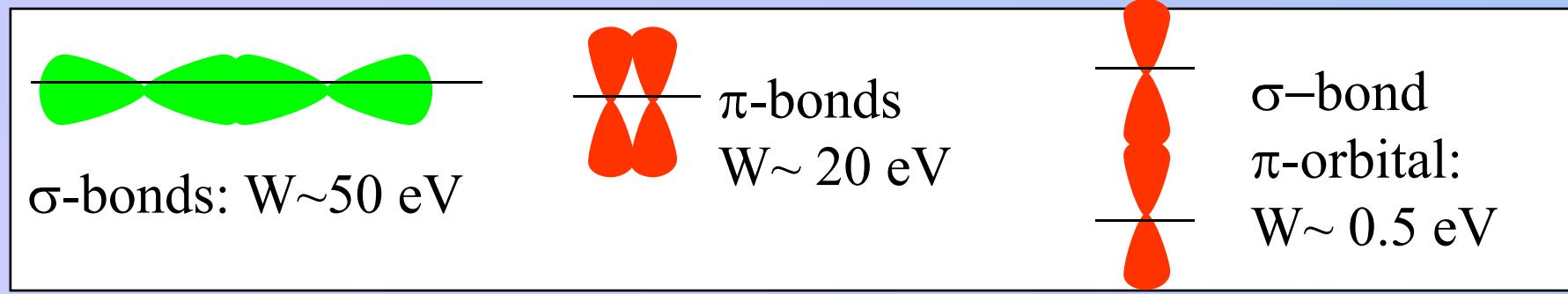
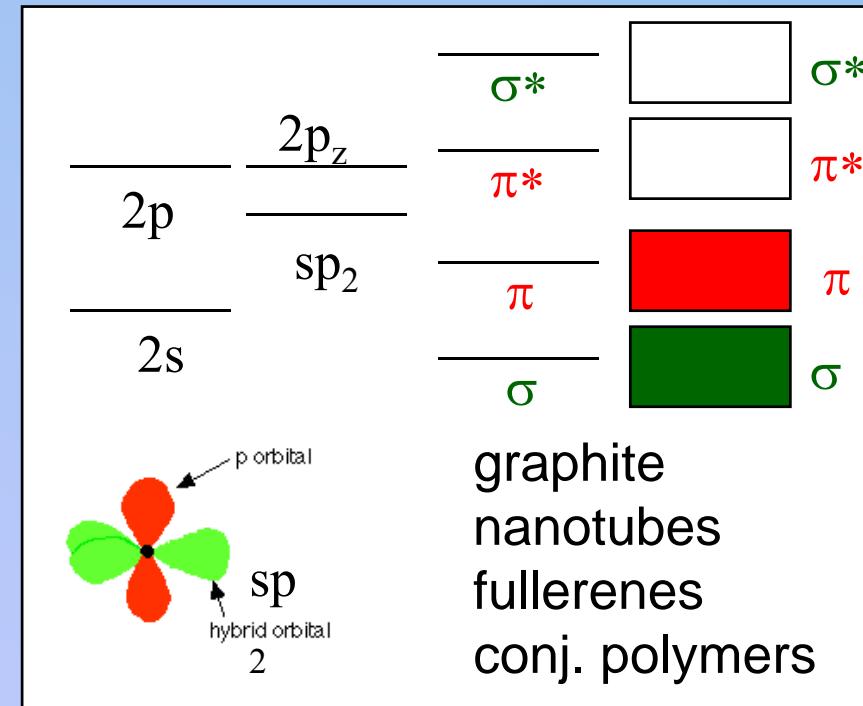
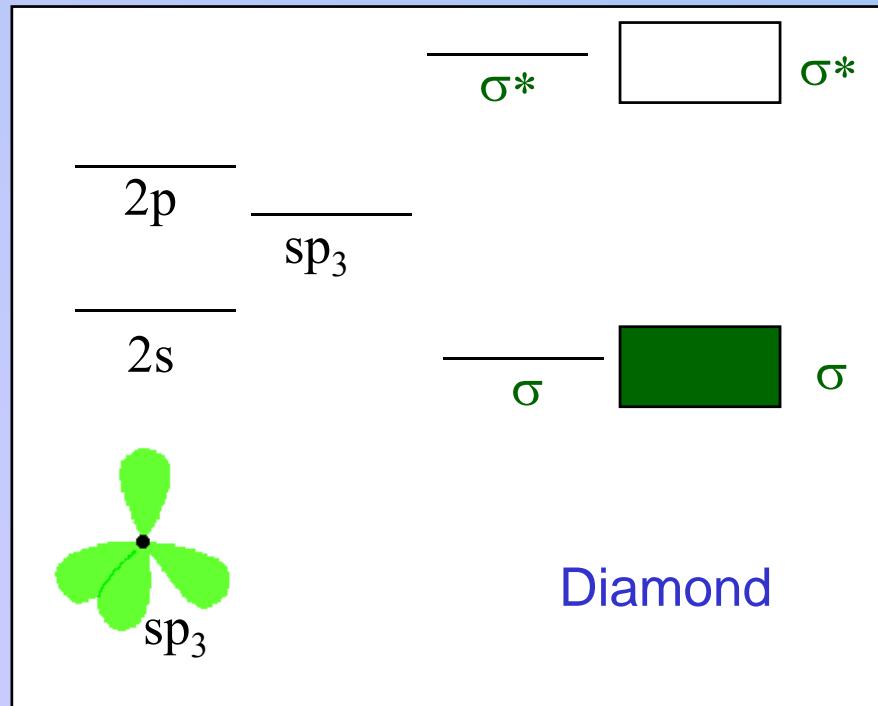
More general Framework

- Electronic/optical properties: Interaction of light and charges with matter
- Novel behavior in micro- and nano-meshed meta-materials.
- How much do surface and/or bulk properties actually contribute?
- What is the needed separation for isolated objects as low D systems?
- Can we tailor the properties?



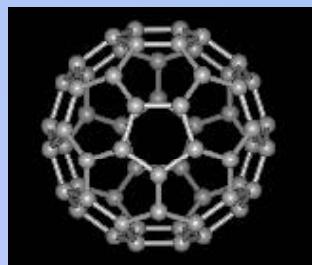
Electronic structure of different carbon allotropes

Two 1s core electrons and 4 valence electrons (two 2s and two 2p)



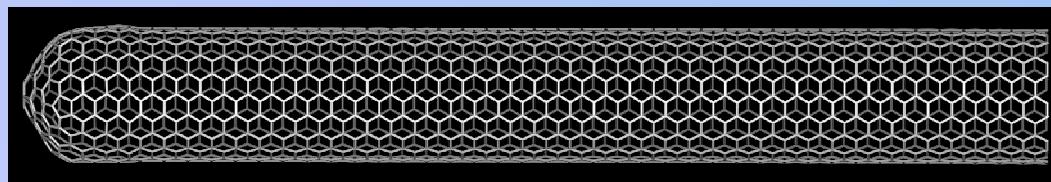
All-carbon Nanostructures

0D



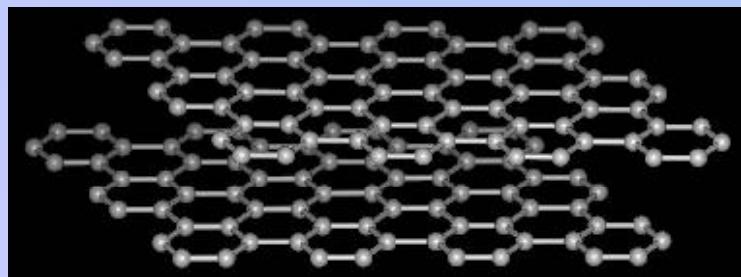
Fullerenes, C_{60}

1D



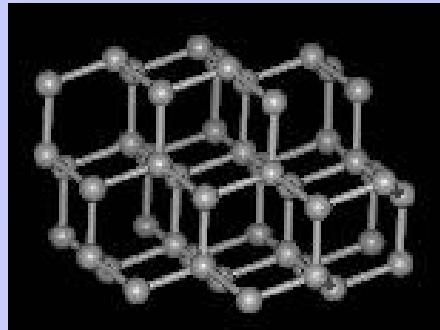
Nanotubes

2D



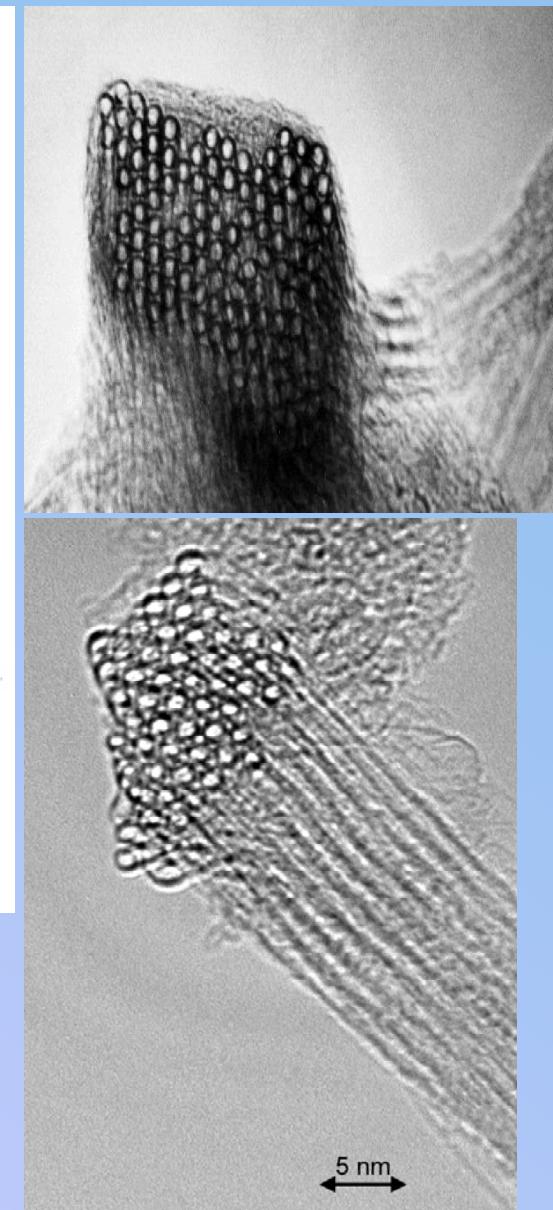
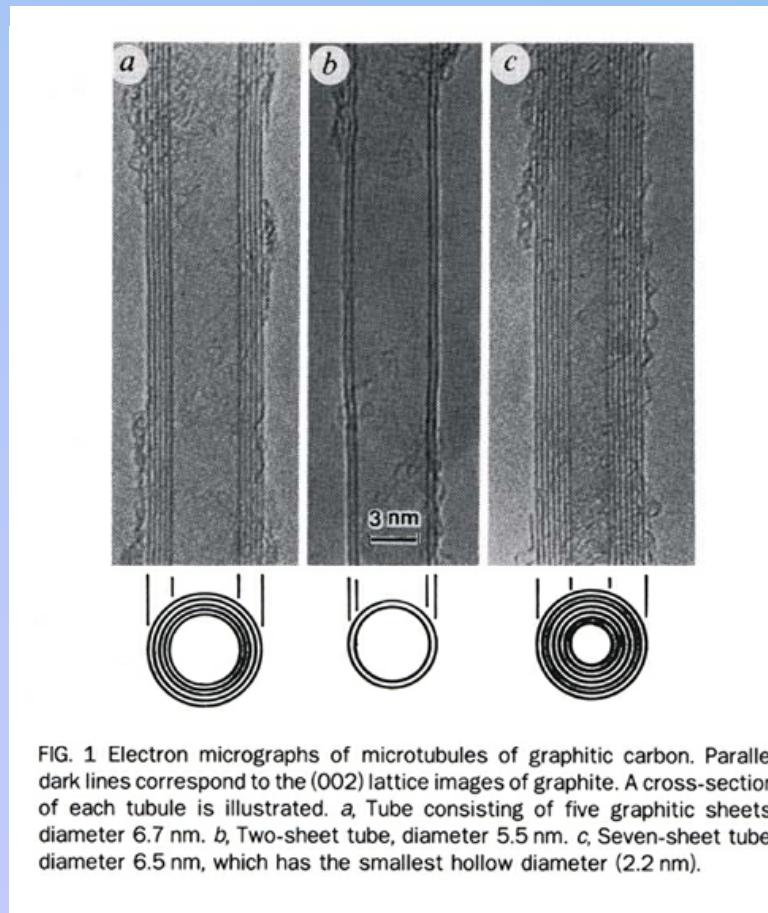
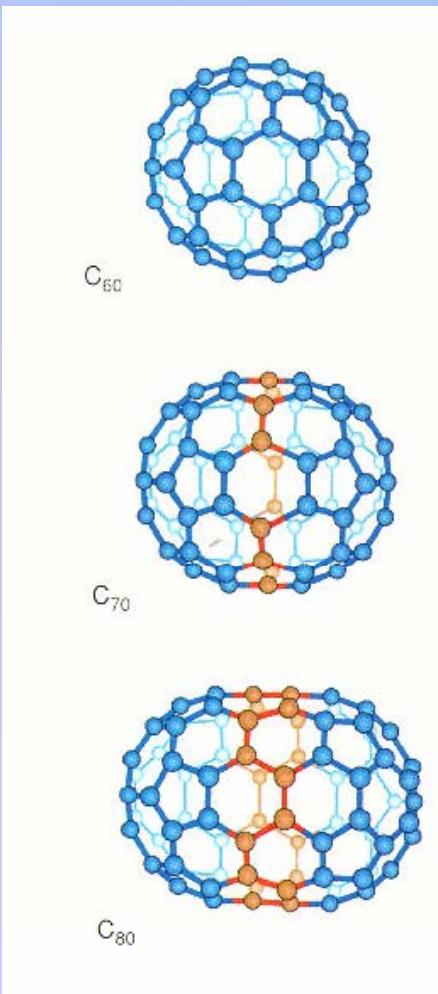
Graphite

3D



Diamond

Carbon nanotubes



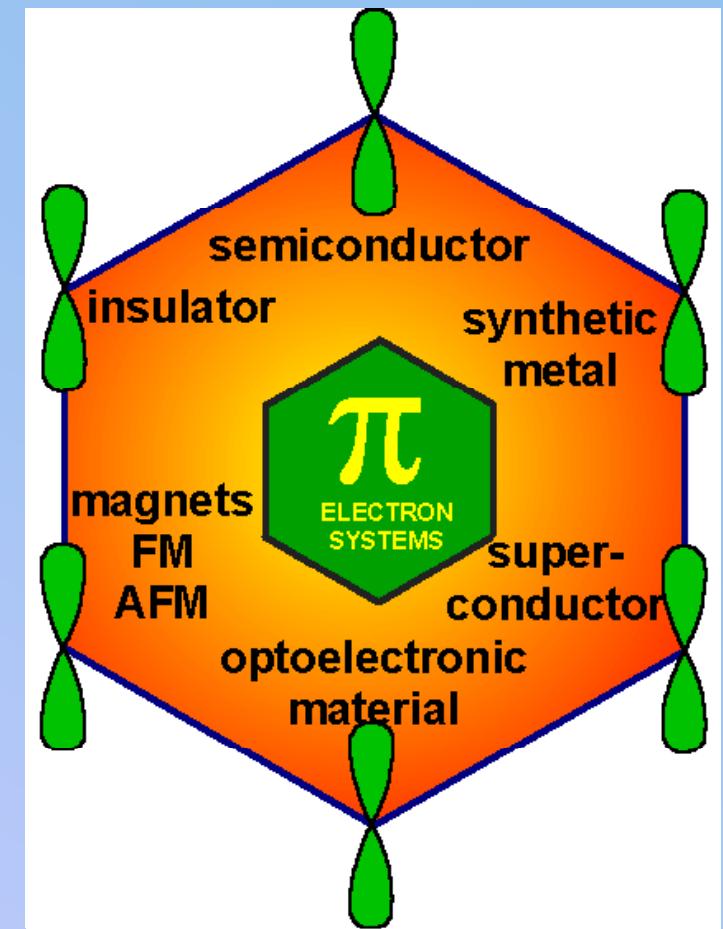
1991 Iijima, TEM of carbon nanotubes
d ~ nm, l ~ 100 nm - 1 mm, MWCNT
SWCNT bundles 5-20 nm diameter

Why carbon nanotubes?

● Archetypical 1D system:

New electronic, optical and structural properties:

- Ideal 1D quantumwires
- Hard, ductile Nanowires
- Molecular Magnets, Spintronics
- Protection for reactive elements and molecules
- Gas storage
- FET, VIAS, optoelectronic
- Sources for field emission
- Organic superconductors
- Li-Ion batteries
- Sensors, actuators, NT-yarns, composites, drug delivery.....



→ Ideal tools for nanoelectronics/optics/mechanics/(bio)chemistry

Special properties of nanotubes

- Electronic properties:

Geometry of SWCNT : 2/3 are semiconducting 1/3 metallic
conductivity tunable by functionalisation

Tremendous current carrying capacity: 1 billion A/cm²
no electromigration

Excellent heat conductor: twice as good as diamond

- Mechanical properties:

Strength at least 100 times higher than steel

Young's modulus 1 Tpa

Flexible  ultimate carbon fibers

- Geometry:

High aspect ratio of 1000 to 10000, 1 nm diameter

High surface area

Why carbon nanotubes?

Ideal one dimensional (1D) objects

Outstanding mechanical and electronic properties

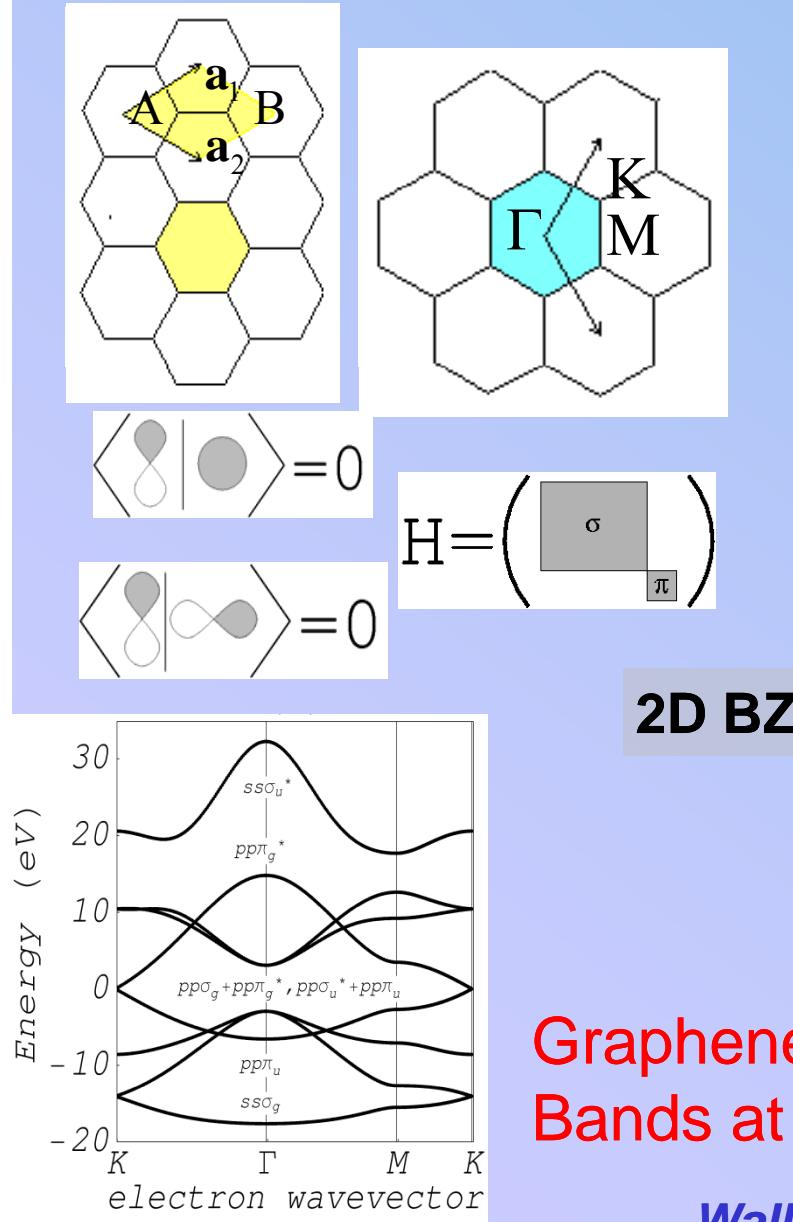
Ideal systems to study basic correlation effects
even at room temperature!!!!!!

Tunable electronic properties and interactions!!!

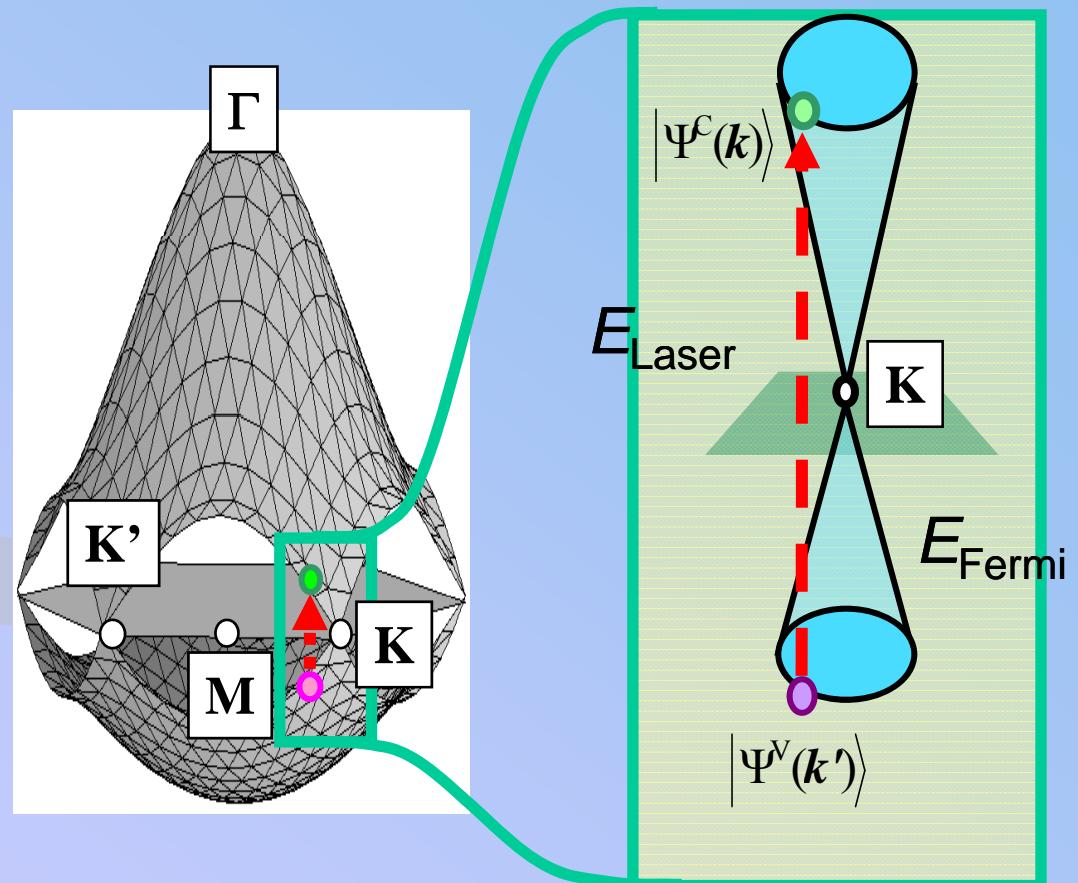


Interesting for both
basic research and applications

ELECTRONIC STRUCTURE OF GRAPHENE



$$H(\mathbf{k})c^j(\mathbf{k}) = E^j(\mathbf{k})S(\mathbf{k})c^j(\mathbf{k})$$



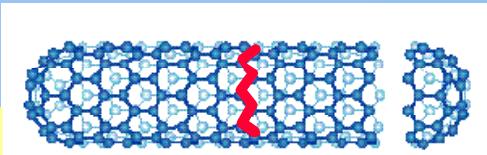
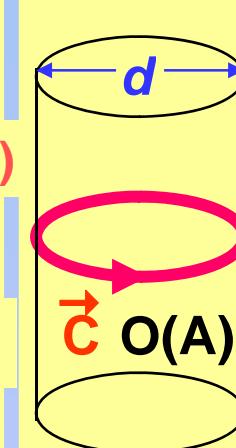
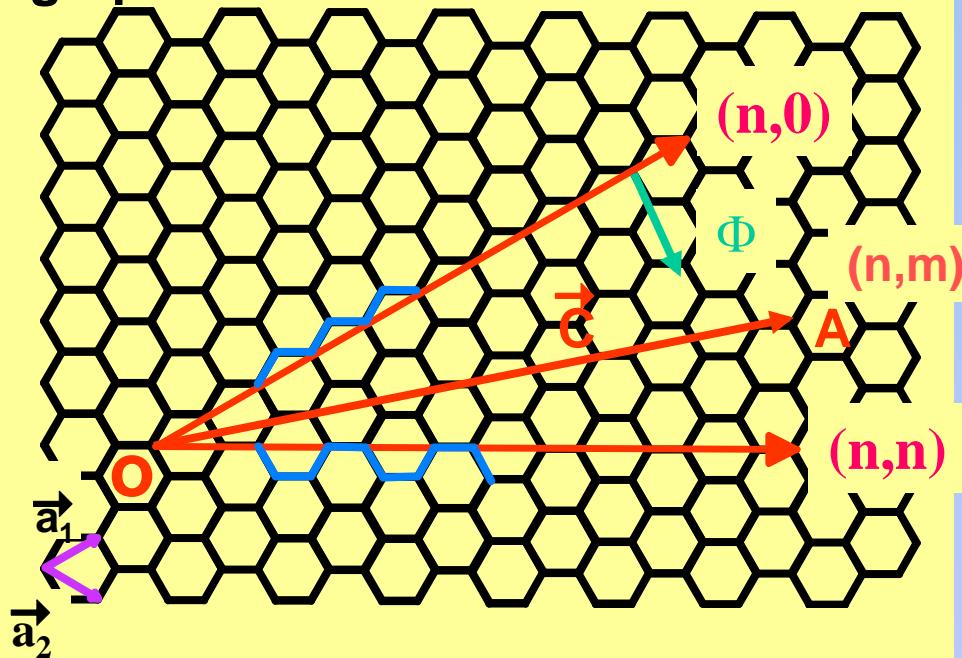
Graphene is a zero gap semiconductor.
Bands at Fermi energy (K point) are linear in k.

Wallace (1947), Painter and Ellis PRB 1, 4747 (1970)

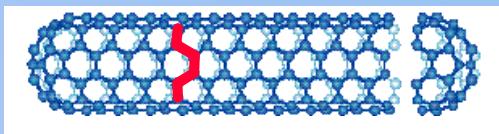
Introduction: formation of single-wall carbon nanotubes

SWCNTs $\vec{C} = n\vec{a}_1 + m\vec{a}_2 = (n, m)$ **Hamada vector**

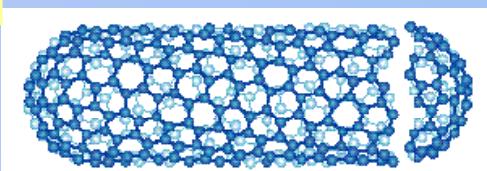
graphene sheet



(n,0) zigzag $\Phi=0$



(n,n) armchair $\Phi=30^\circ$



(n,m) chiral $0 < \Phi < 30^\circ$

The diameter of SWCNT (n,m):

$$d_{n,m} = |\vec{C}|/\pi = a(n^2 + m^2 + mn)^{1/2}/\pi$$

$$a = |a_1| = |a_2| = 2.49 \text{ \AA}$$

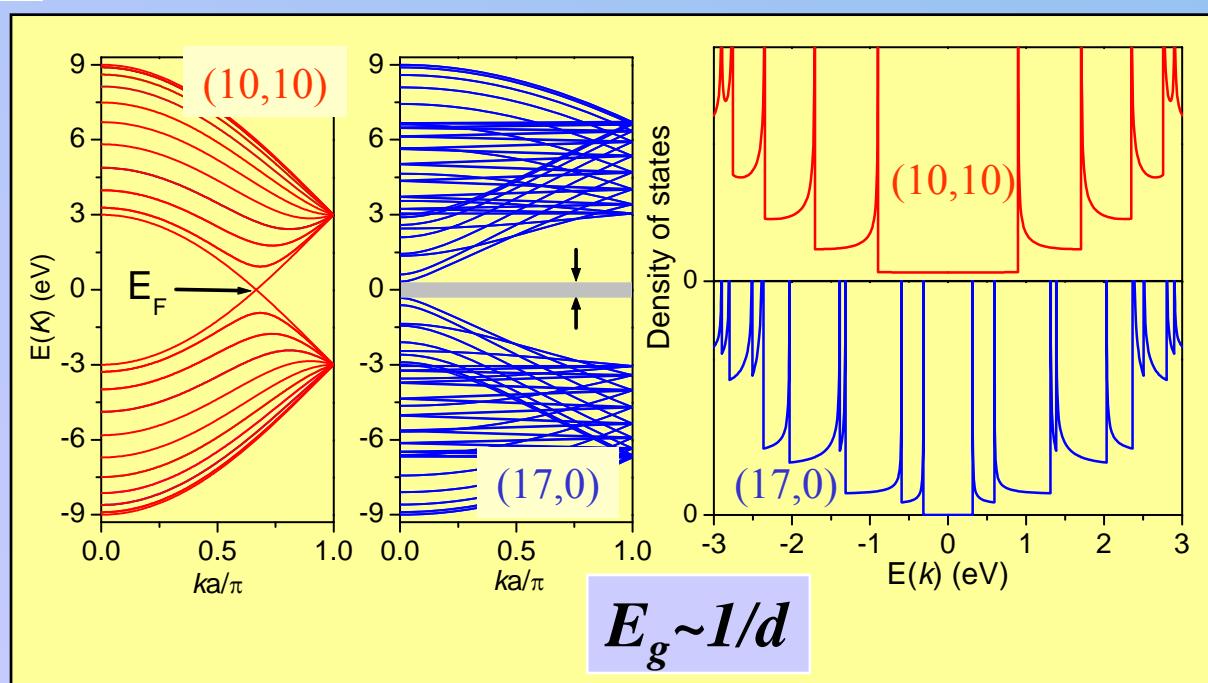
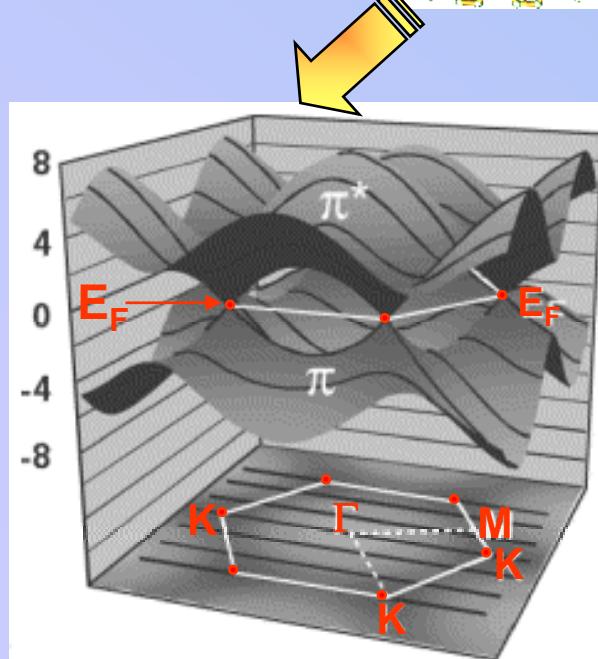
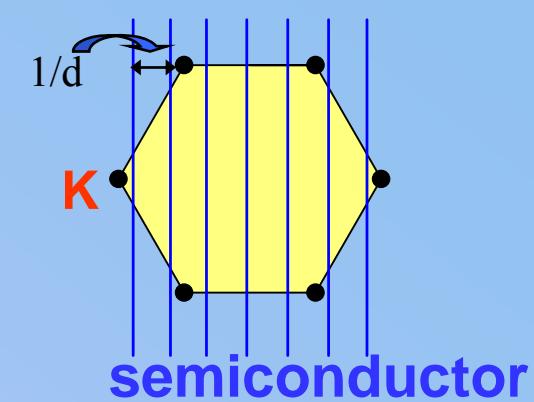
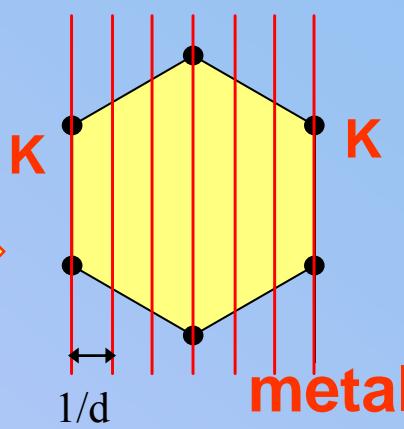
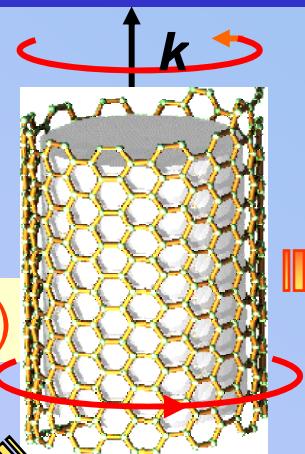
e.g. (10,10) $d=1.37 \text{ nm}$

Introduction: electronic properties of SWCNTs

periodic boundaries

$$2\pi d \cdot k_n = 2\pi \cdot n$$

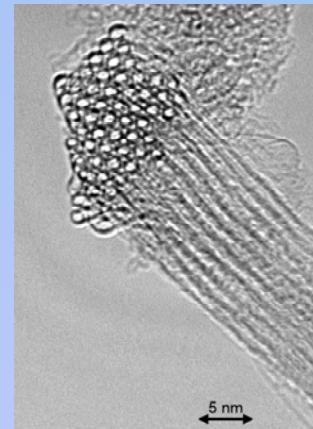
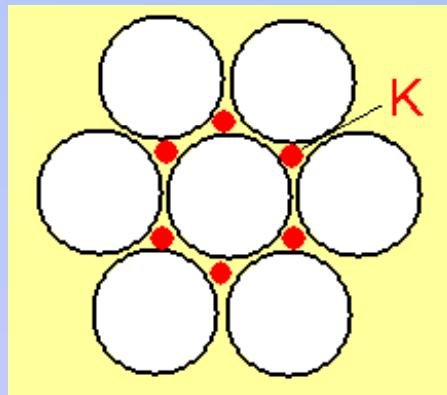
(k_n -quantization)



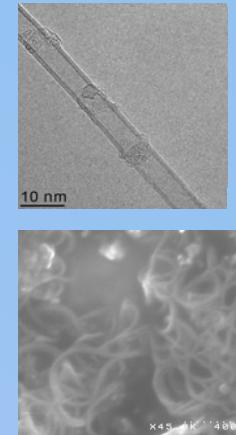
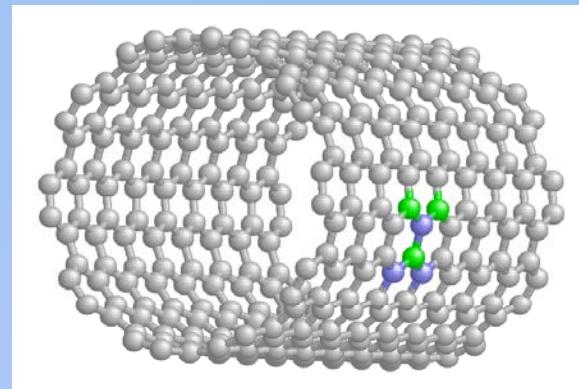
Electronic properties are determined by the structure of SWCNT
using tight binding approximation

Functionalization: doping and filling of SWCNT

● Intercalation:



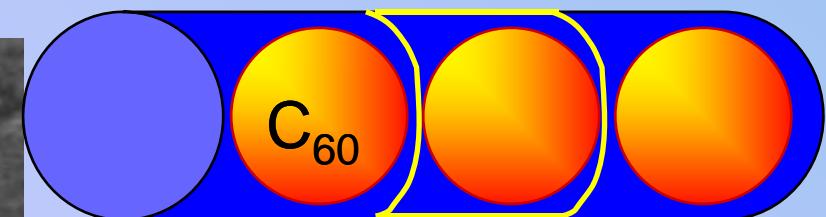
Substitution:



e.g. K_x -SWCNT

BC_3 -SWNT, MWBNT, SiCNT

● Filling: C_{60} Peapods, DWCNT....

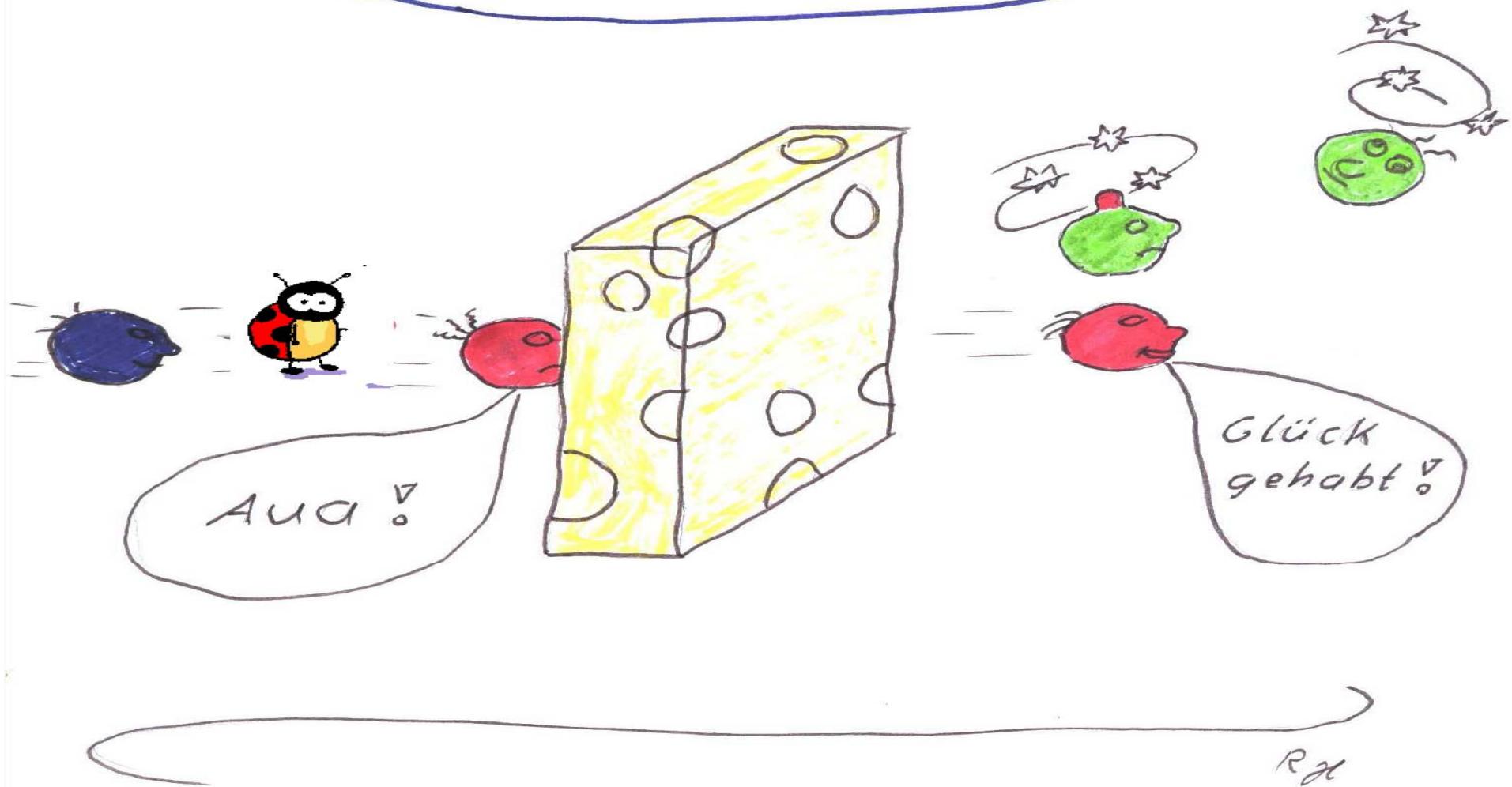


→ Optical and high energy spectroscopy are ideal tools!

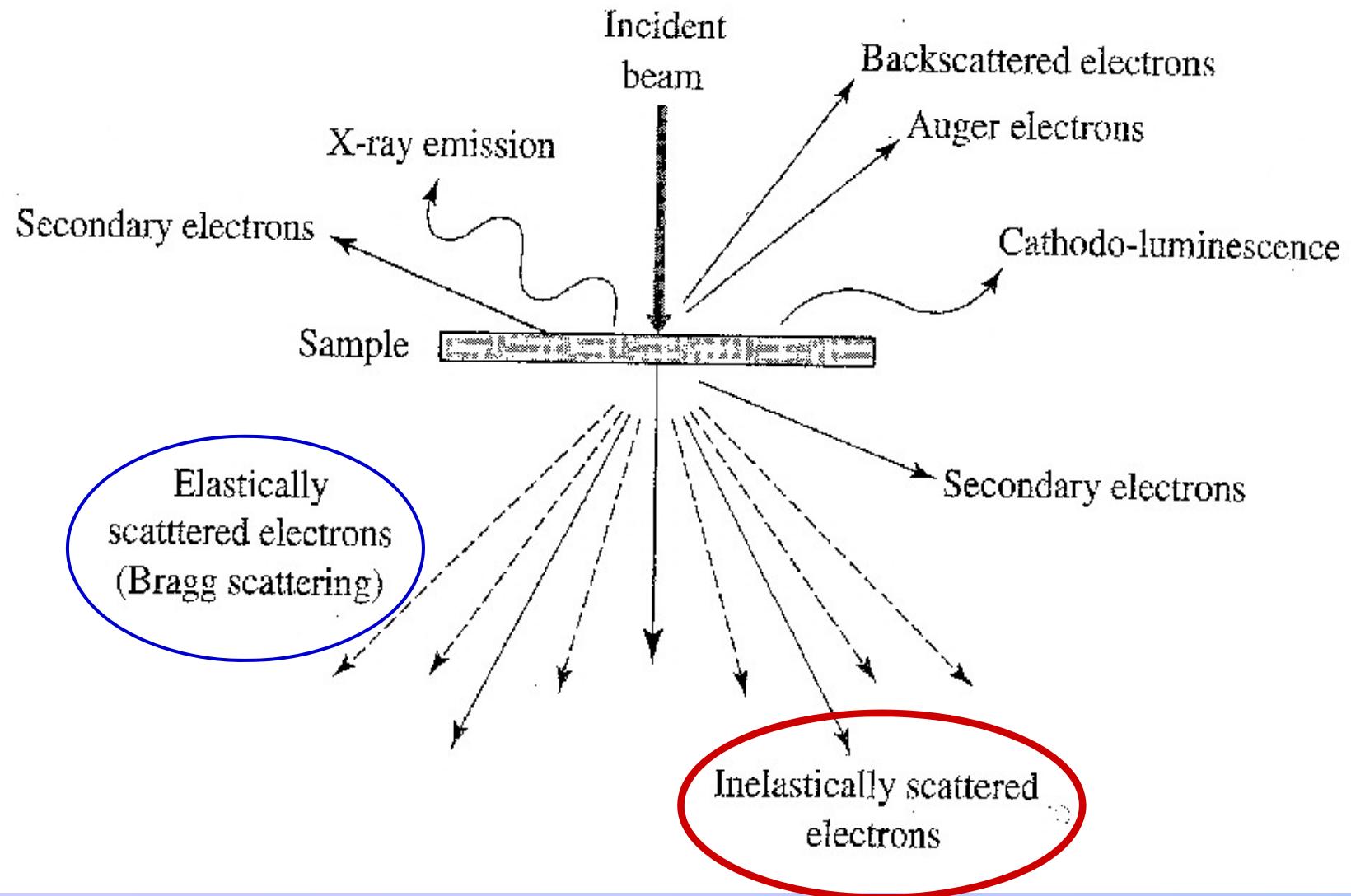
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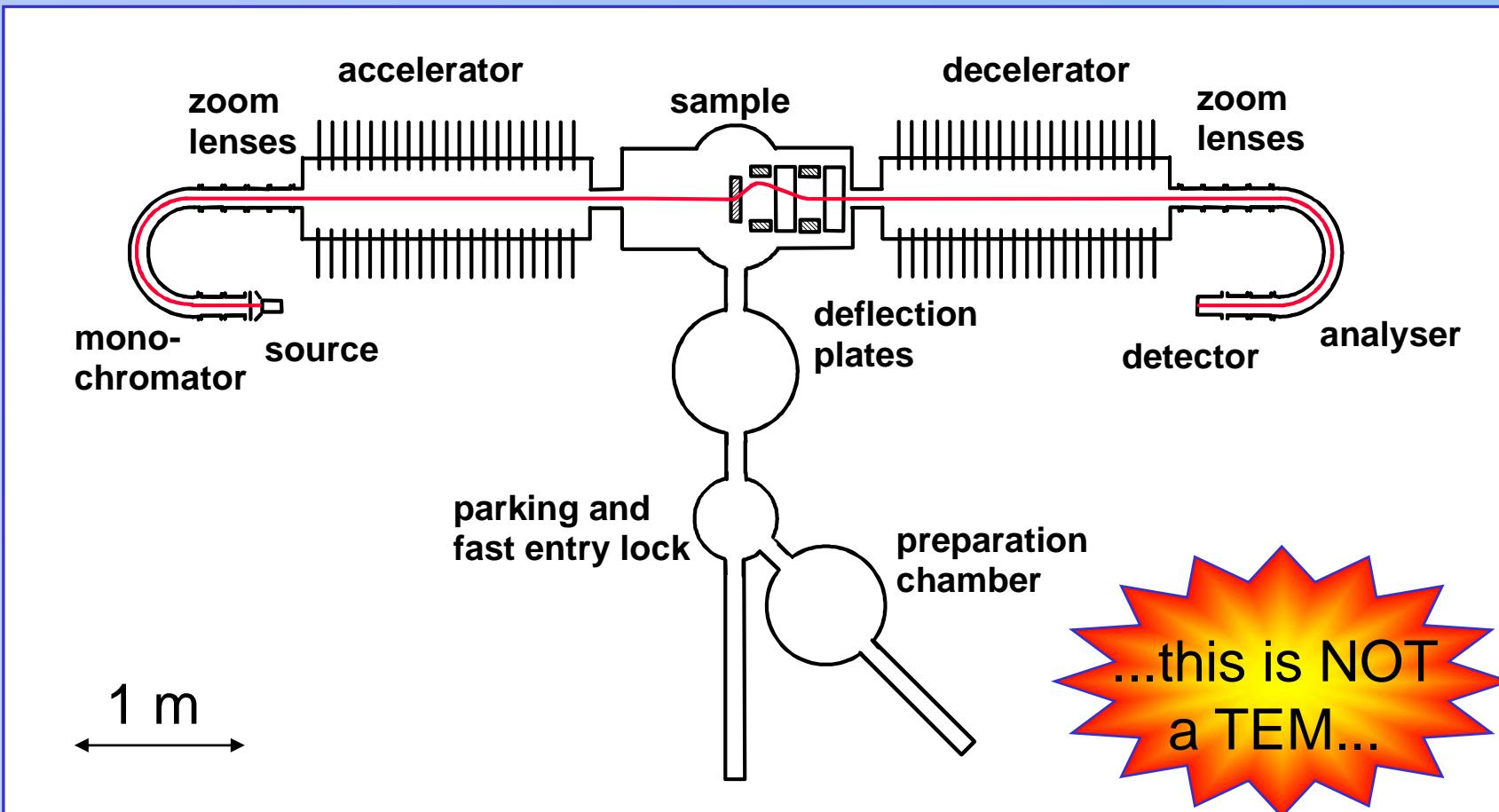
EEELS



Transmission of fast electrons through a solid: *what can happen ?*



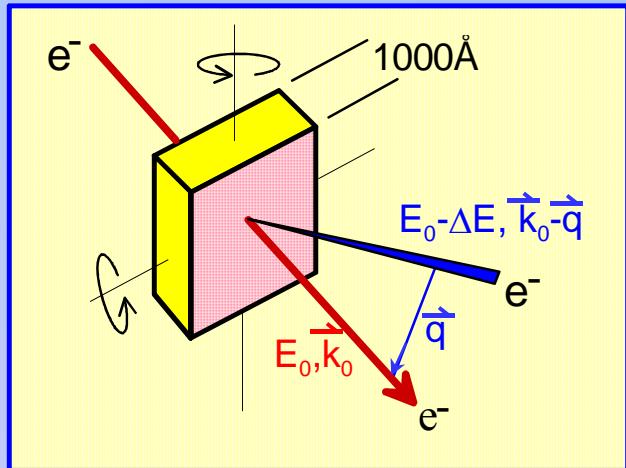
EELS in transmission system: purpose built



$$E_0 = 170 \text{ keV} \quad \Delta E = 100 \text{ meV} \quad \Delta q = 0.03 \text{\AA}^{-1}, \\ k_0 = 228.4 \text{\AA}^{-1} \quad \lambda = 0.0275 \text{\AA}, \quad \theta = 0.25^\circ$$

J. Fink, Adv. Electron. Electron Phys., 75, 121 (1989)

What is measured: EELS in transmission

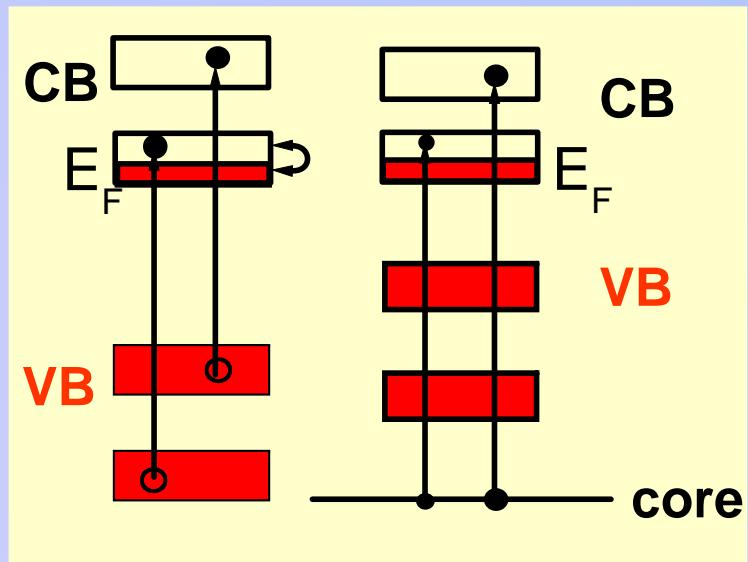


- **Transmission:**
 $d \sim 100\text{nm}$
bulk sensitive probe

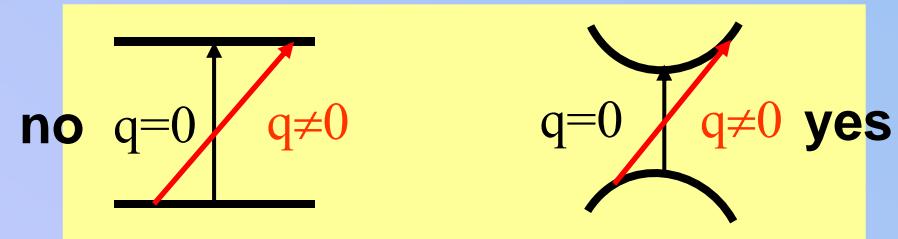
- **Structural analysis**

$\Delta E = 0$: Electron diffraction

Excitations



- Electronic properties
loss function = $\text{Im}[-1/\epsilon(\omega, q)]$
- Optical properties vs. q dependence



dispersion: curvature of bands

$$E_0 = 170 \text{ keV} \quad \Delta E = 180 \text{ meV} \quad \Delta q = 0.03 \text{ \AA}^{-1}$$

Some theoretical aspects

The cross section

$$\frac{d^2\sigma}{d\Omega dE} = \left(\frac{d\sigma}{d\Omega} \right)_{\text{Ruth}} S(q, \omega)$$

$$\left(\frac{d\sigma}{d\Omega} \right)_{\text{Ruth}} = \frac{4}{a_0^2 q^4}$$

Rutherford scattering

$$S(q, \omega) = \frac{1}{2\pi\hbar N} \int e^{-i\omega t} \langle n_q(t) n_{-q}(0) \rangle dt$$

**dynamic structure factor
(system response)**

$n_q(t)$ = density operator

compare to elastic scattering (form factor):

$$S(q, \omega=0) = S(q) = (Z - F(q))^2$$

Some theoretical aspects

Connection to the dielectric properties

$$\frac{d^2\sigma}{d\Omega dE} = \left(\frac{d\sigma}{d\Omega} \right)_{\text{Ruth}}$$

$S(q, \omega)$

$$= \frac{\hbar}{(\pi e a_0)^2} \frac{1}{q^2}$$

$$\text{Im} \left[-\frac{1}{\epsilon(q, \omega)} \right]$$

loss function

$\epsilon(q, \omega)$ = dielectric function = $\epsilon_1 + i\epsilon_2$

$$\text{Im}[-1/\epsilon] = \frac{\epsilon_2}{\epsilon_1^2 + \epsilon_2^2}$$

Some theoretical aspects

EELS probes longitudinal excitations

- ♦ the quantity measured in EELS is $\frac{1}{q^2} \text{ Im} \left[-\frac{1}{\epsilon(q, \omega)} \right]$
- ♦ $\text{Im} \left[-\frac{1}{\epsilon(q, \omega)} \right]$ has a maximum when ϵ small (zero)

|| → fast electrons excite longitudinal waves !

in addition: $\mathbf{k} \times \mathbf{E} = \omega \mathbf{H} = \mathbf{0}$

$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{H}}{\partial t} = \mathbf{0}$ || → no magnetic contributions

Some theoretical aspects

Kramers-Kronig relations and sum rule

- Kramers-Kronig relation:

$$\operatorname{Re} \left[\frac{1}{\varepsilon(\mathbf{q}, \omega)} \right] - 1 = \frac{1}{\pi} P \int_{-\infty}^{\infty} \frac{d\omega'}{\omega' - \omega} \operatorname{Im} \left[\frac{1}{\varepsilon(\mathbf{q}, \omega)} \right]$$

- sum rule:

$$\int_0^{\infty} \omega \operatorname{Im} \left[\frac{1}{\varepsilon(\mathbf{q}, \omega)} \right] d\omega = \frac{\pi}{2} \omega_p^2, \quad \omega_p^2 = \frac{n e^2}{\varepsilon_0 m}$$

Plasma frequency

BUT!!! Macroscopic not always = Microscopic ε

Macroscopic ε

Regular case:

macroscopic isotropic bulk electron density, crystal symmetry
cancels off-diagonal elements

$$\varepsilon = \varepsilon_M (\lVert \vec{r}_i - \vec{r}_j \rVert)$$

$$E(\vec{q}, \omega) = \varepsilon_M^{-1}(\vec{q}, \omega) D(\vec{q}, \omega)$$

Random Phase Approximation

Quasi free particles in a mean field. Works well if electrons are reasonably screened

$$\varepsilon_M(\vec{q}, \omega) \approx \varepsilon_{RPA}(\vec{q}, \omega) = 1 - V_q \chi_0(\vec{q}, \omega)$$



The Drude-Lorentz dielectric function

Solution for damped and driven harmonic oscillator

$$x = \frac{eE_0 e^{i\omega t}}{m(\omega_T^2 - \omega^2 - i\omega/\tau)}$$

and

$$P = n e x = \chi_e \epsilon_0 E$$

and

$$\epsilon = (1 + \chi_e)$$

$$\epsilon(\omega) = 1 + \frac{ne^2}{\epsilon_0 m} \frac{1}{\omega_T^2 - \omega^2 - i\omega/\tau}$$

Dielectric function with **real** and **imaginary** part

$$\omega_p^2 = \frac{ne^2}{\epsilon_0 m}$$

plasma frequency scales with charge density

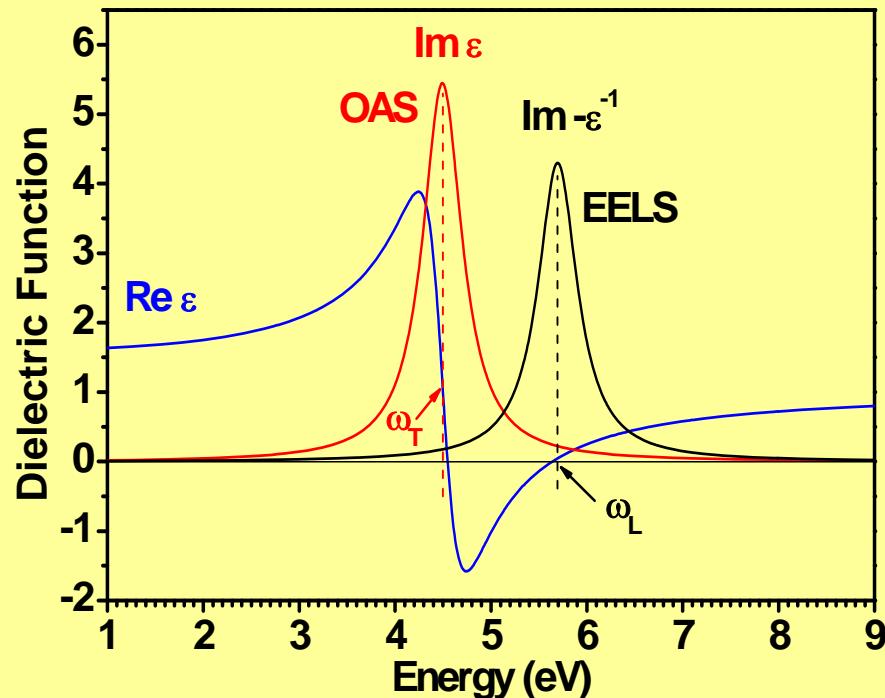
$$\omega_T = 0$$

no spring constant, free electrons, **Drude** model

Drude-Lorentz dielectric function and spectroscopy

High plasma frequency ω_p

$$\varepsilon_\infty = 1 \quad \omega_p = 3.5 \quad \omega_T = 4.5 \quad \Gamma = 0.5$$



ω_T transversal plasma frequency
resonant absorption of light

Optical Absorption Spectroscopy
For thin samples OAS $\sim \text{Im}(\varepsilon)$

ω_L longitudinal plasma frequency
resonant energy loss of e^-

Electron Energy Loss Spectroscopy
EELS $\text{Im}(-\varepsilon^{-1})$

$$\varepsilon = \varepsilon_\infty + \sum_i \frac{\omega_{p,i}^2}{(\omega_{T,i}^2 - \omega^2) + i\omega\Gamma_i}$$

$$\omega_L^2 = \omega_T^2 + \omega_p^2$$

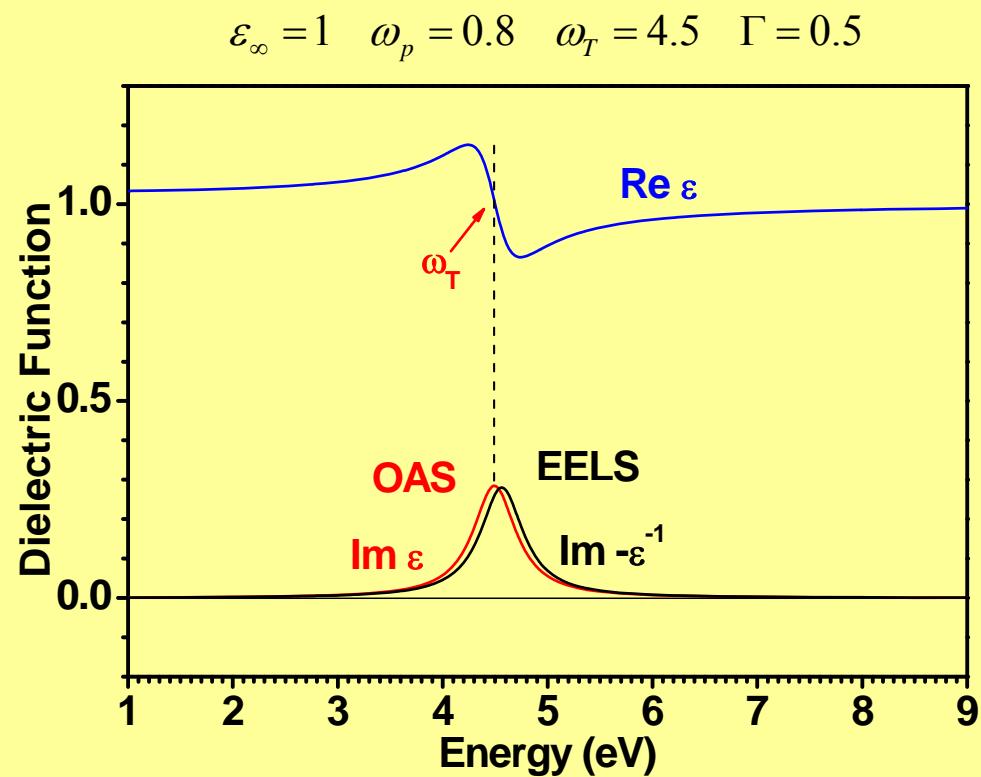
$$\omega_L > \omega_T$$

$$\text{EELS} > \text{OAS}$$

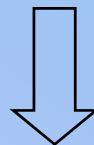
$$\Gamma = 1/\tau$$

Isolated systems and free objects

Low plasma frequency ω_p
(everything else the same)



Weak ω_p (low e^- density)
causes only weak
modulation in $\text{Re}(\varepsilon)$



ω_L and ω_T are very close

$$\omega_L^2 = \omega_T^2 + \omega_p^2$$

$\omega_L \sim \omega_T$

EELS~OAS

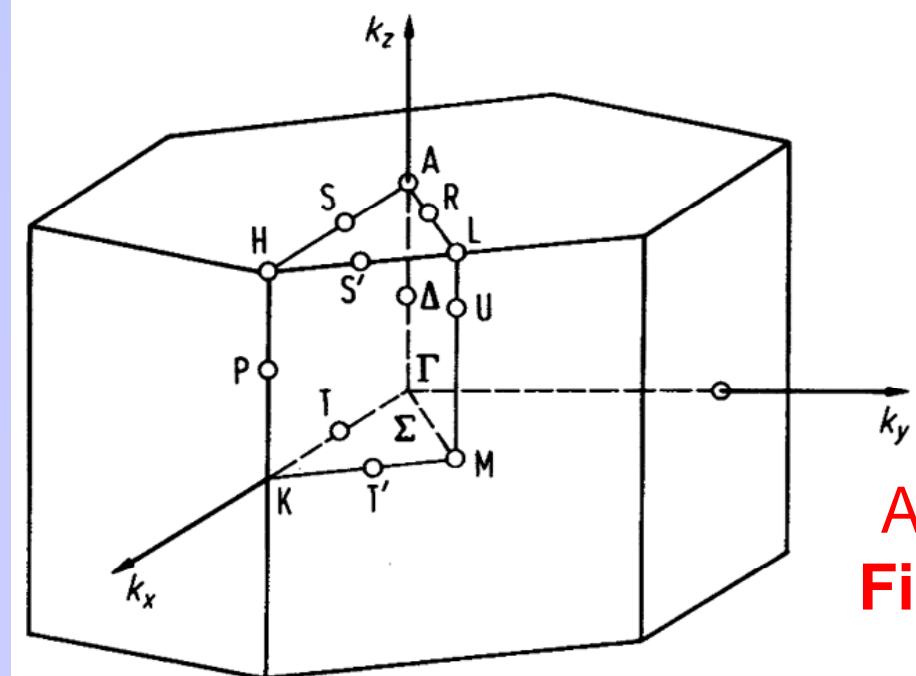
OAS (ω_L) and EELS (ω_T) are
degenerate in diluted systems

EELS Beyond the optical range ($q>0$)

visible light 0.0001 \AA^{-1} direct transitions

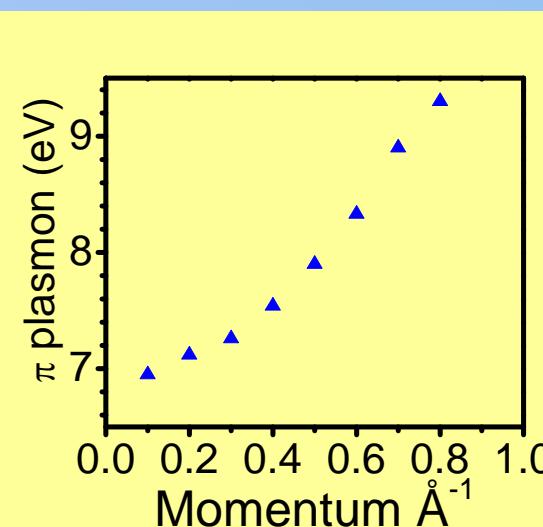
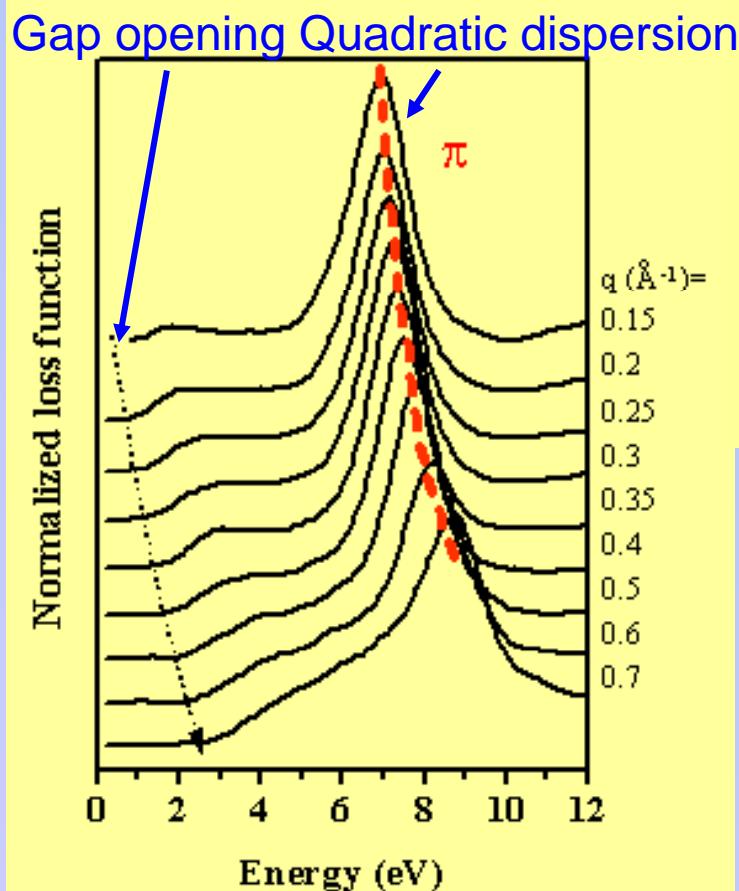
Brillouin zone 3 \AA^{-1} (100) Bragg

172 keV e⁻ 228 \AA^{-1} **whole zone**

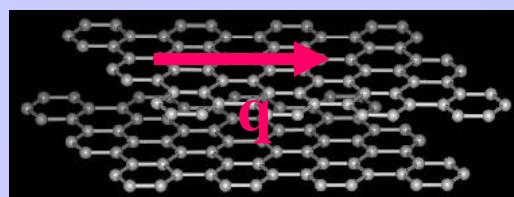
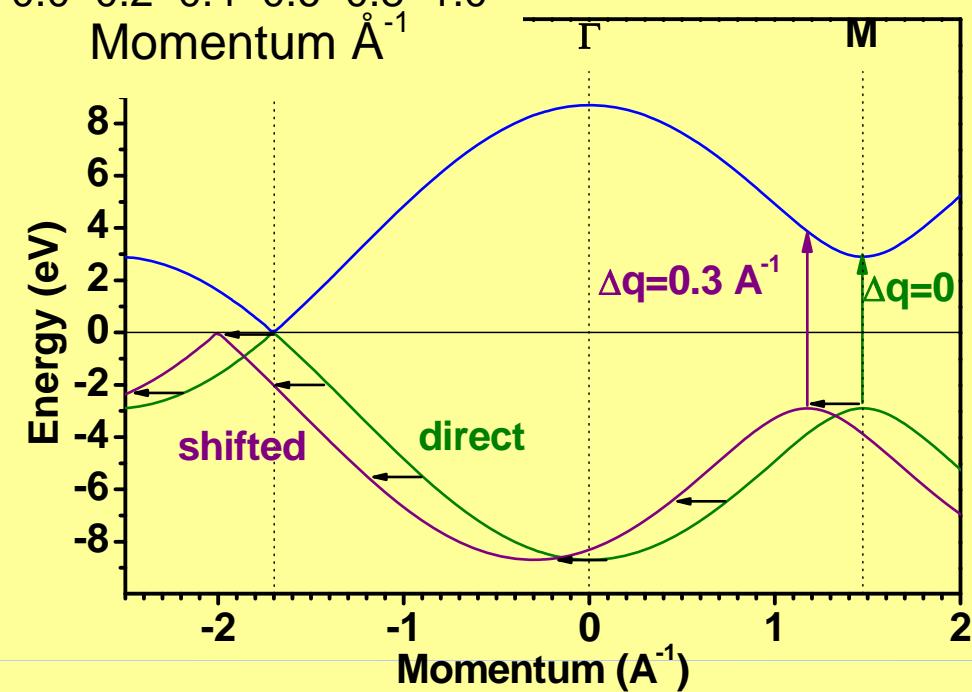


AR-EELS probes the dispersion
Fingerprint of solids/molecules

Example 1: Graphite



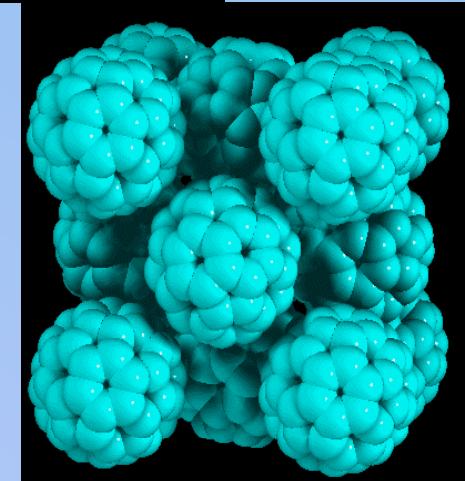
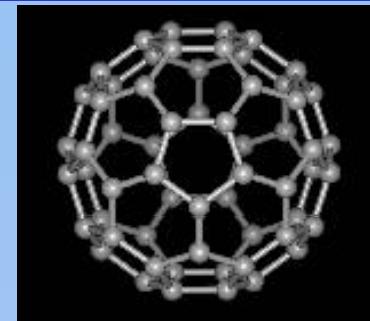
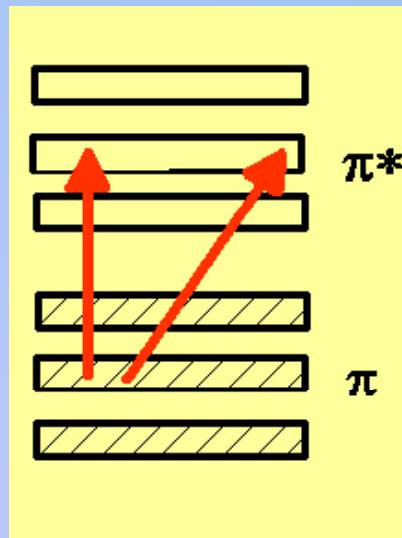
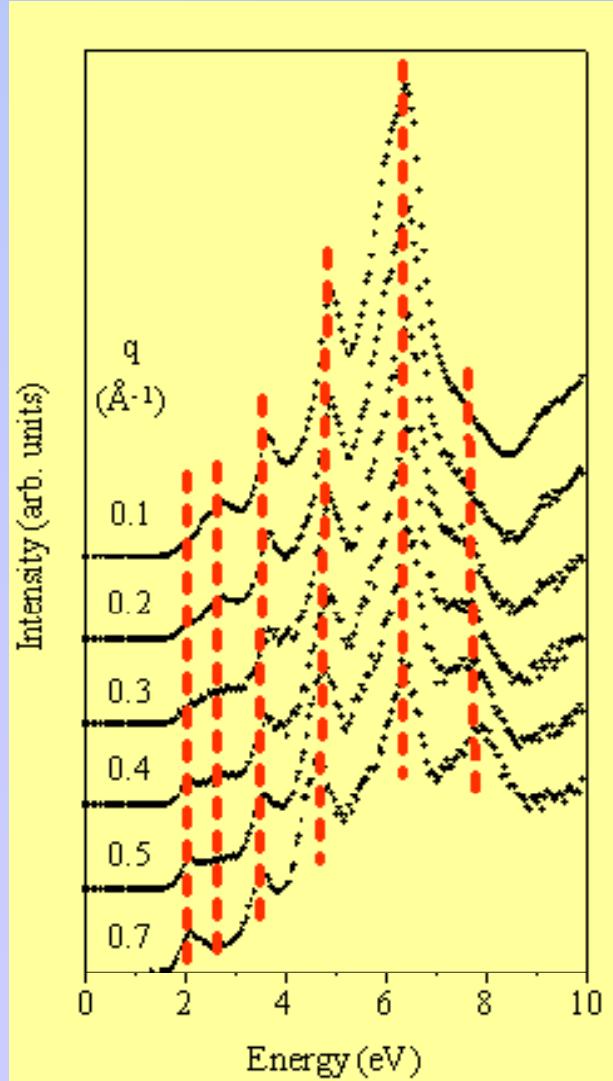
π plasmon
(M point)
3D system



A. G. Marinopoulos *et al*, PRL, 89, 076402, (2002)

Example 2: Fullerenes

EELS of C₆₀: „0-D“ system

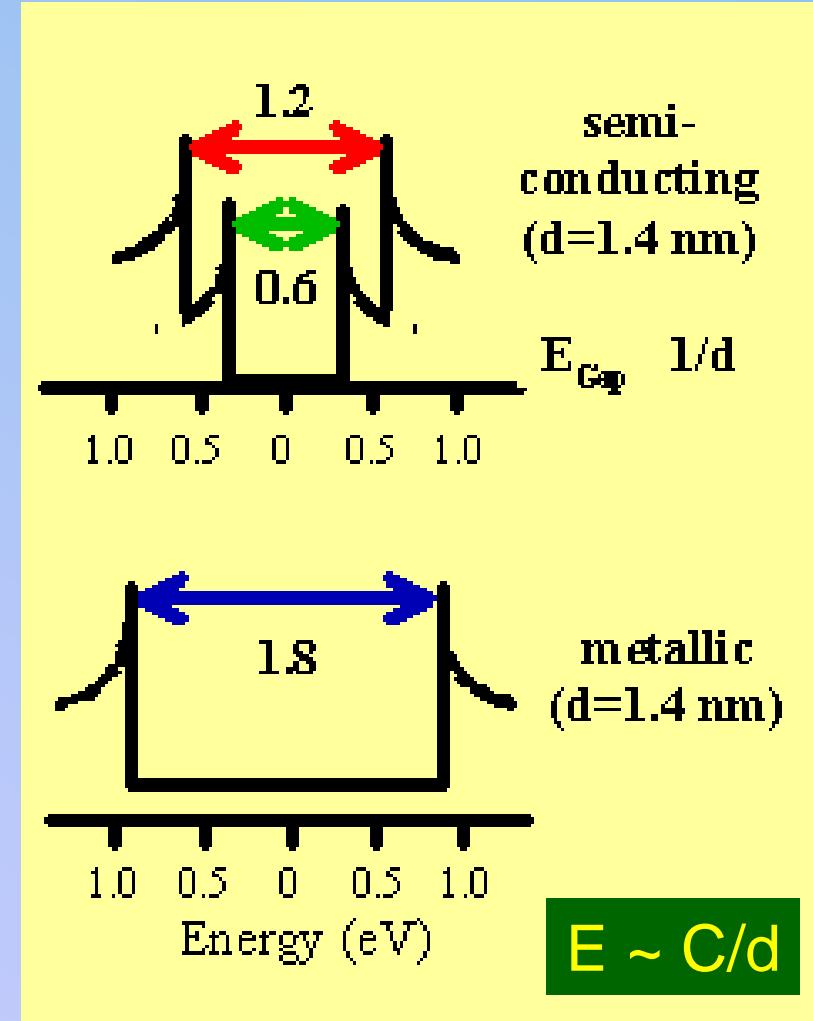
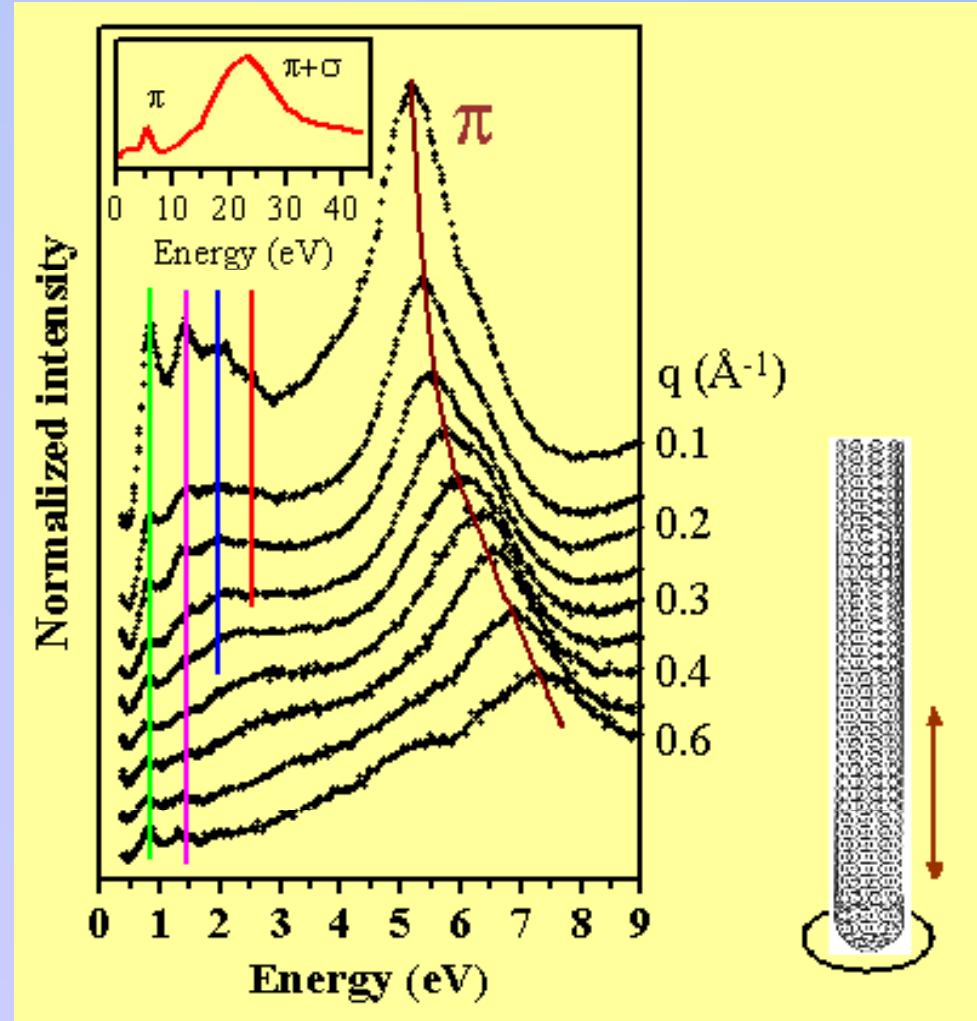


Transitions between flat
 π -Bands - no dispersion.

fcc C₆₀ ist a moleculares
(0D) system.

Example 3: bundled SWCNT

EELS of SWCNT bundles : quasi 1D system



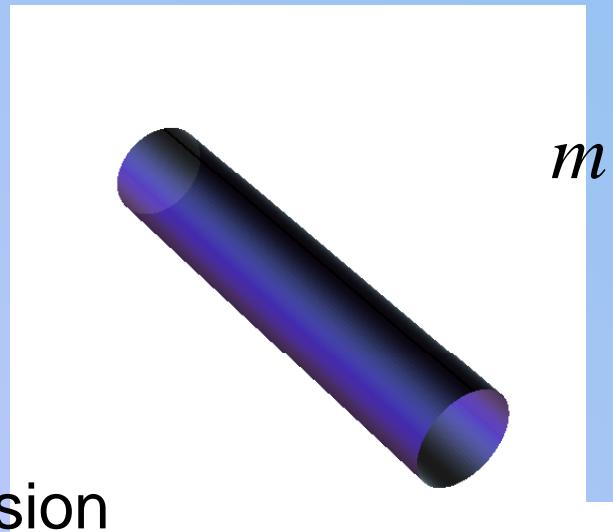
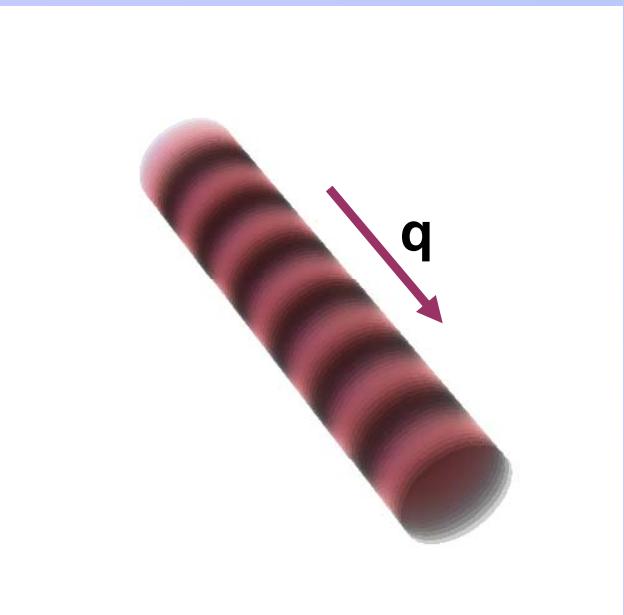
Plasmons in 1D

On-axis

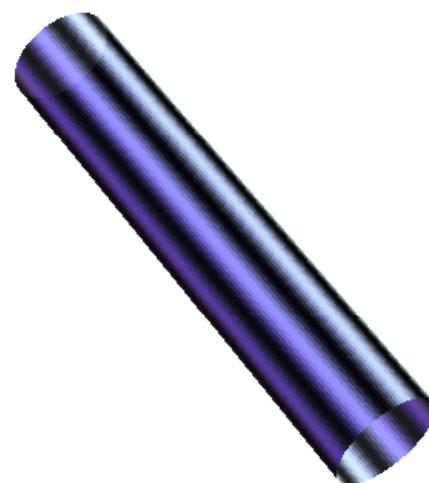
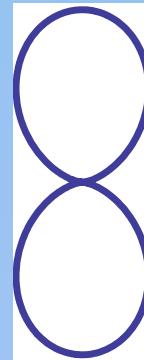
1D solid, dispersion

Circumferencial

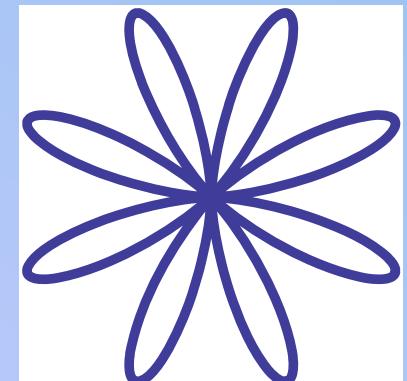
multipoles, no dispersion



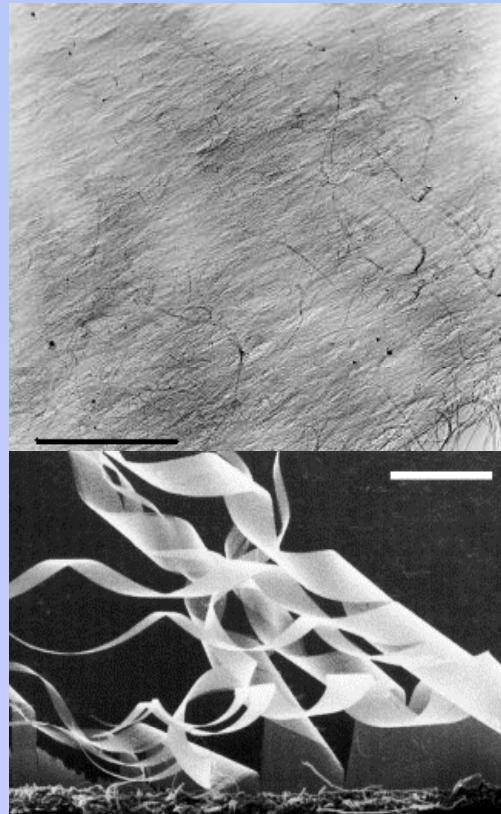
$$m = 2$$



$$m = 8$$



MA-SWCNT

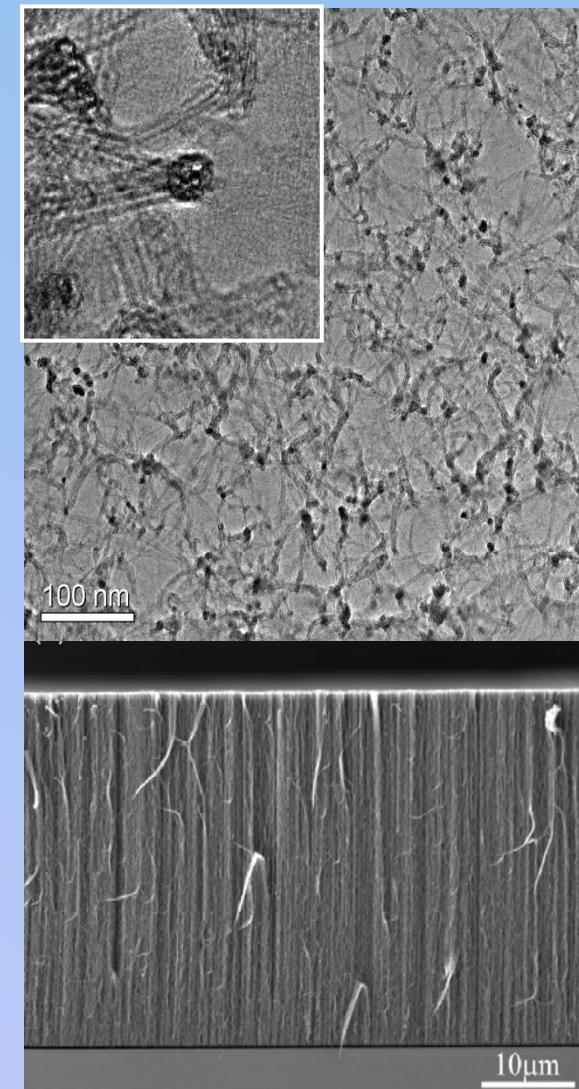


Aligned-SWCNT

Laser ablation
In plane
 $\sim 1.3 \text{ g/cm}^3$

CVD
Out of plane
 $\sim 0.05 \text{ g/cm}^3$

VA-SWCNT



Batch	Orientation	Method	Domain	σ	β	ξ
MA-SWNT	in-plane	ED	bulk	20°	19.2°	0.84
VA-SWNT	perpendicular	XAS	surface	37°	32.7°	0.56
VA-SWNT	perpendicular	OAS	bulk	26°	24°	0.75

E. Einarsson, C. Kramberger, *et al.* J.Phys. Chem. C, **111**, 17861, (2007)
X. Liu, *et al.* Synth. Met. **121**, 1183, 2001

sp² carbon at different densities

High Density

ω_p

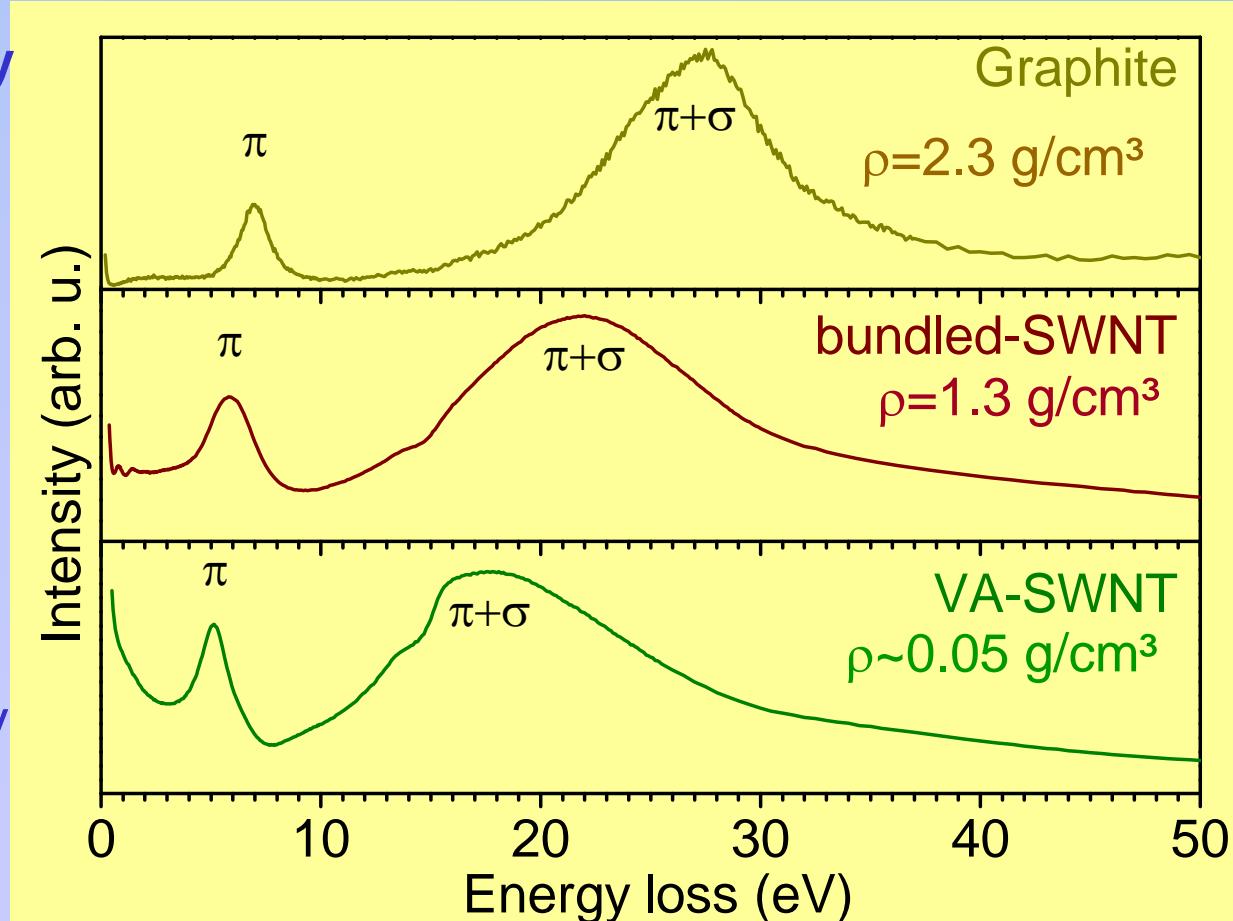
Bulk Graphite

Low Density

ω_p

Bundled SWCNT

Isolated SWCNT

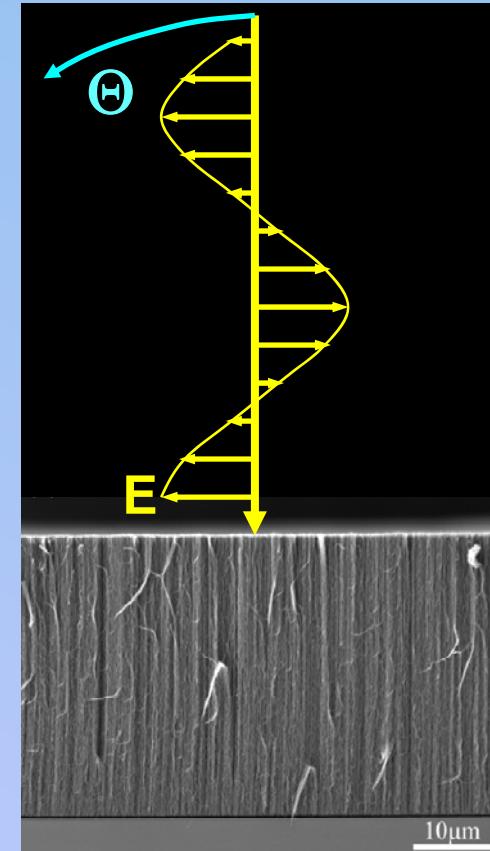
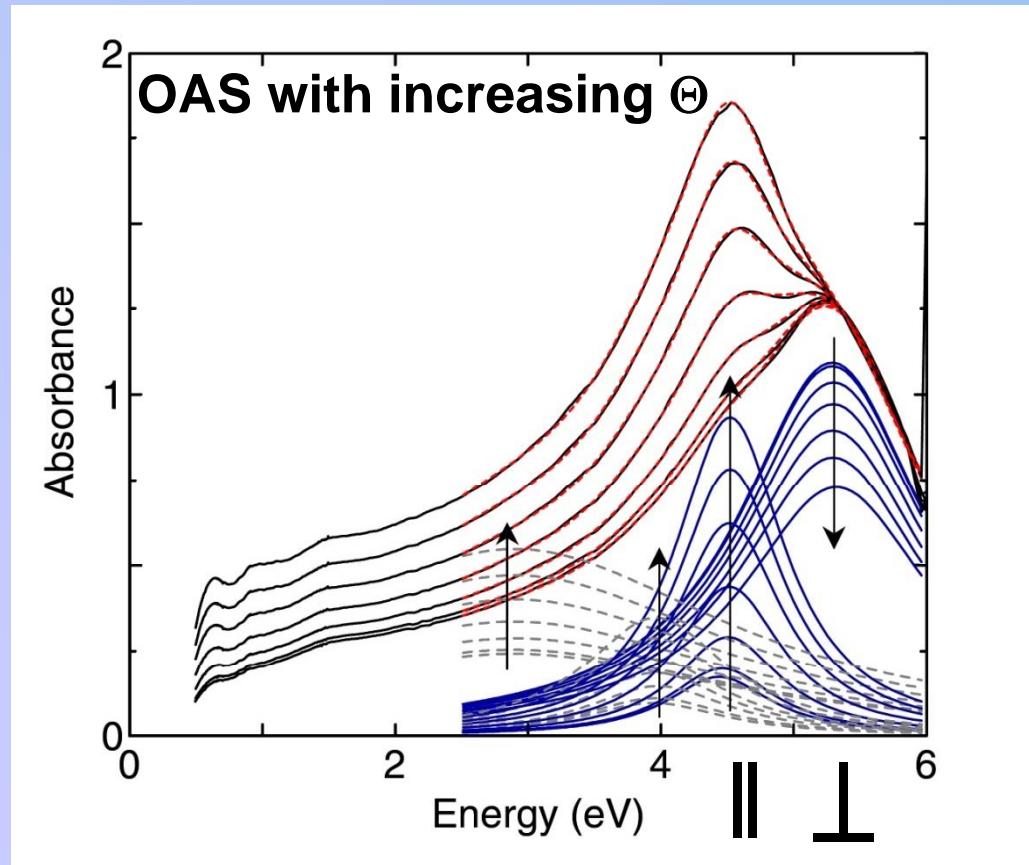


$$\omega_L^2 = \omega_T^2 + \omega_p^2$$

$$\omega_p^2 = \frac{ne^2}{m^* \epsilon_0}$$

But only one π plasmon!

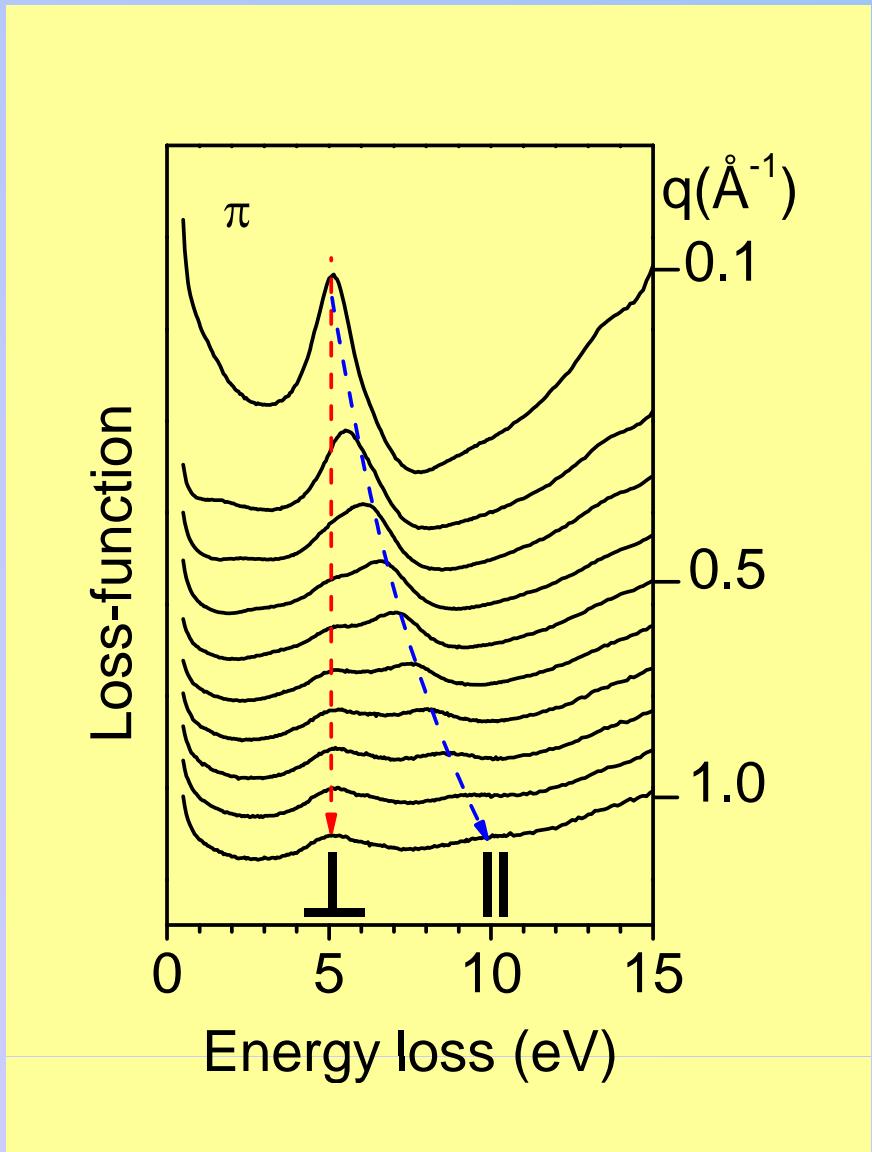
Optical absorption on VA-SWCNT



on-axis at 4.5 eV, crossed at 5.25 eV

Y. Murakami et al, PRL, 94, 087402, (2005)

π plasmon dispersion in isolated VA tubes

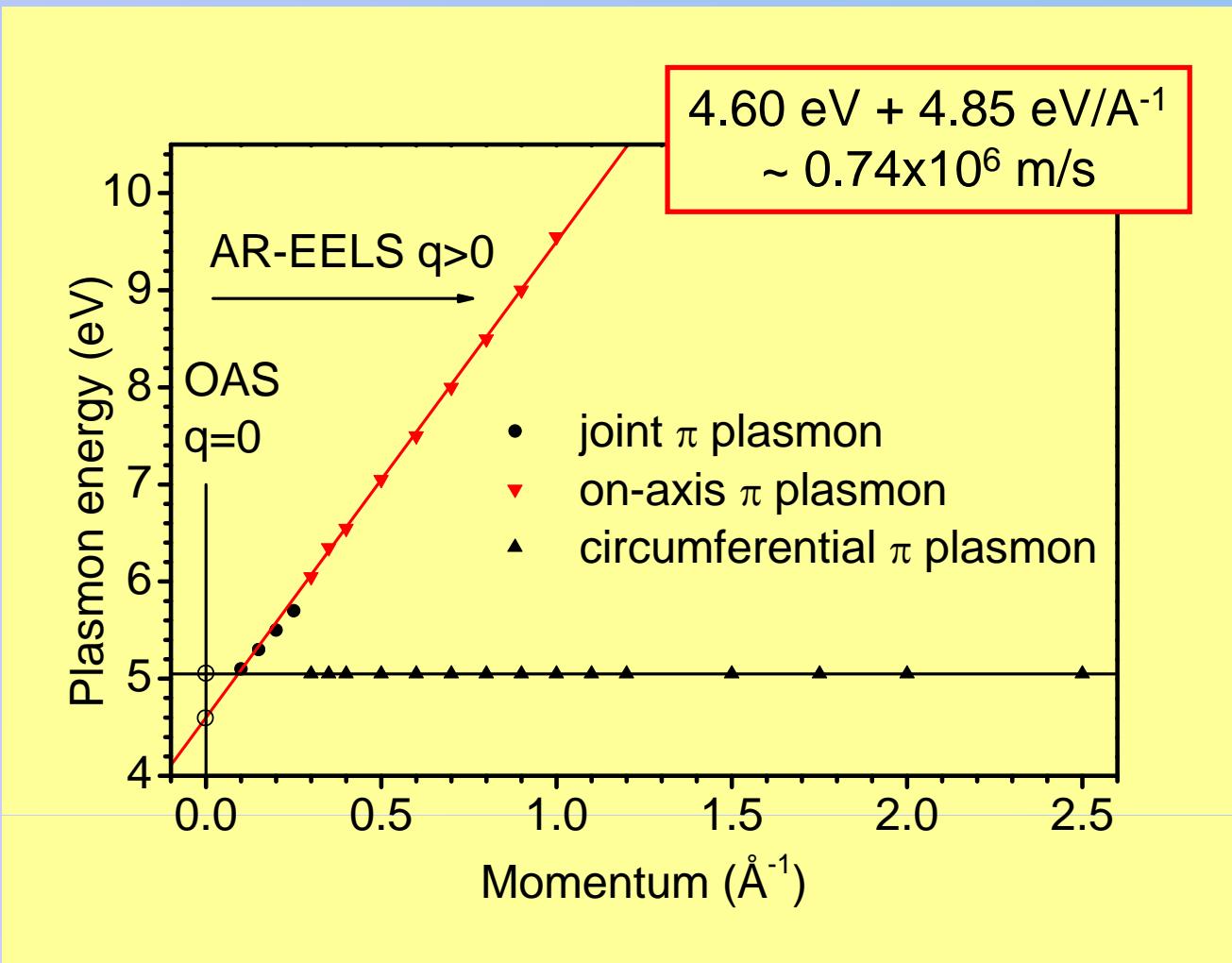


twofold π plasmon dispersion!

Perpendicular plasmon
No dispersion

On-axis plasmon
Dispersion

Twofold π plasmon dispersion paradox!



Linear dispersion!

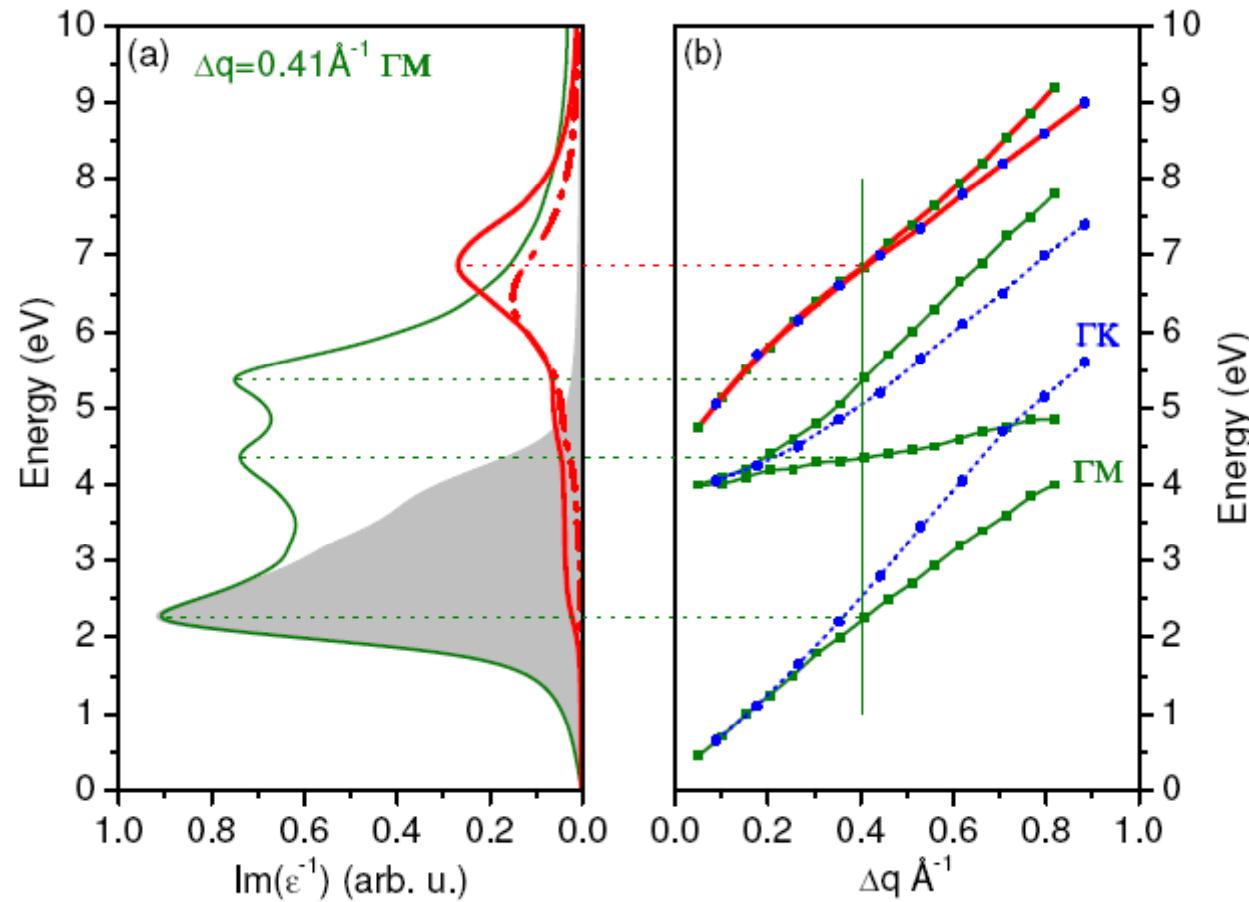
Is this π plasmon dispersion **REALLY** related to the K point in graphene?

optical limit $\sim 4.6 \text{ eV}$ parallel
 $\sim 5.1 \text{ eV}$ perpendicular



Excellent match to OAS

DFT calculations of loss function of graphenes including LFE

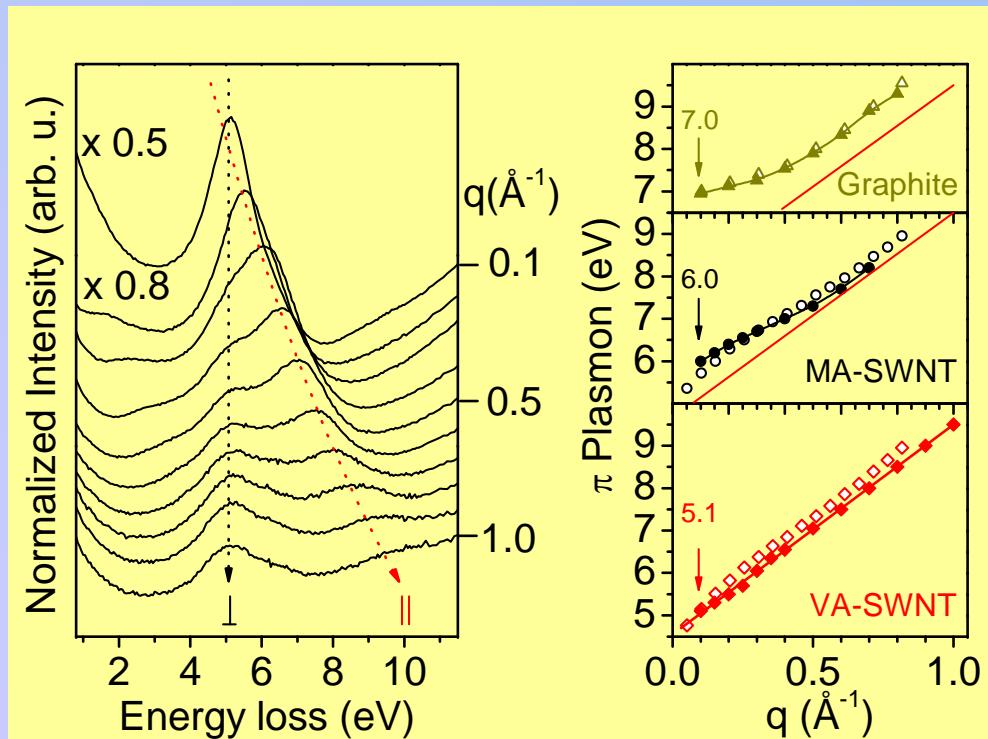


bare RPA (green line) and including depolarization effects (red line).

→ strong intermixing of excitations close to the Dirac point yields linear dispersion

C. Kramberger *et al*, PRL, 100, 196893 (2008)

Quantitative microscopic description



Sample

bulk graphite

bundled SWCNT

isolated SWCNT

Model

bulk graphite

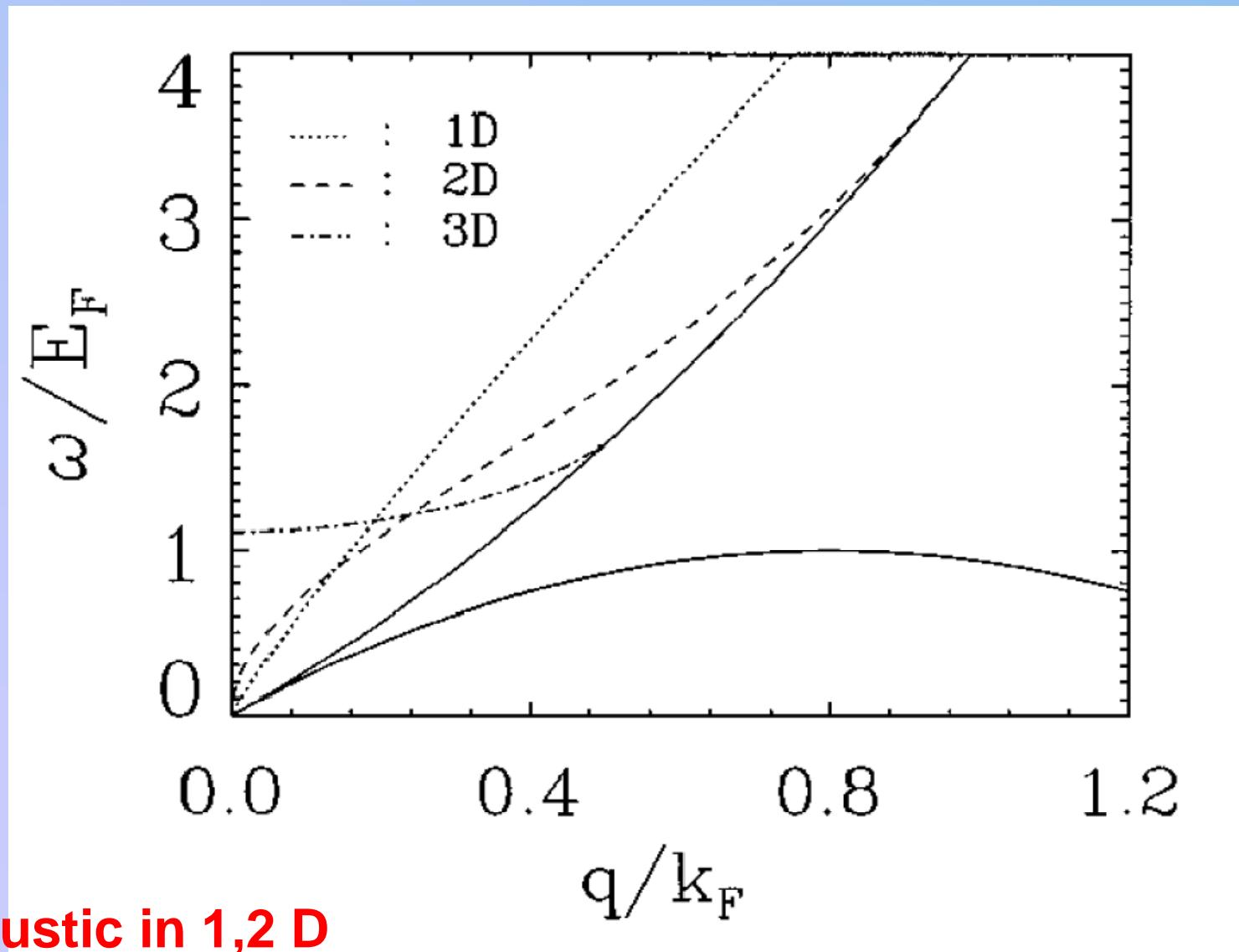
bi-layer

graphene

On-axis interband excitations are equivalent to in-plane excitations of isolated graphene!

C. Kramberger *et al*, PRL, 100, 196893 (2008)

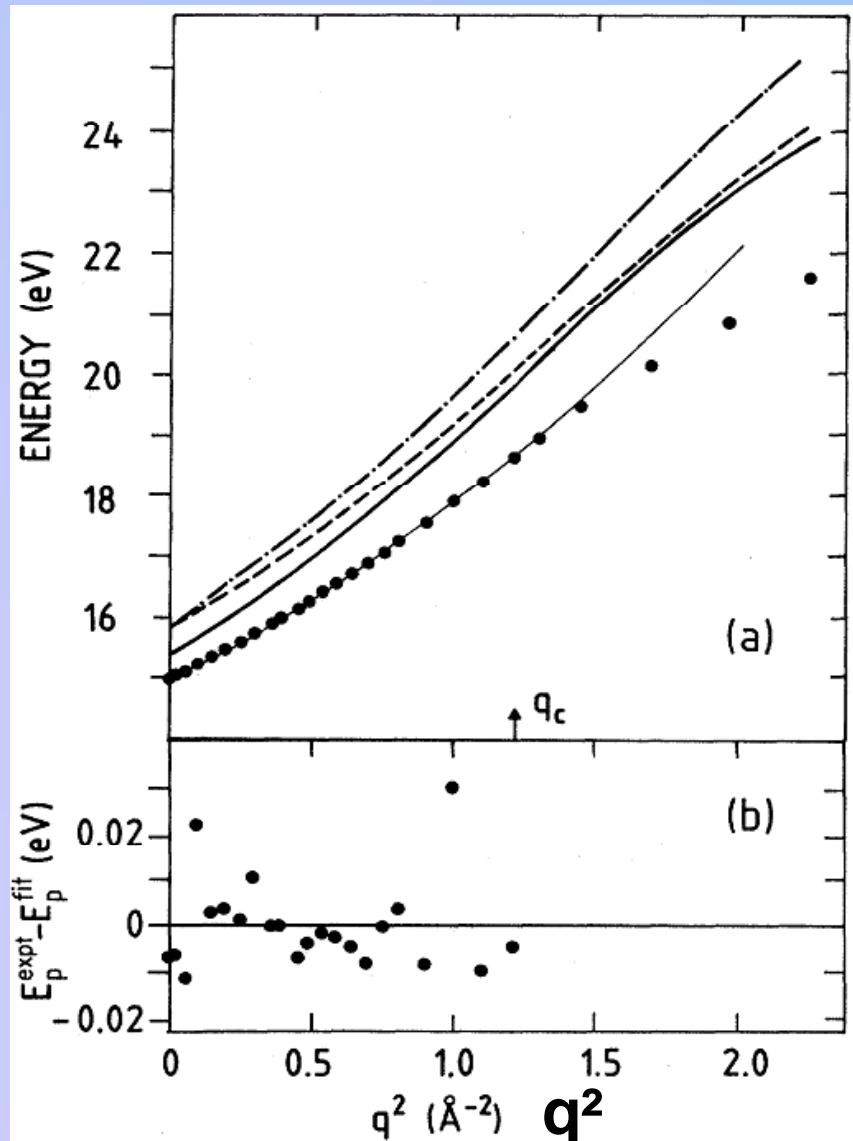
Charge carrier plasmon dispersion in different dimensions



Acoustic in 1,2 D

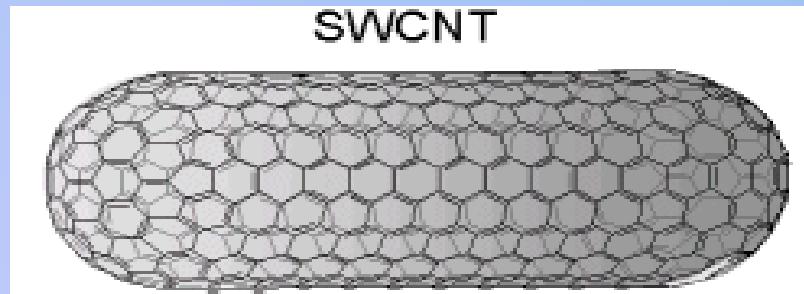
S. Das Sarma and E. H. Hwang PRB 54, 1936 (1996)

Charge carrier plasmon in 3D: Example Al

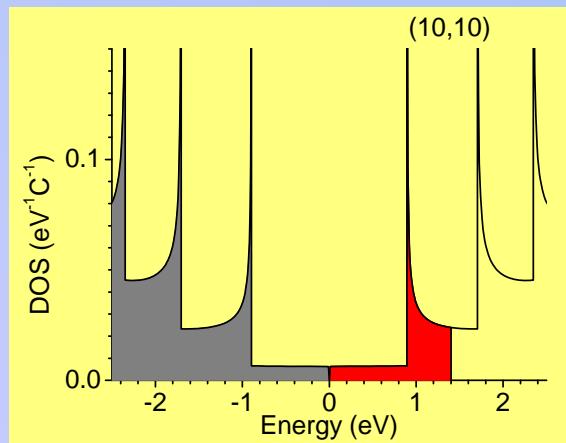
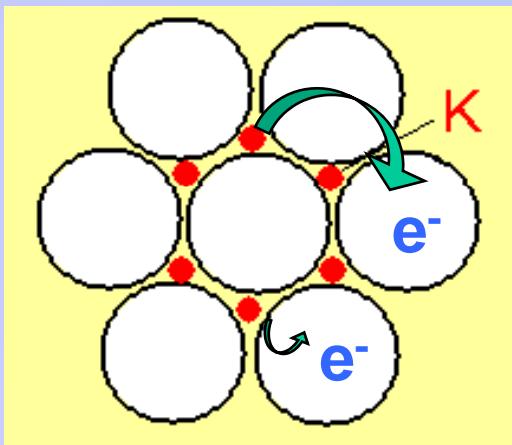
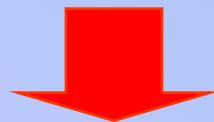


Quadratic dispersion in 3D metal!

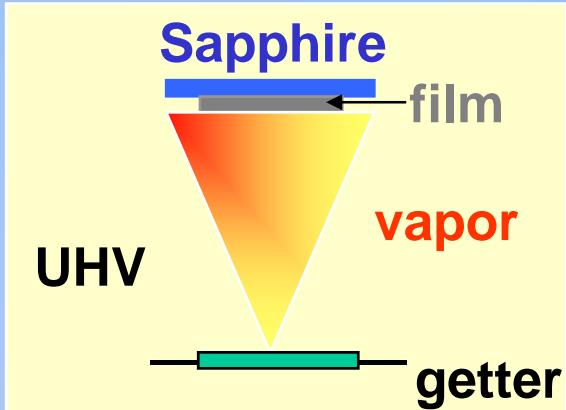
Intercalation



doping



Alkali-metal intercalation

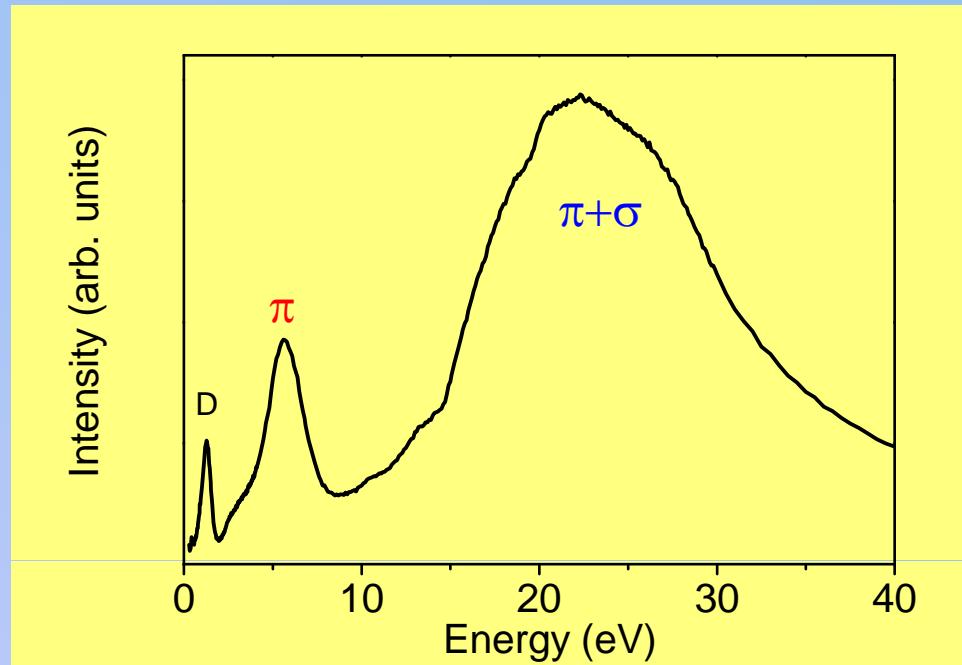


in-situ doping:
UHV evaporation
(5×10^{-10} mbar)
Na, K, Rb, Cs and
Ba SAES getters

Analysis of the loss function in sp² carbon systems

Drude-Lorentz model

Example KC₈:



Drude

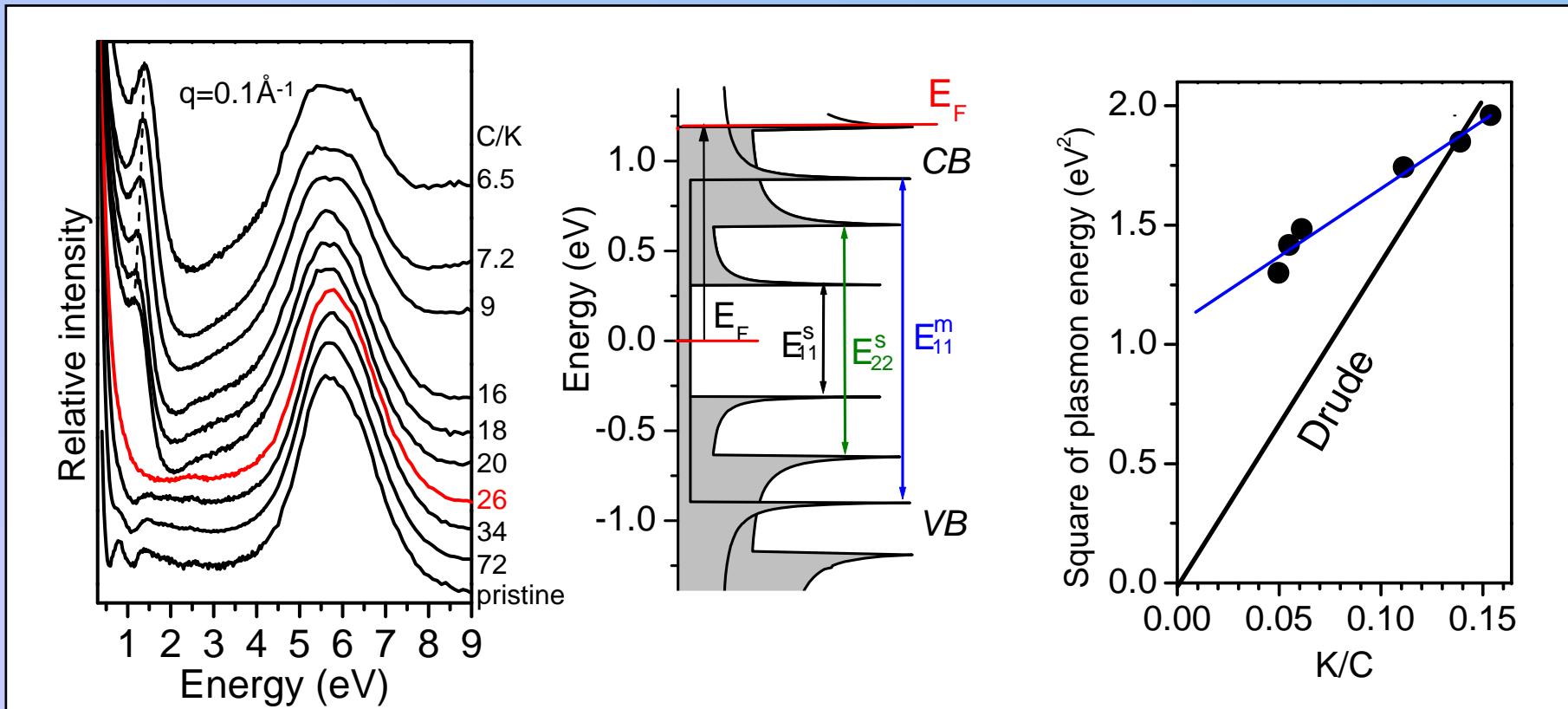
$$\varepsilon(\omega) = 1 - \frac{\omega_p^2}{\omega(i\gamma + \omega)} + \frac{\omega_{p,\pi}^2}{(\omega_{T,\pi}^2 - \omega^2) - i\gamma_\pi\omega} + \frac{\omega_{p,\pi+\sigma}^2}{(\omega_{T,\pi+\sigma}^2 - \omega^2) - i\gamma_{\pi+\sigma}\omega}$$

π plasmon

π+σ plasmon

- Screening effect: → downshift of Drude plasmon
- π, π+σ plasmons = effective dielectric background (ε_∞)
- $\varepsilon_\infty = \varepsilon_\infty^\pi + \varepsilon_\infty^{\pi+\sigma}$

EELS response upon doping (K-doping)



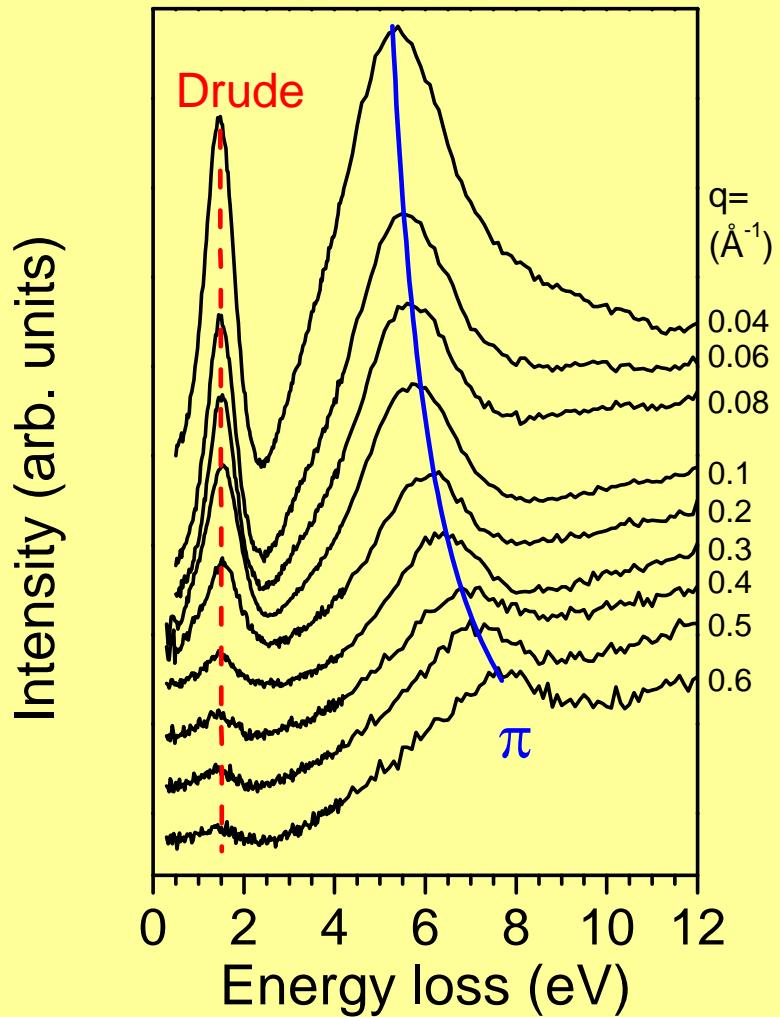
- Interband transitions disappear (up to $C/K=26$)
Free charge carrier (Drude) plasmon above $C/K=26$, Fermi liquid behaviour.
- Plasmon energy (E_p^2) ~ charge density (n):

$$E_p^2 = \frac{\hbar n e^2}{(m^* \epsilon_\infty \epsilon_0)}$$

$m^* \epsilon_\infty \sim$ doping dependent (not constant)

X. Liu et al., PRB 67, 125403 (2003)

Plasmon dispersion in a doped SWCNT bundle



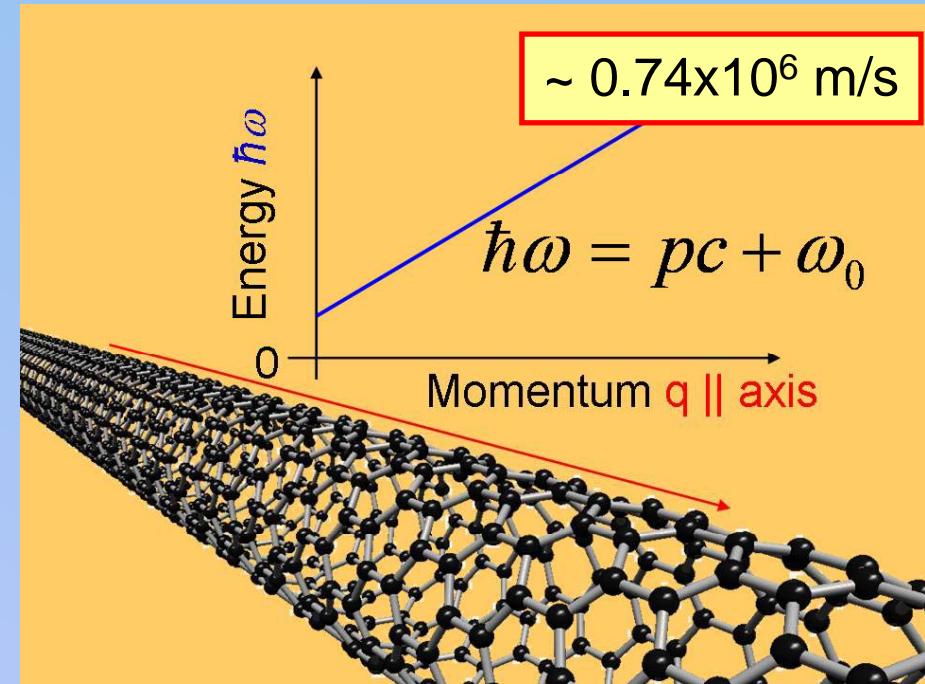
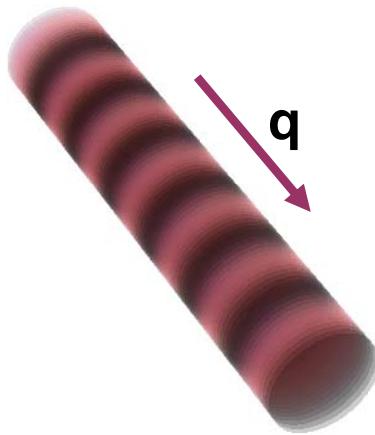
Unchanged π dispersion

Localized Drude plasmon

No 3D dispersion

So far **bulk**
What about isolated?

Summary: Part 1



Isolated VA-SWCNT are truly 1D!

Ist all surface properties, no dielectric screening

a) **Dispersive π plasmon is linear!**

→ mimic K-point excitation in graphene.

b) Charge carrier plasmons in isolated SWCNT are tunable between 0.85 eV and 1.1 eV.
→ similar limits are expected for doped graphene.

Open issue: Why is there no plasmon dispersion?

Where is the acoustic plasmon in 1D,2D? Connection to localized VHS?

Outline

- Introduction/Motivation/Experimental
- Part 1: Angle resolved dielectric response of SWCNT
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 - c) 1D „Drude“ plasmon in intercalated SWCNT?
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Photoemission principles



1905

Albert Einstein

..... 1 Meuro question in
Austria in march (2008)

$$E_{kin} = h\nu - \phi$$



The Nobel Prize in Physics 1921

"for his services to Theoretical Physics, and especially for his discovery of the law of the photoelectric effect"



Albert Einstein

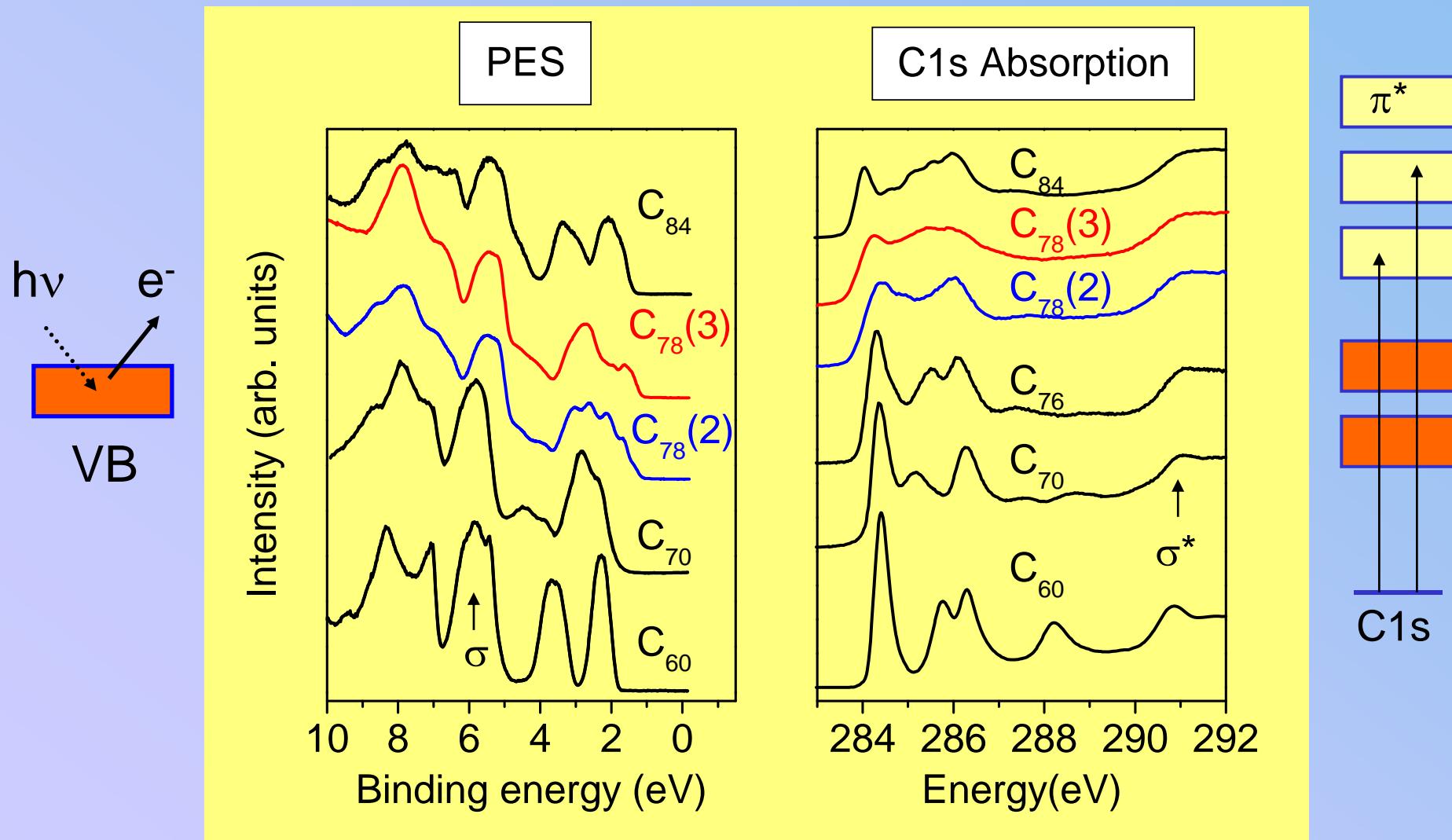
Germany and Switzerland

Kaiser-Wilhelm-Institut (now Max-Planck-Institut) für Physik
Berlin, Germany

b. 1879
(in Ulm, Germany)
d. 1955

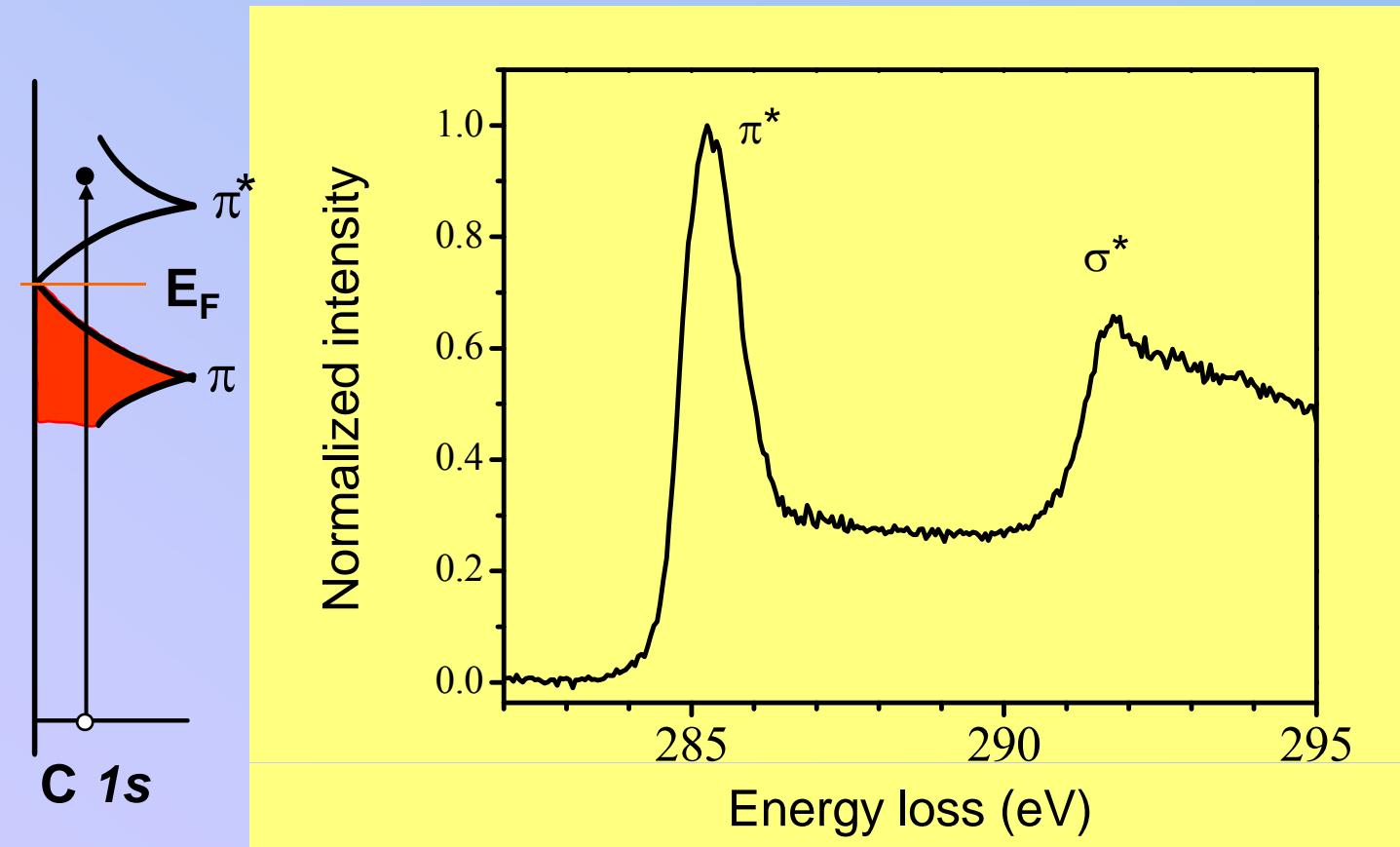
Example: PES, XAS on fullerenes

DOS of molecular orbitals revealed:



Additional gain from EELS, XAS with high energy resolution

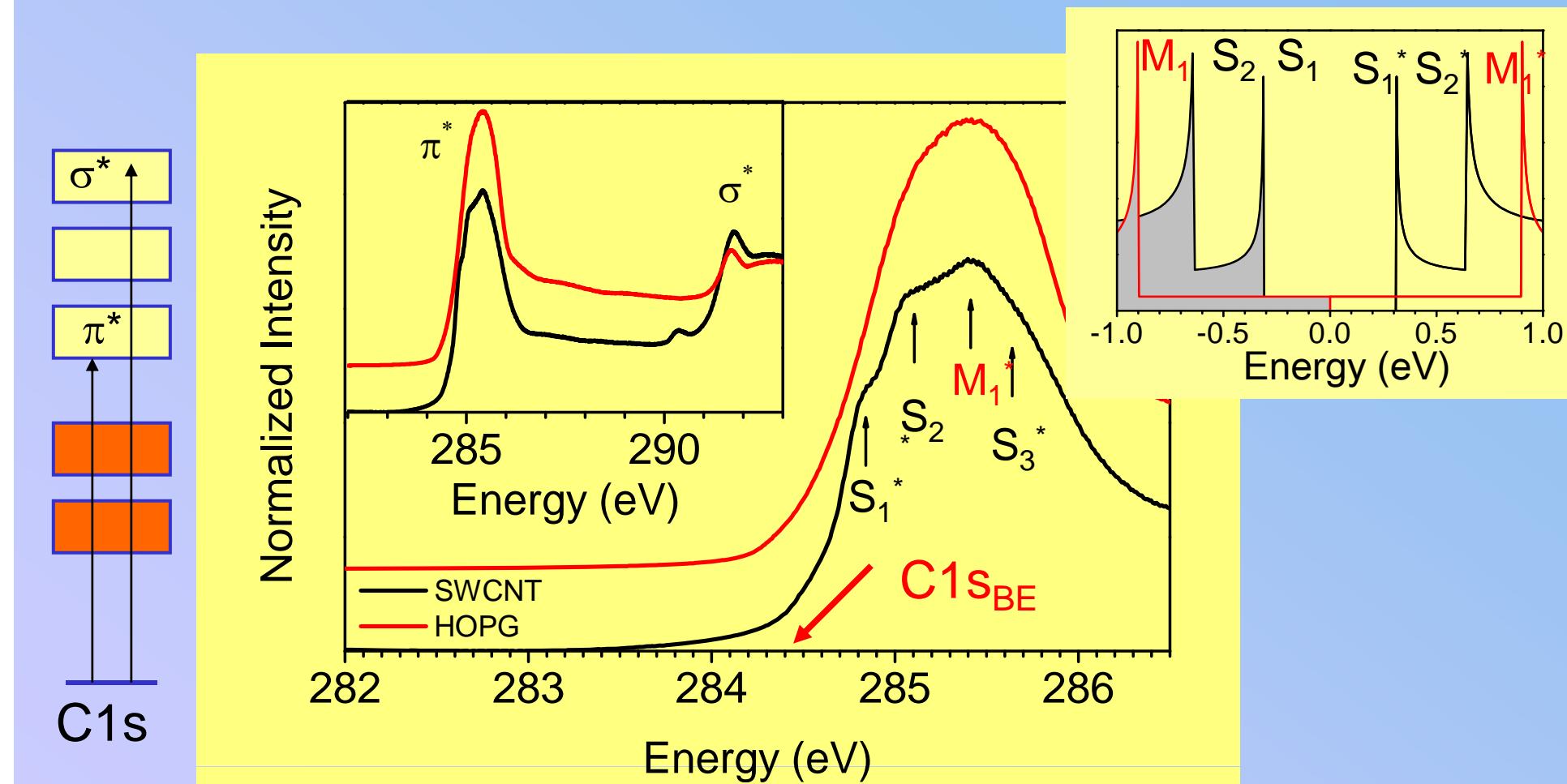
Example: C1s XAS/EELS of graphite:



→ Site selective projected DOS of the conduction band
strongly excitonic by core hole effects:

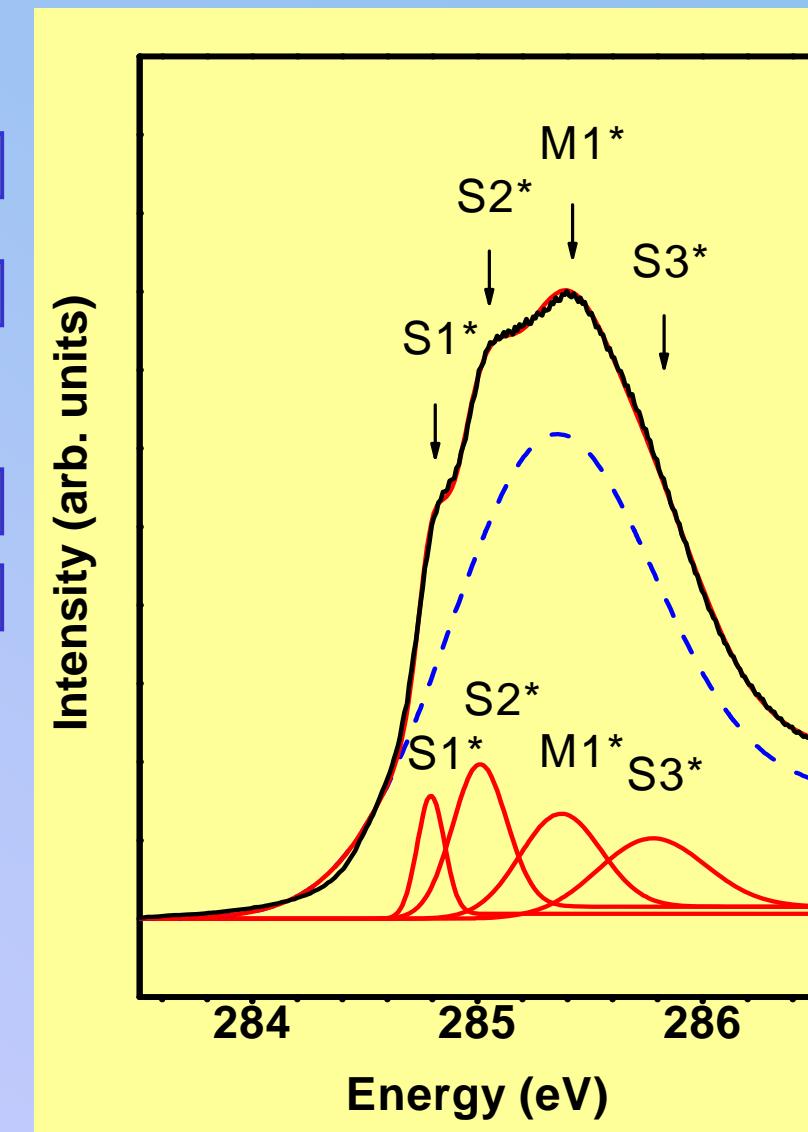
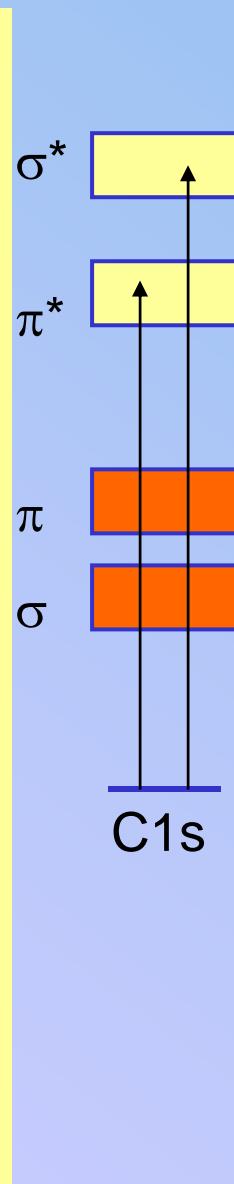
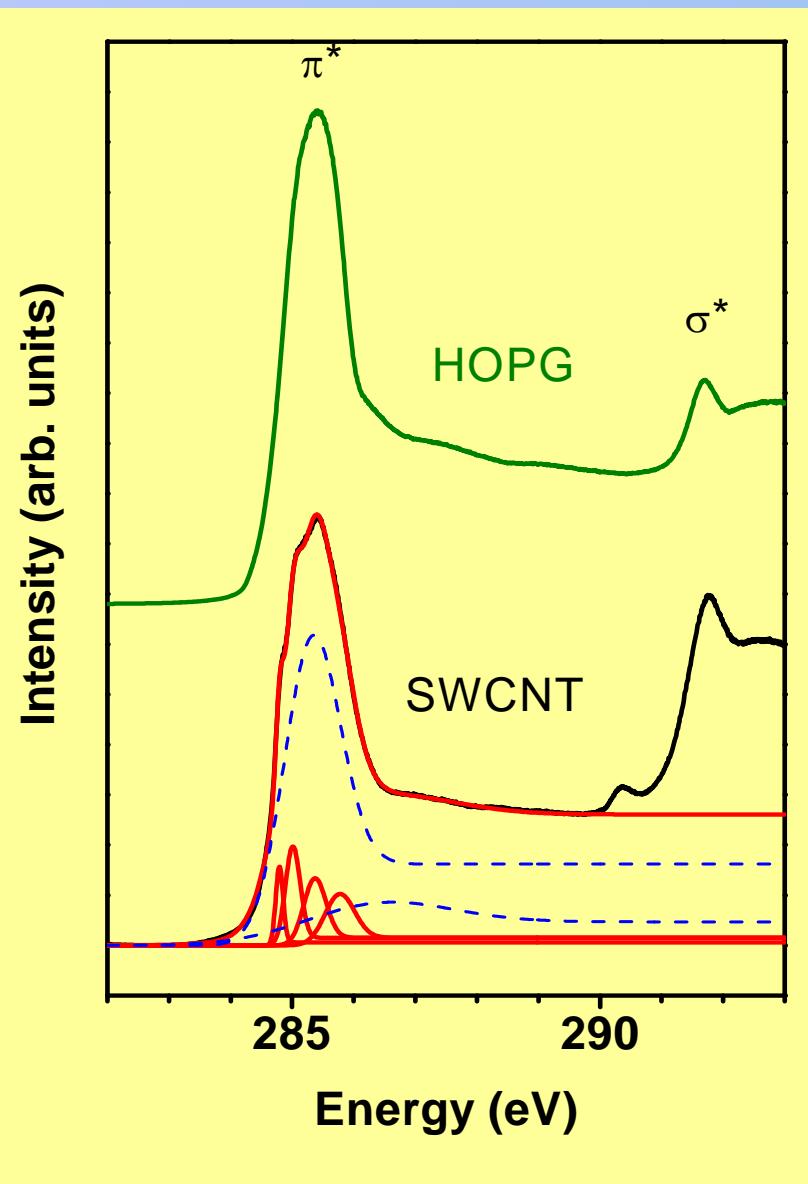
e.g. Bruhwiler PRL 76, 1761 (1996)

XAS: Conduction band of metallicity mixed SWCNT bundles



- ➡ Site selective weighted projected DOS of the conduction band
- ➡ Fine structure in π^* resonance: S_1^*, S_2^*, S_3^* , and M_1^* .
- But: exitonic effects: $\Delta S = 0.25$ eV

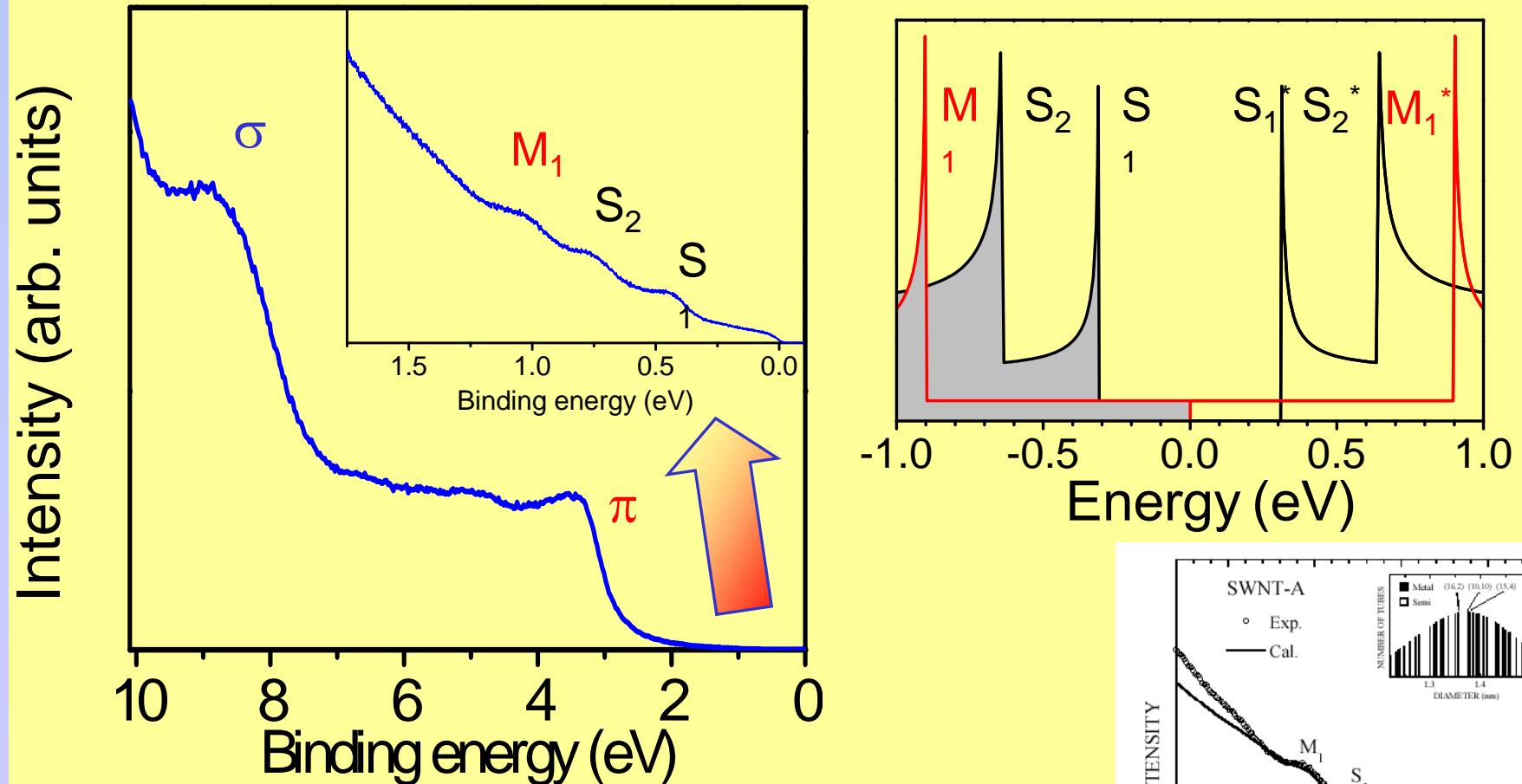
High Resolution XAS of metallicity mixed SWCNT



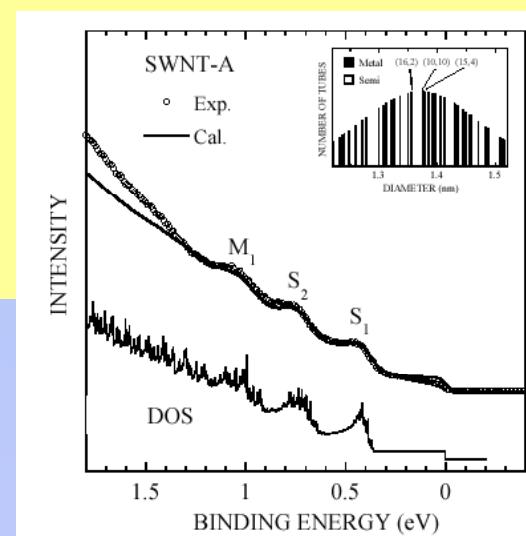
C. Kramberger et al. PRB, 75, 235437 (2007)

Valence band of a metallicity mixed SWCNT bucky paper

High resolution angle integrated photoemission at T=35 K:



vHS of tubes: S_1 , S_2 , M_1
good agreement with TB calculations
e.g. H. Ishii et al, Nature, 426, 540 (2003), H. Rauf et al.



Scaling in 1D and 3D metals: gain from photoemission

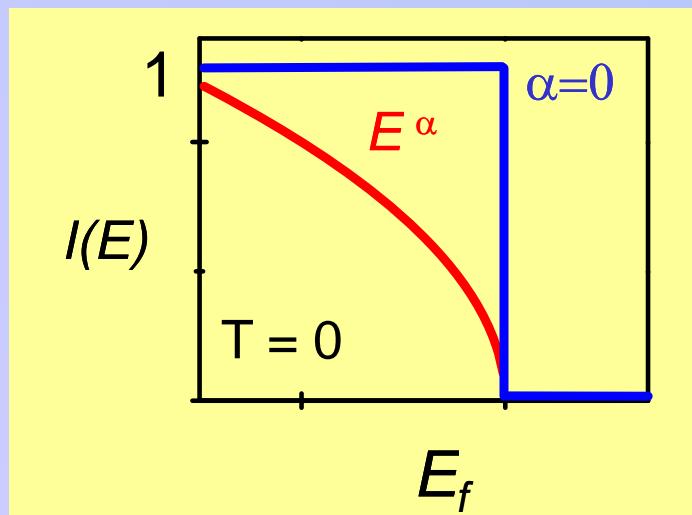
Photoemission intensity: $I(E) \sim n(E) f(E)$

3D metal: Fermi-liquid

- PES shows Fermi edge

$$f(E) = \frac{1}{\exp[(E-E_f)/kT] + 1}$$

$n(E) \sim \text{const}, \alpha=0, g=1$



1D metal: Tomonaga-Luttinger liquid

- 1D paramagnetic metal
- correlated electron state interaction parameter g

$g \ll 1$ repulsive Coulomb interaction

$g > 1$ attractive interaction

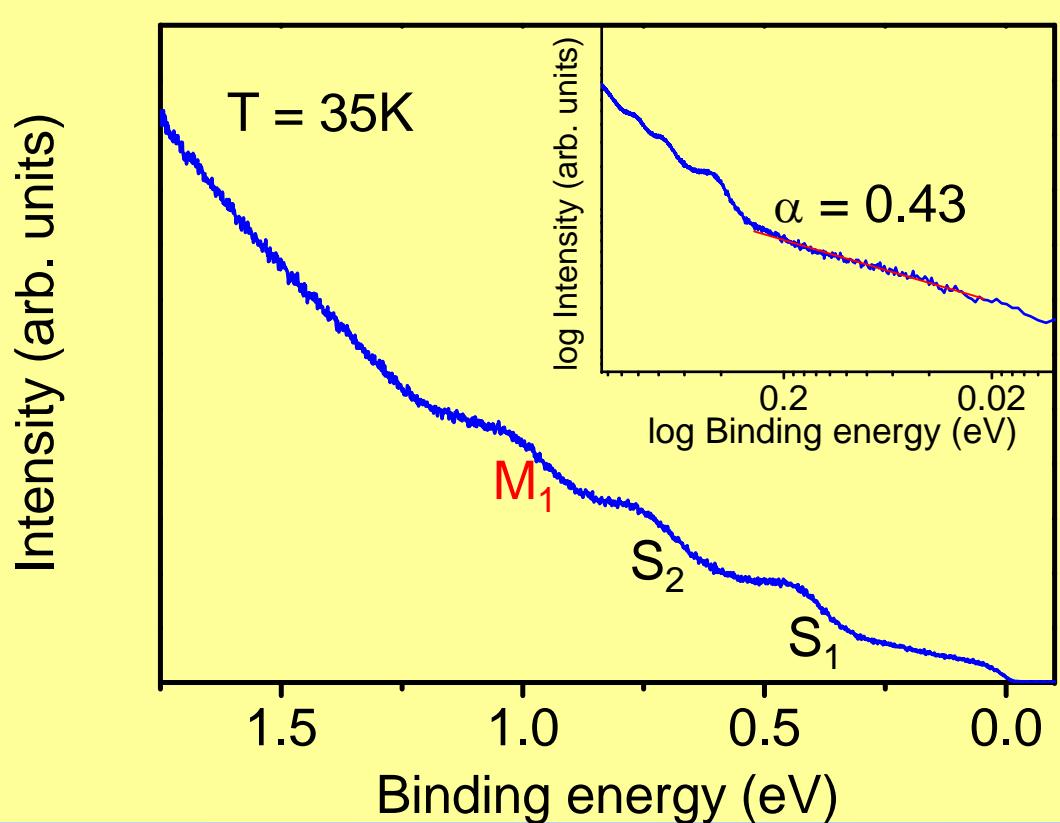
power law dependence of $n(E)$:

$$n(E) \sim E^\alpha$$

For SWCNT:

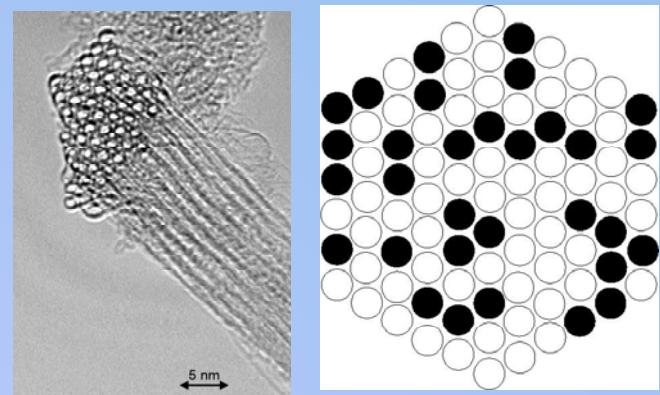
$$\alpha = (g+1/g-2)/8 = 0.3-0.5$$

Effect of dimensionality: Photoemission



→ $\alpha = 0.43, g=0.18$

Metallic SWCNT
in a bundle of SWCNT
have the renormalization
of a Tomonaga-
Luttinger liquid

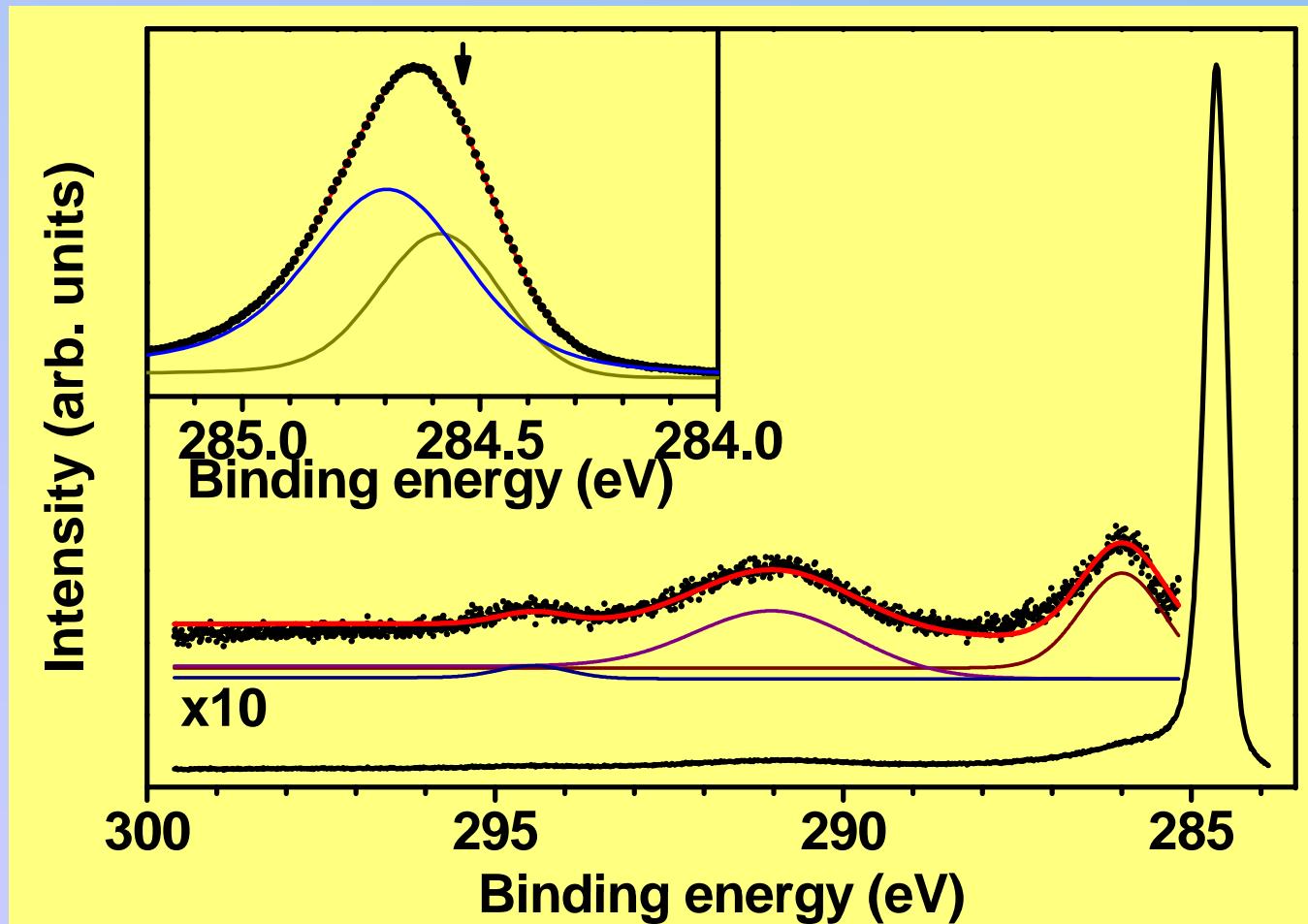


○ semiconducting
● metallic

good agreement with theoretical predictions and results from transport measurements

H. Ishii et al, Nature, 426, 540 (2003); H. Rauf et al, PRL 93, 096805 (2004)

High resolution XPS of metallicity mixed SWCNT



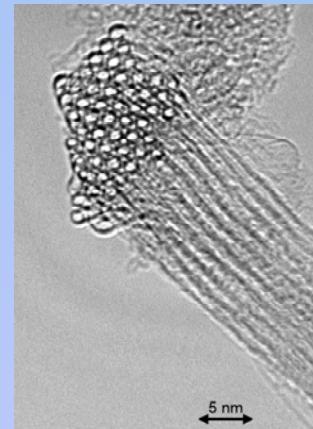
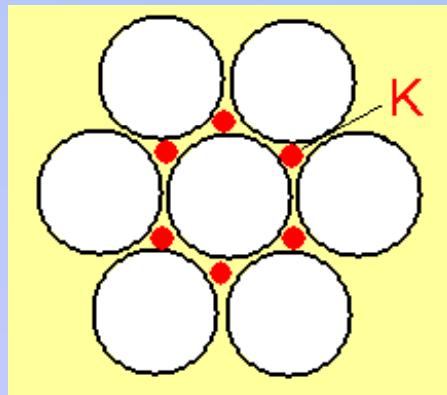
- C1s for SWCNT has two components at different binding energy
- One Doniac-Sunjic lineshape (metallic SWCNT)
One Voigthian (semicond. SWCNT)

Outline

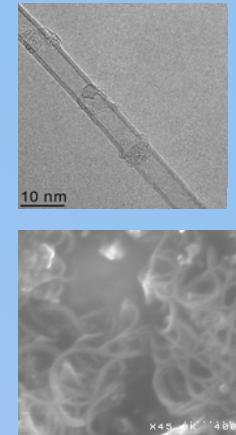
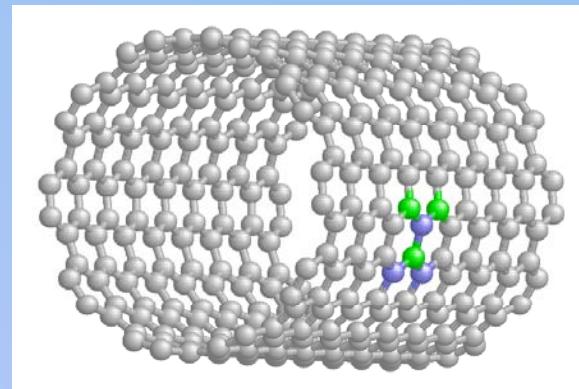
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Functionalization: doping and filling of SWCNT

● Intercalation:



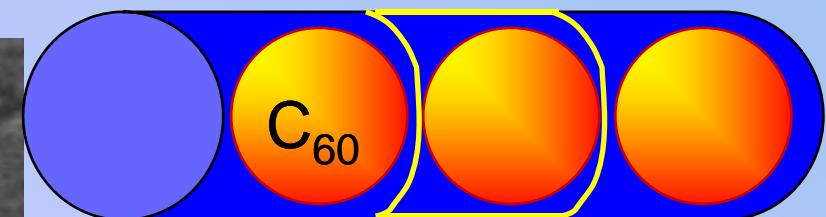
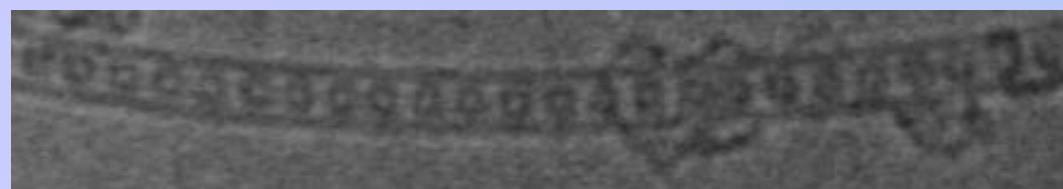
Substitution:



e.g. K_x -SWCNT

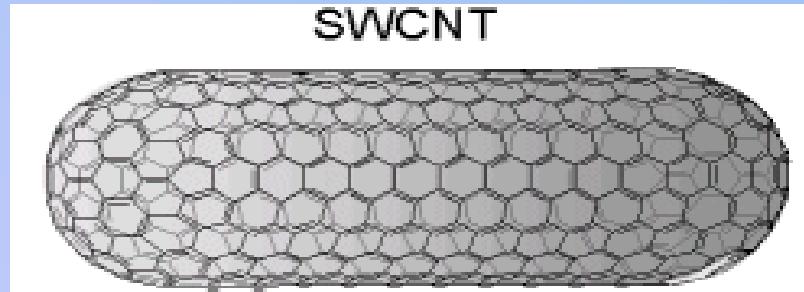
BC_3 -SWNT, MWBNT, SiCNT

● Filling: C_{60} Peapods, DWCNT....

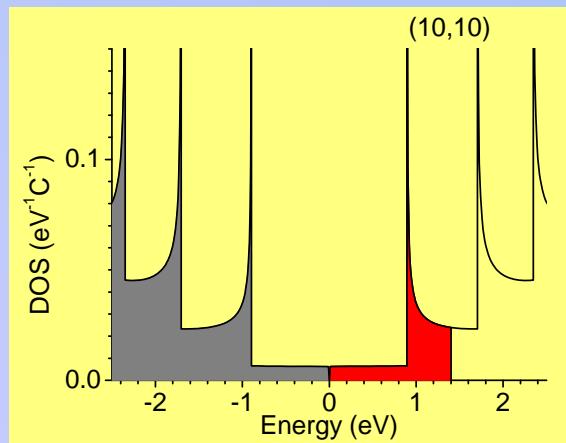
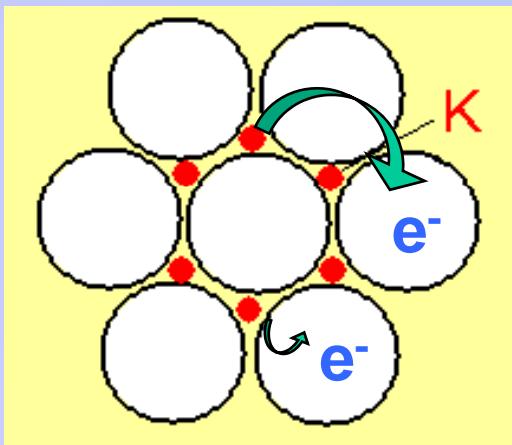
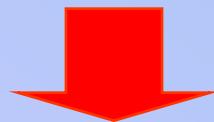


→ Optical and high energy spectroscopy are ideal tools!

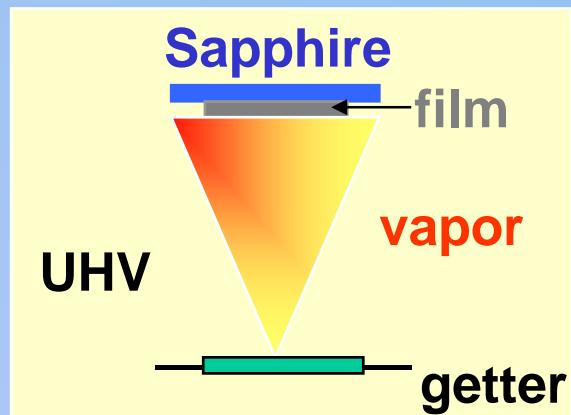
Intercalation



doping

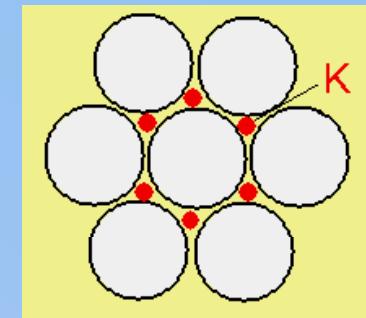
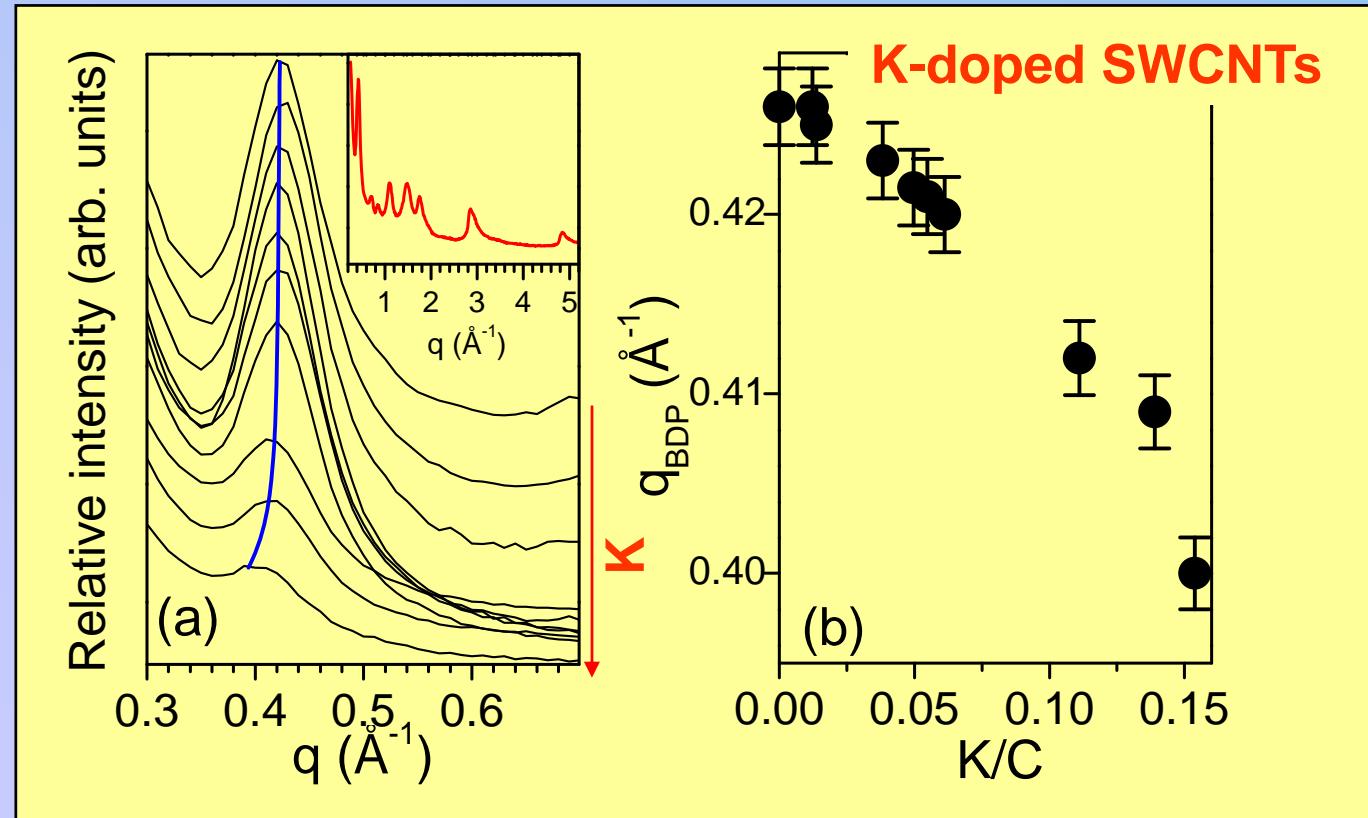


Alkali-metal intercalation

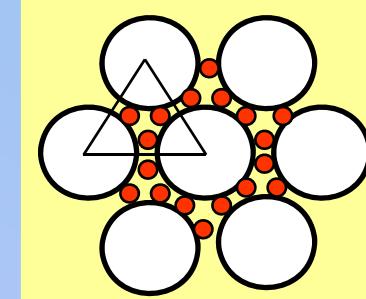


in-situ doping:
UHV evaporation
(5×10^{-10} mbar)
Na, K, Rb, Cs and
Ba SAES getters

Structural analysis of alkali-metal doped SWCNTs



no expansion



expansion

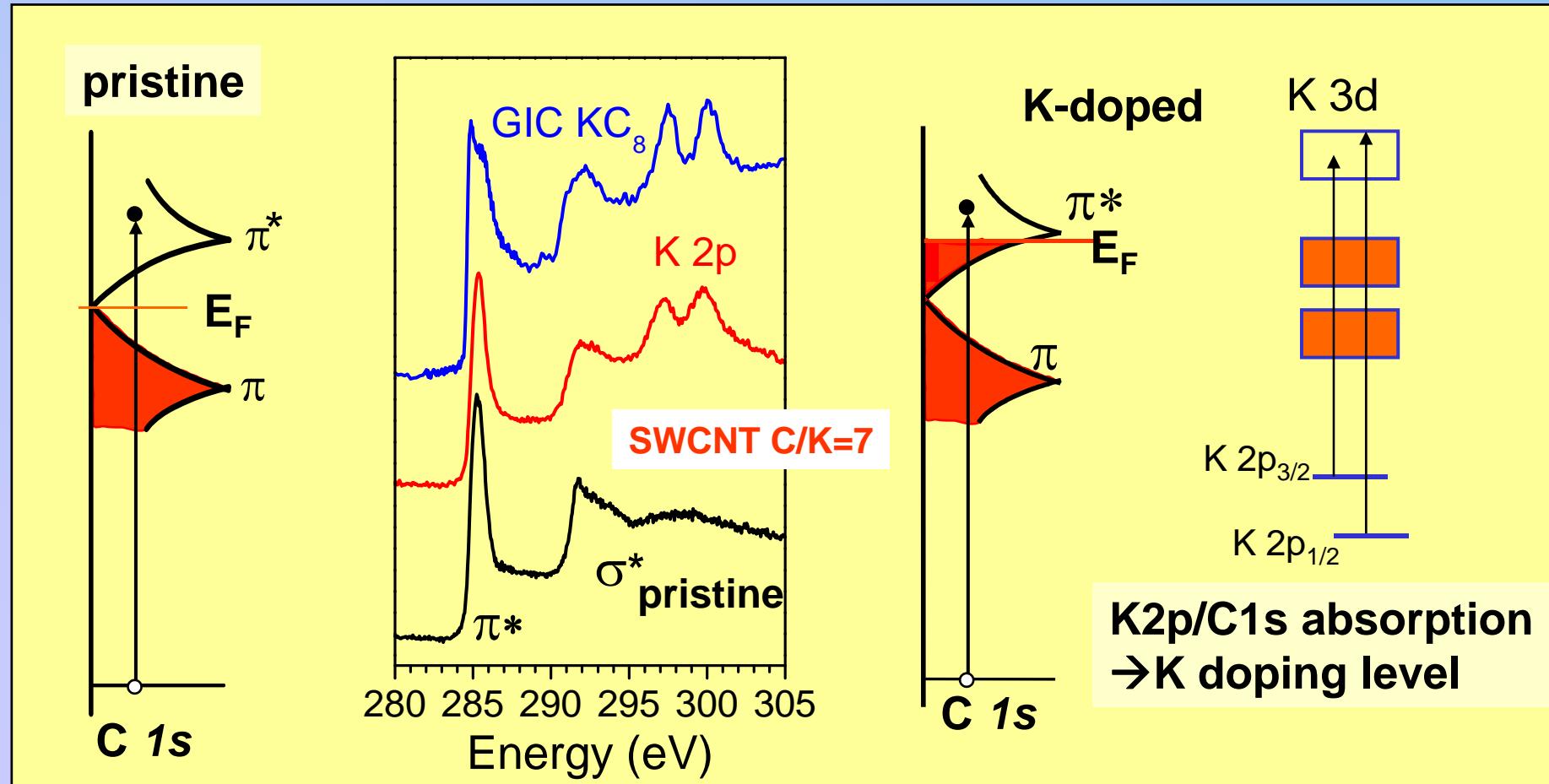
$d=1.37\text{nm}$

	(\AA)	Na	K	Rb	Cs
Ionic radius:	0.95, 1.33, 1.48, 1.69				
Expansion:	0	1.1	1.7	2.7	

alkali intercalates between the SWCNT
lattice expansion depends on doping level,
diameter & size of dopant. No staging

T. Pichler et al. SSC 109, 721 (1999);
X. Liu et al. PRB 67, 125403 (2003)
X. Liu et al. PRB 70, 245435 (2004)

Intercalation: C,K edge excitations (alkali-metal doping)



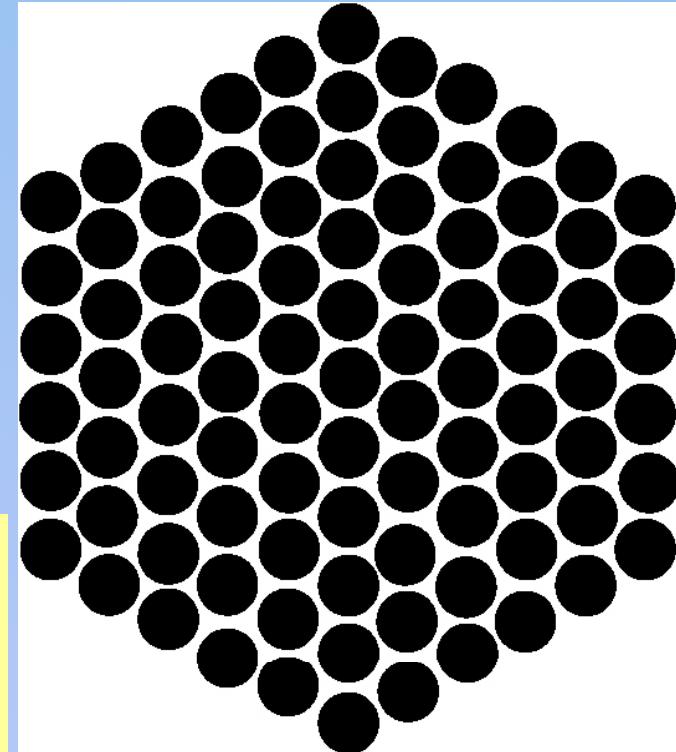
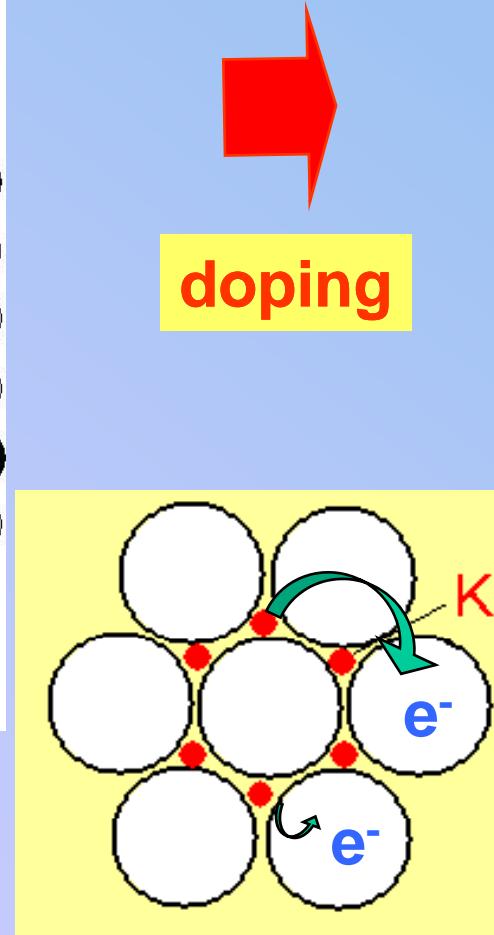
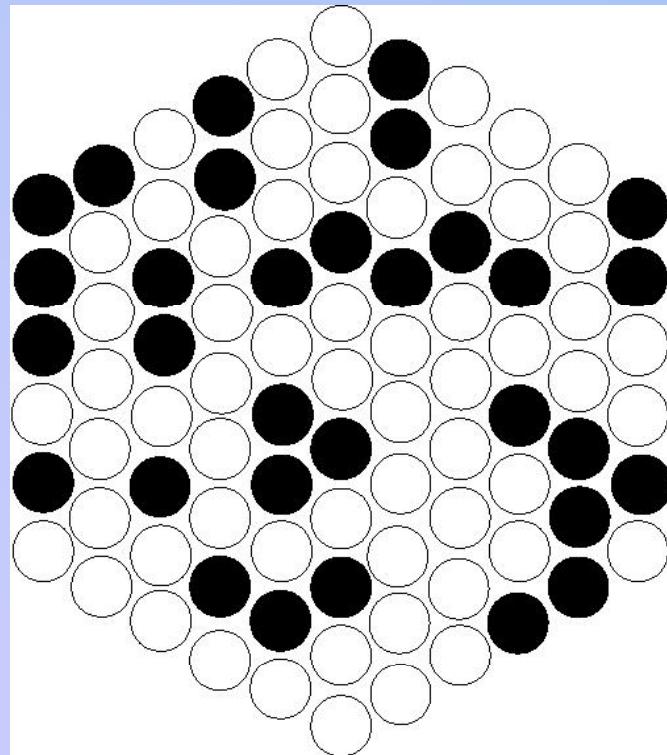
Ratio C1s/K2p edges:

Maximum intercalation ratio C/K=7

Ionic charge transfer, π^* resonance decreases by 15%

T. Pichler et al. SSC 109, 721 (1999), X. Liu et al. PRB 67, 125403 (2003)

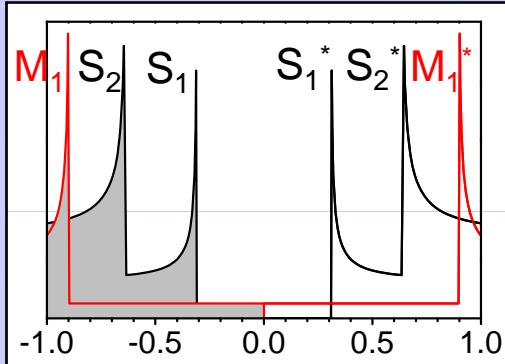
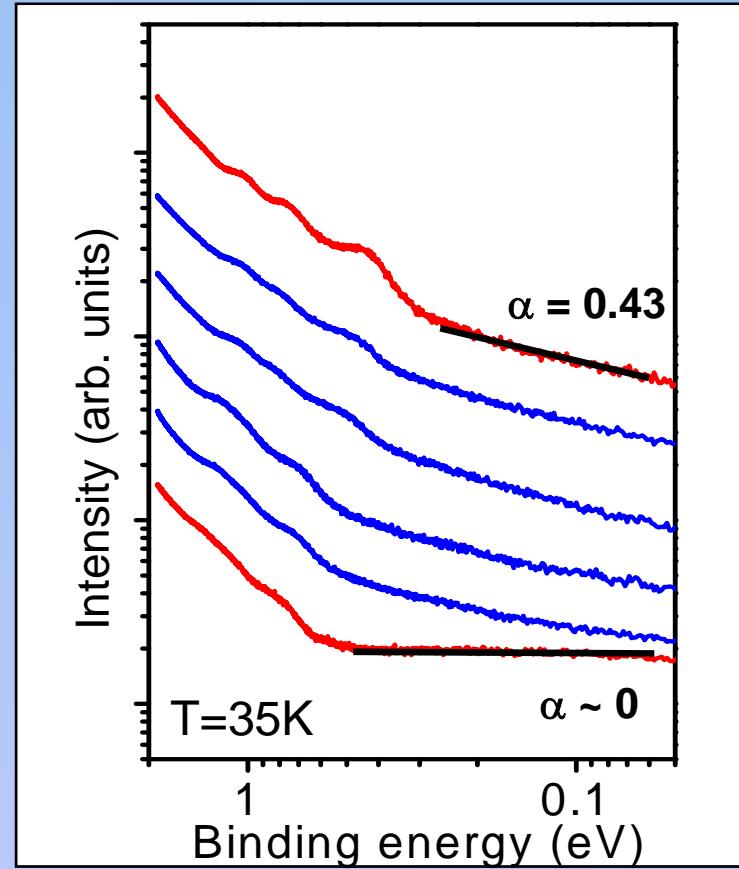
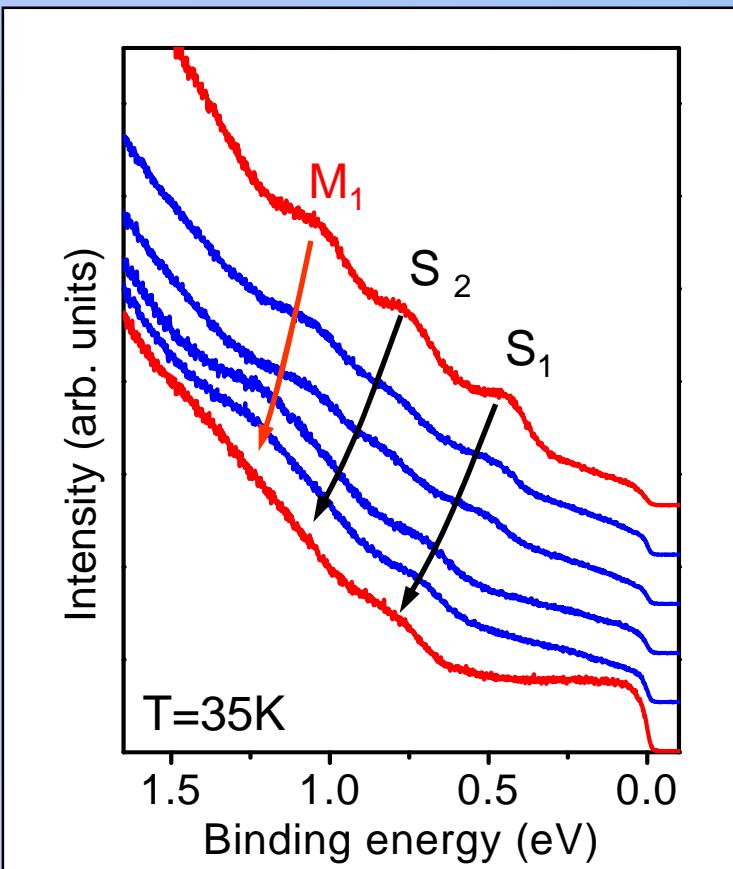
Intercalation: what happens at low doping?



○ semiconducting
● metallic

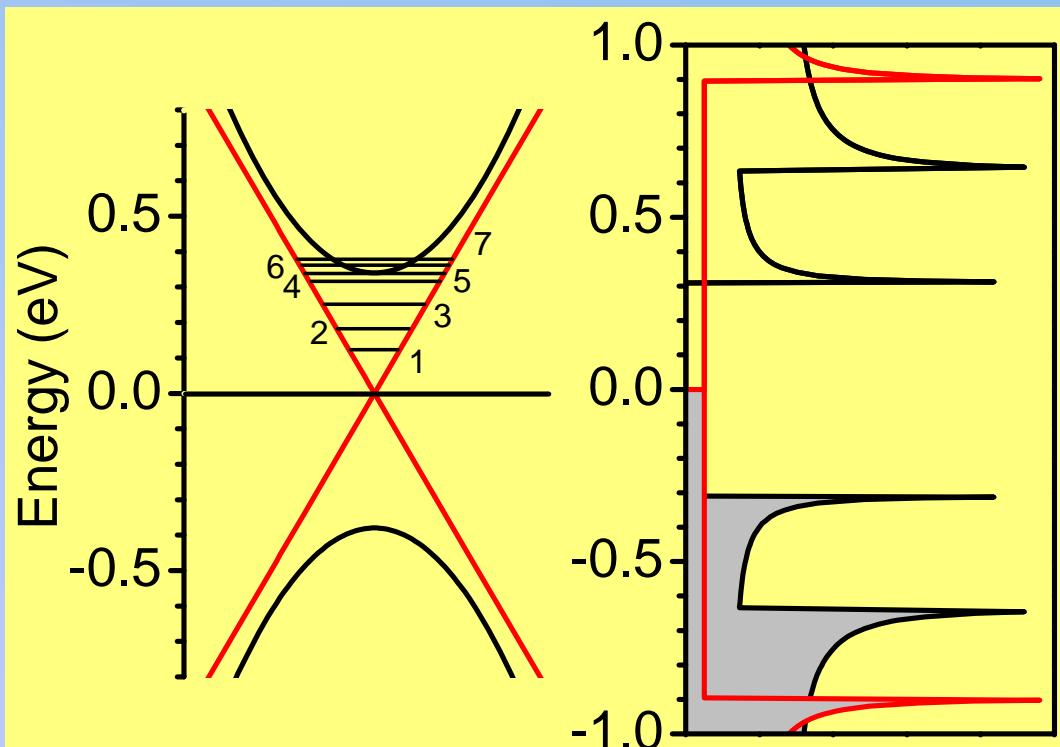
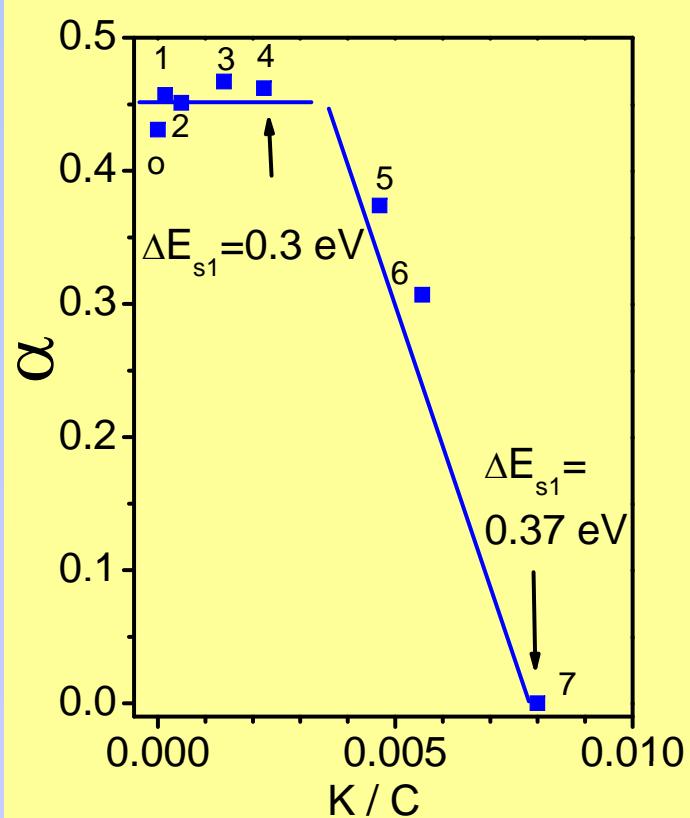
Is a bundle of metallic tubes a
normal Fermi liquid or a Luttinger liquid?

Rigid band shift



- S_1 , S_2 and M_1 shift to higher binding energies
- power law behavior at first stable
- abrupt change at K:C = 1:125, $\Delta E(S_1) = 0.37\text{eV}$

Doping dependence of the LL parameter α



Power law scaling vs. intercalation:

$\alpha = \text{const. up to } K/C = 1/500, \Delta E_{s1} = 0.25 \text{ eV}$

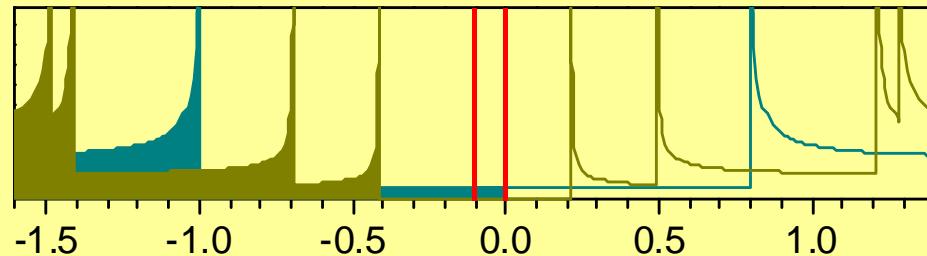
$\alpha = 0, g = 1 \text{ above } K/C = 1/150, \Delta E_{s1} > 0.35 \text{ eV}$

→ a bundle of only metallic tubes is a normal Fermi liquid!

Electronic structure of doped SWCNT bundles: 4 regions

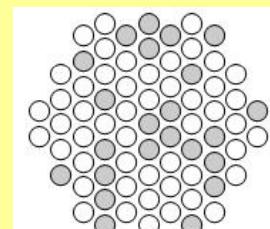
pristine

s/m



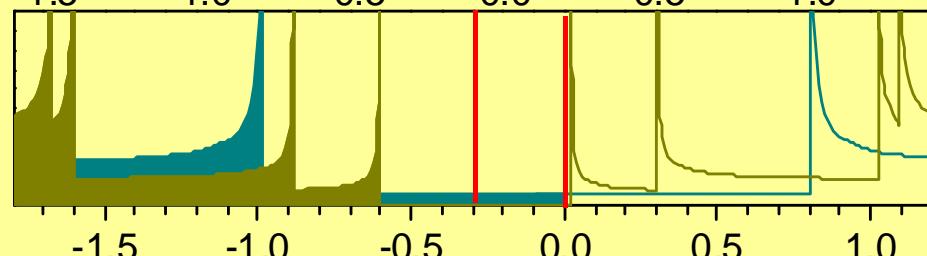
K/C

0 ‰

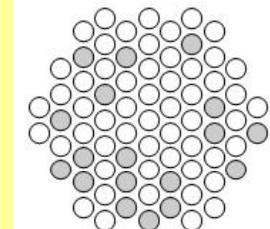


low

s/m

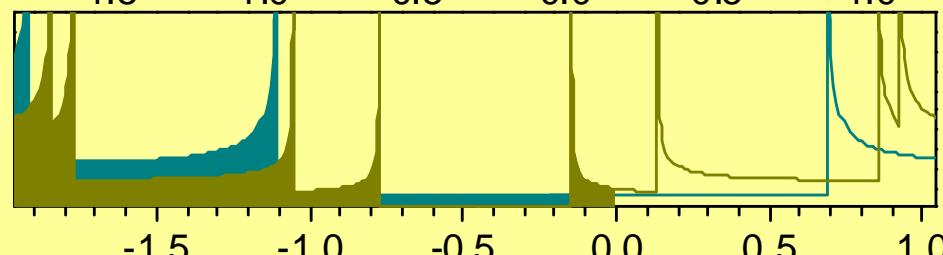


< 10 ‰

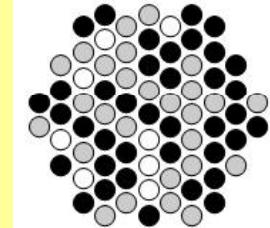


intermediate

M/m

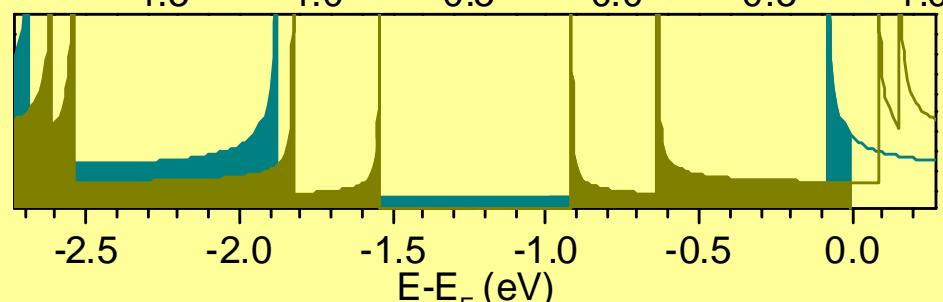


< 50 ‰

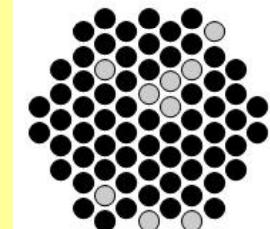


high

M/M



> 50 ‰



Filling of the inner space of SWCNTs with molecules

● Modified electronic properties?

What happens upon the formation of the peapods?

Charge transfer, chemical bonding

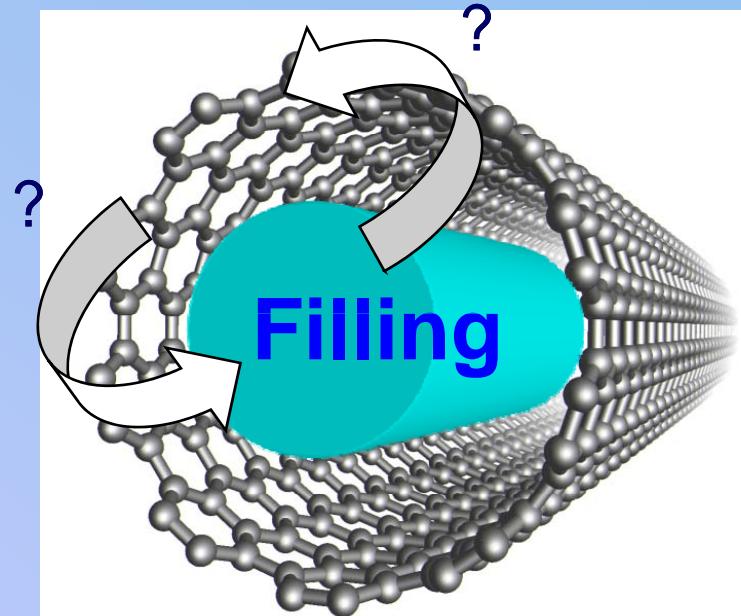
or orbital hybridisation?

Conversion to double-wall carbon nanotubes (DWCNT)?

● Potential applications

Carrier tuning of SWCNT, protection for molecules and even drugs, stabilizing reactive compounds and nanocrystals, **chemical reaction inside a 1D confined nano reaction tube**, 1D spinchain for spin electronics, ...

Filling@SWCNT
so-called “peapod”



Examples

● Fullerenes@SWCNT

C_{60} @SWCNT :

1D C_{60} chain inside SWCNT (TEM,Raman)

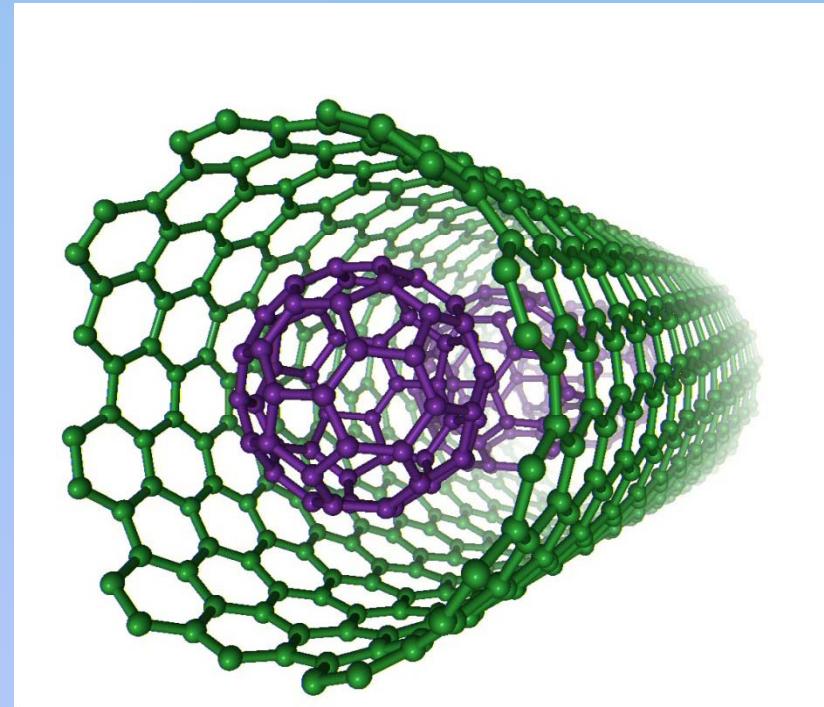
Gd@ C_{82} @SWCNT :

Band gap modulation (STM,STS)

Dy_3N @ C_{80} @SWCNT :

Charge transfer (PES)

Conversion to DWCNT(TEM, Raman)!

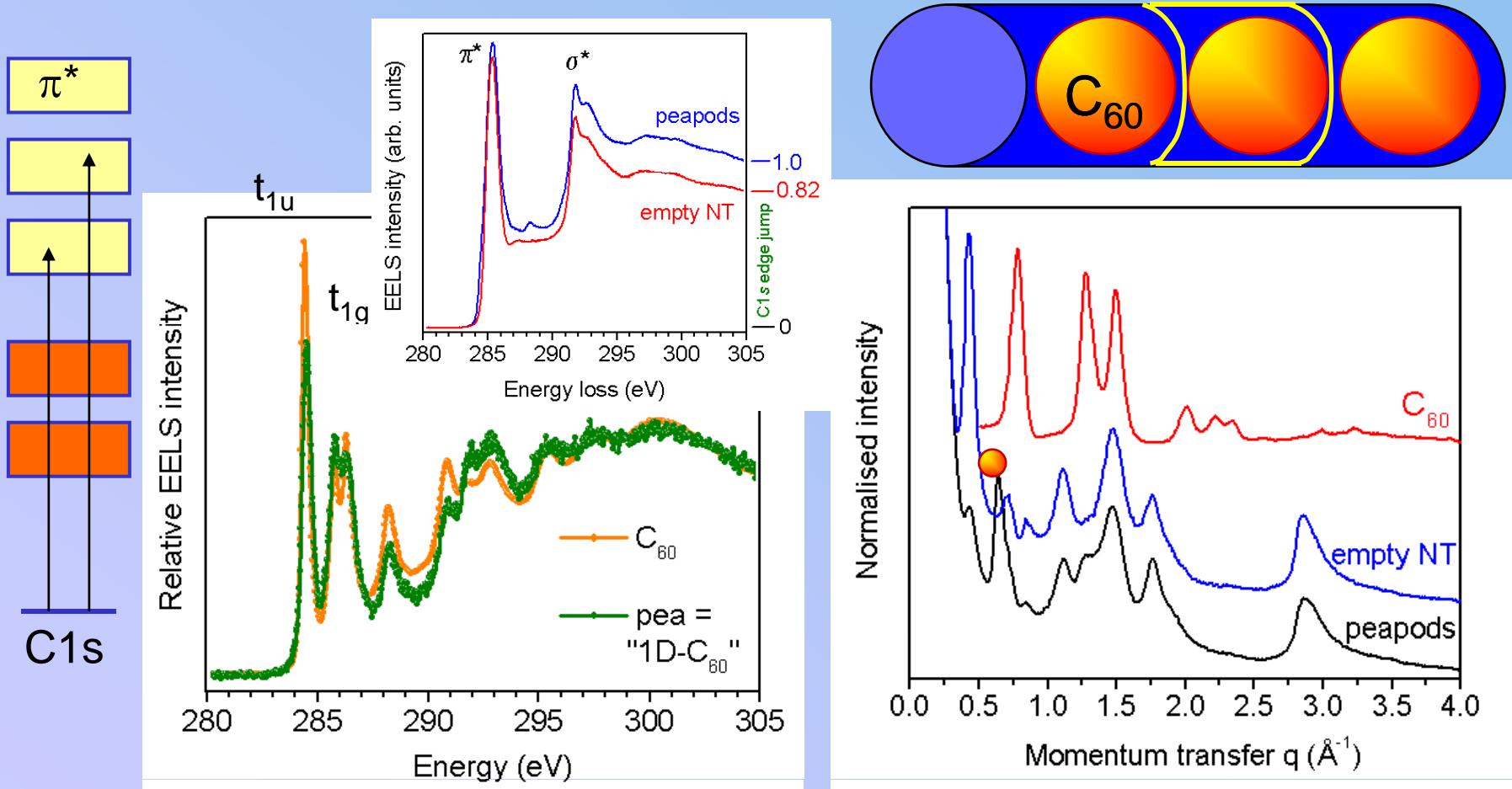


C_{60} @SWCNT

● Organic molecules @SWCNT

TCNQ, TDAE, Anthracene@SWCNT,
Metallocenes p- and n- type doping

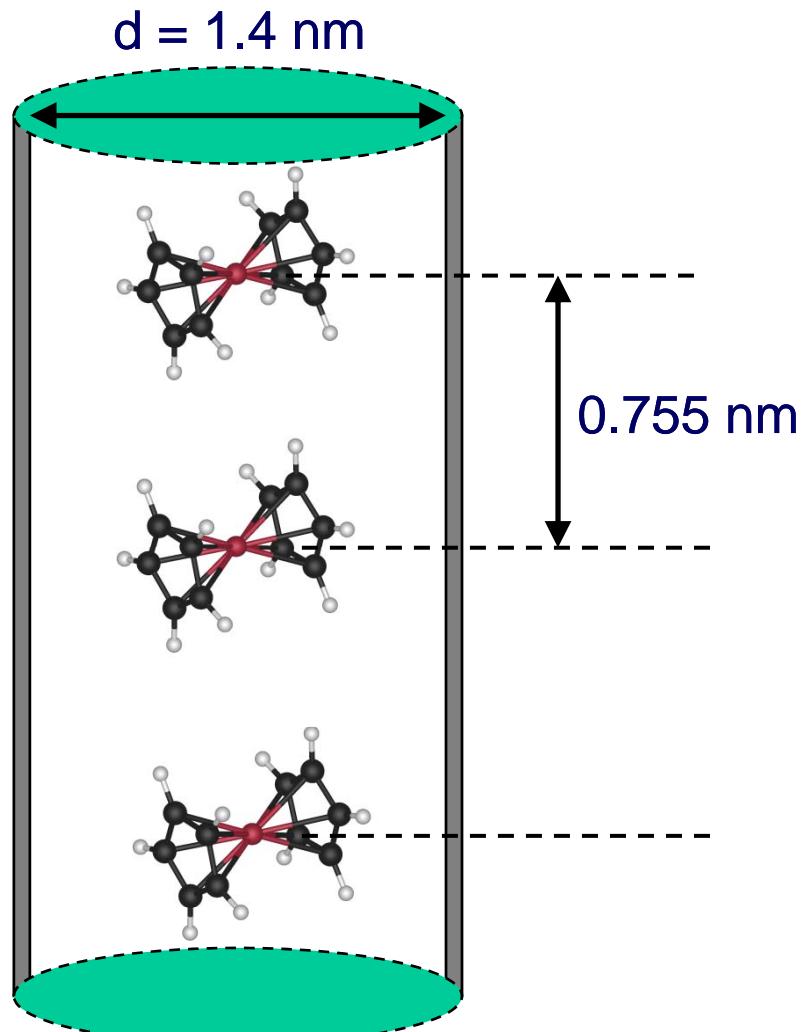
Bulk filling ratio, structure in peapods



Maximum filling ratio with C₆₀: 100%
no charge transfer to C₆₀

C₆₀ - C₆₀ distance 9.7 Å
➡ no C₆₀ polymer in SWCNT

Metallocene filling: Filling factor



Assuming a 1D ferrocene chain inside SWCNT ($d = 1.4 \text{ nm}$) with adjacent ferrocene distance $\sim 0.755 \text{ nm}$

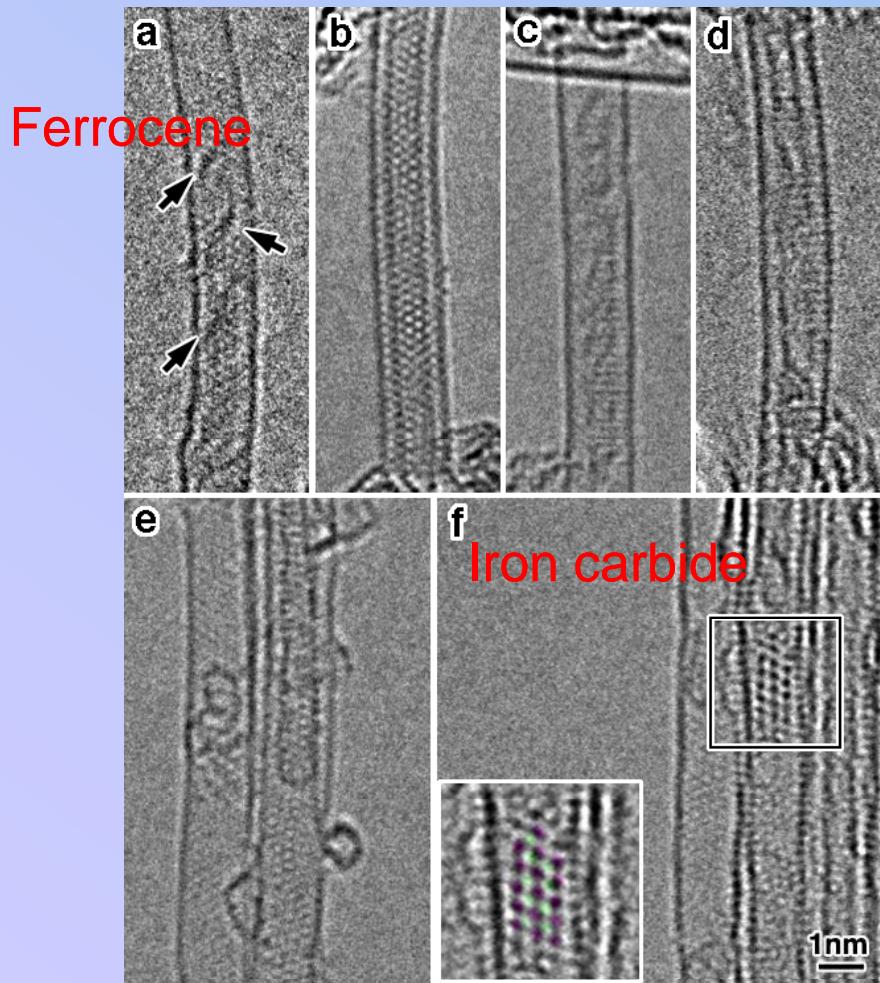
Fe/C ratio ~ 0.003



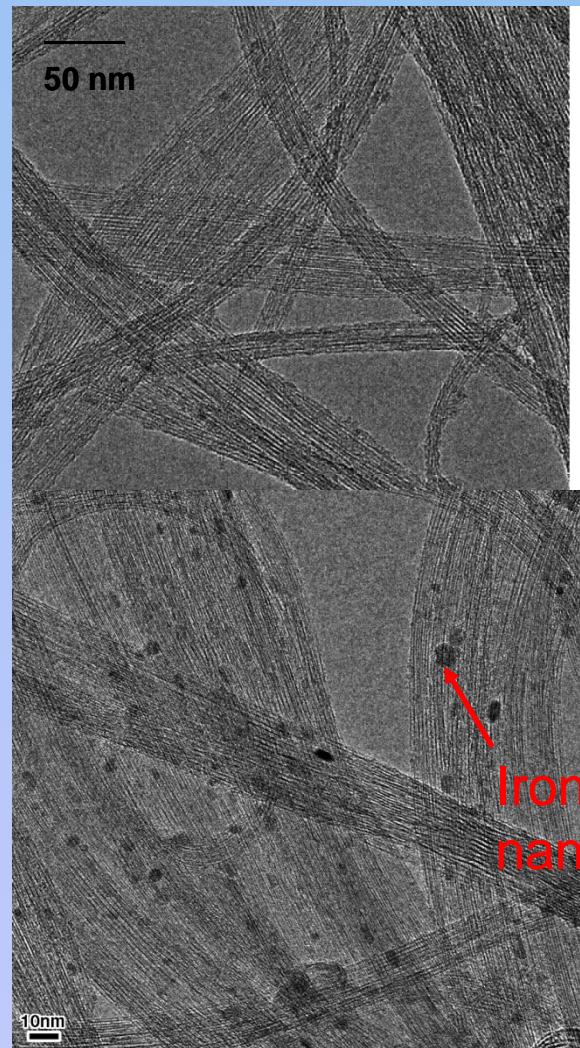
Filling factor $\sim 35 \%$

Nanochemistry in 1D: Gain from TEM

High resolution

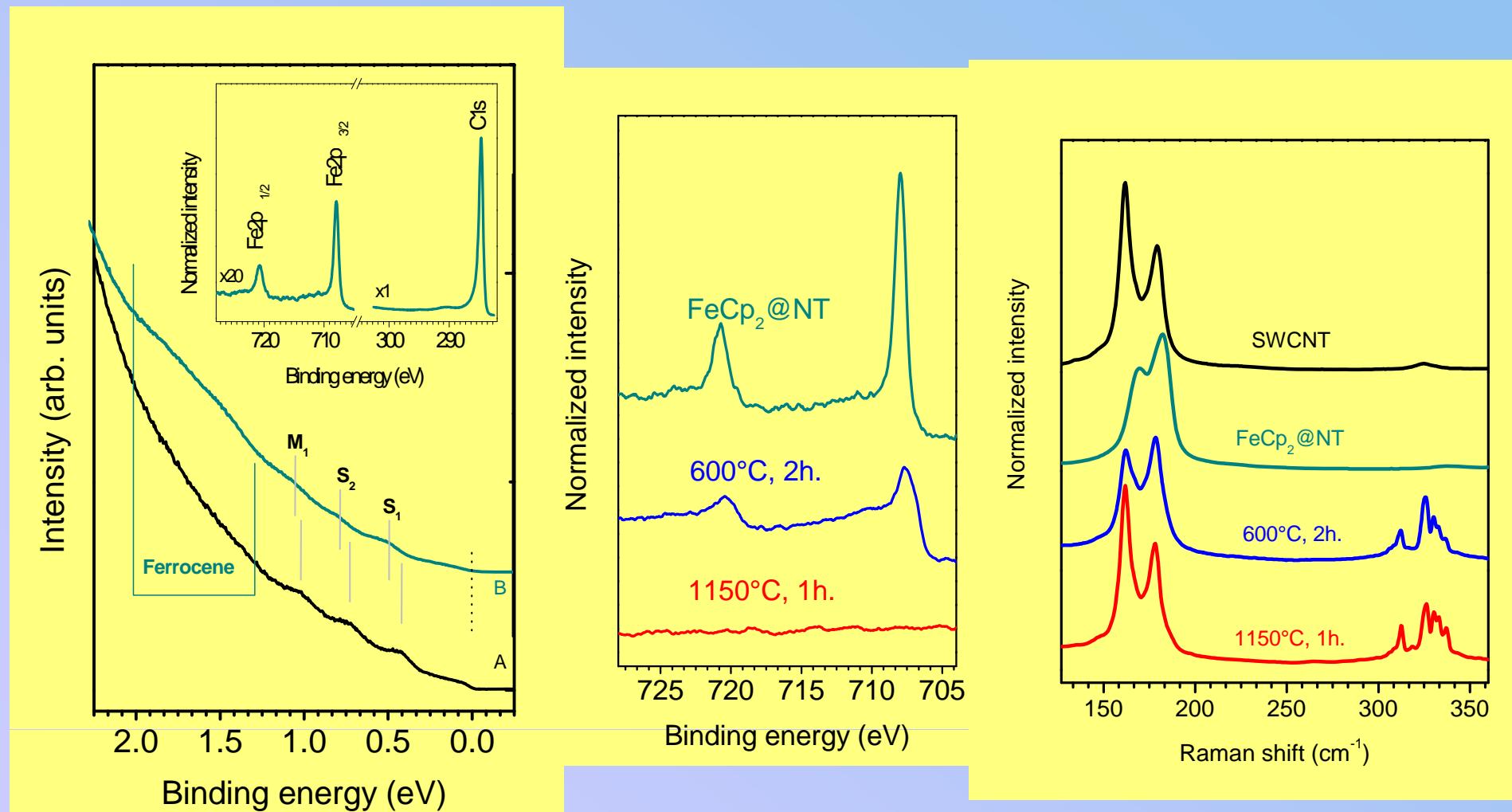


Overview



H. Shiozawa et al. Adv. Mat. (2008)

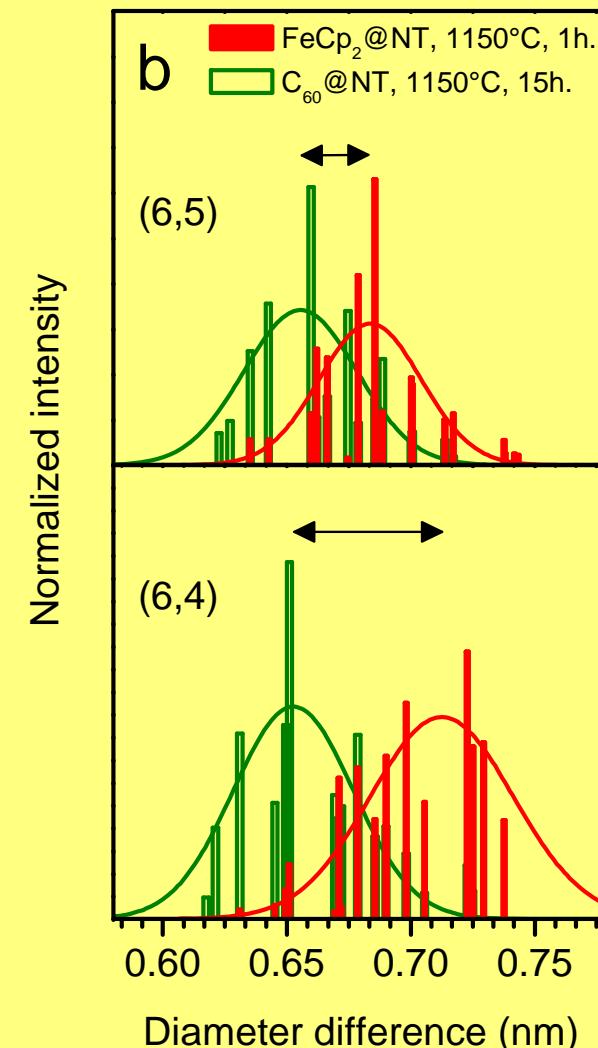
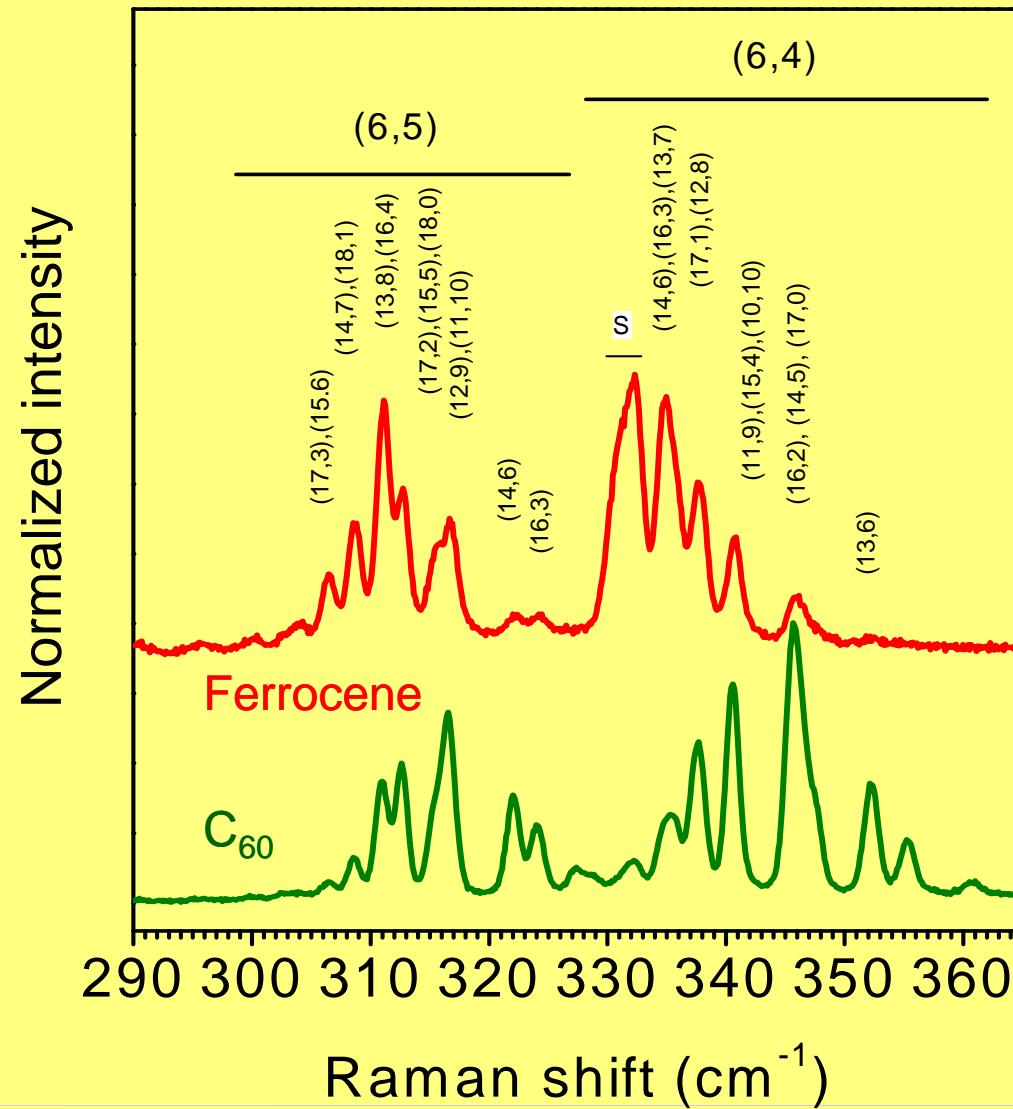
Ferrocene filled SWCNT and growth of DWCNT



- Charge transfer $\Delta E_F = 0.1$ eV, $\sim 0.2 e^-$ per ferrocene
- Power law behavior $\alpha = 0.55$
- growth of inner tubes at $\sim 600^\circ C$

H. Shiozawa et al.
Adv. Mat. (2008); PRB (2008)

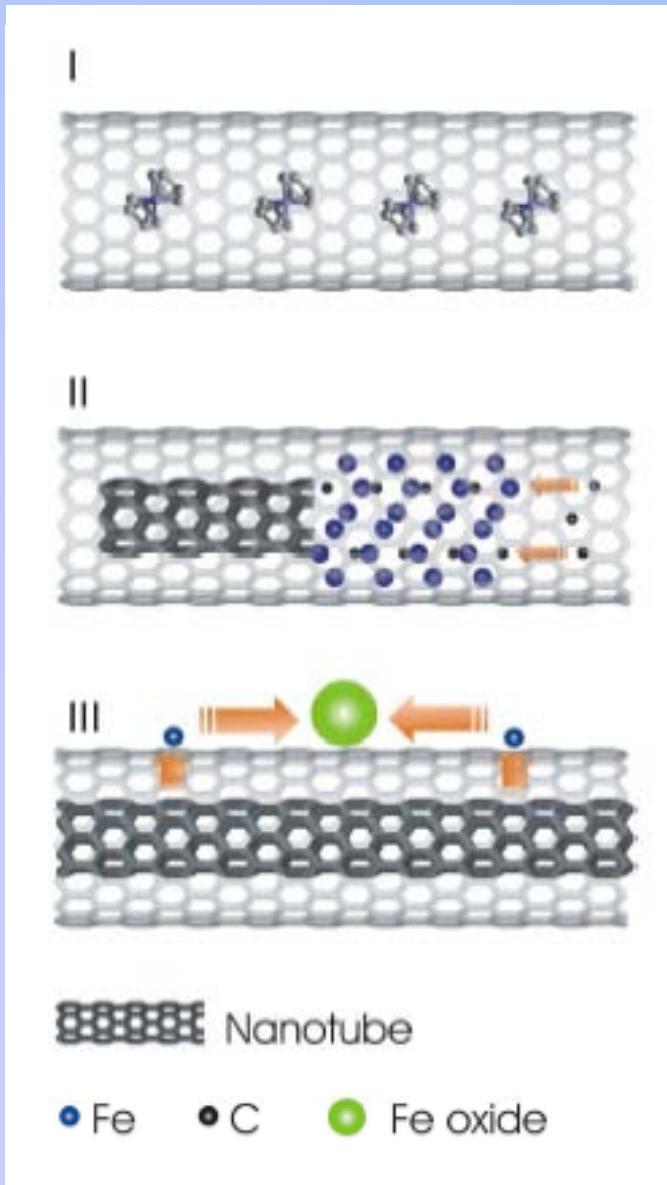
DWCNT/Inner tubes: Gain from high resolution Raman



Catalytic growth of inner DWCNT
differs from non-catalytic!

H. Shiozawa et al. Adv. Mat. (2008)

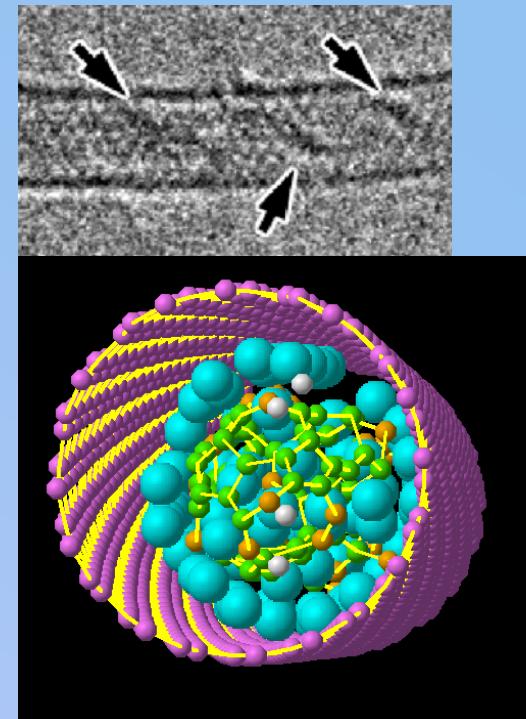
Temperature induced chemical reaction in a SWCNT nanoreactor



1. Ferrocene filling

2. Catalytic growth from
Iron carbide particles

3. Iron release via defects
Iron oxide nanocrystals

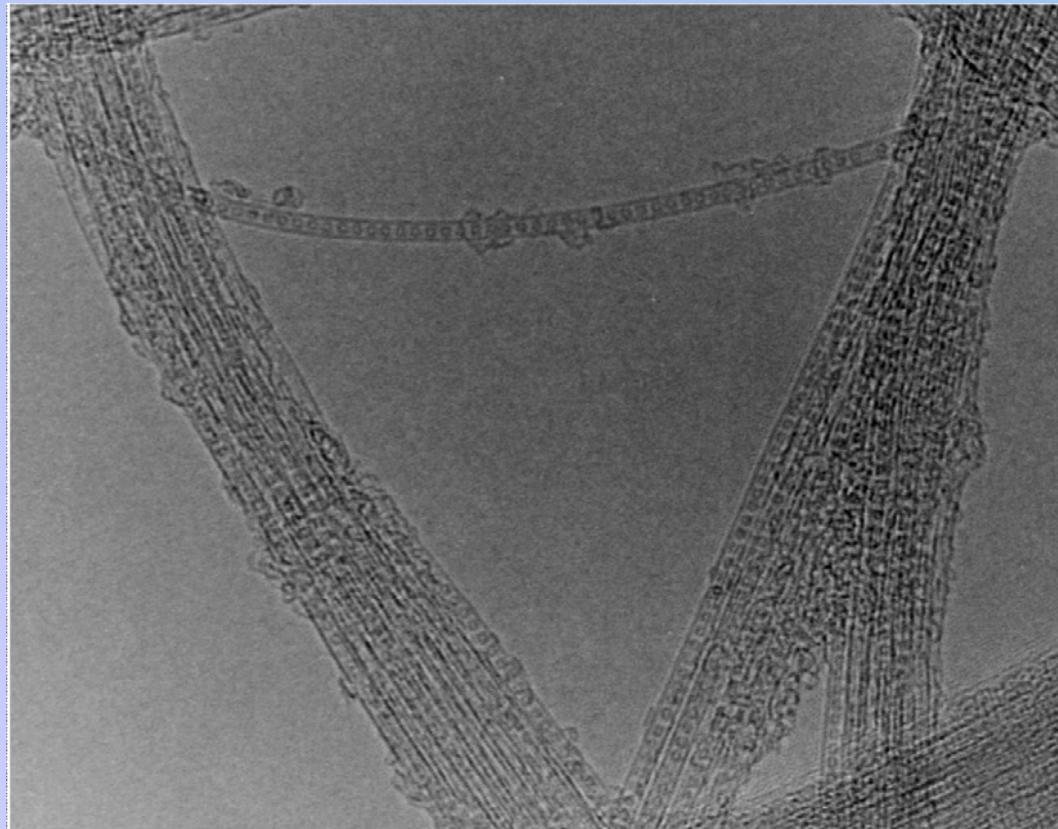


MD simulation
from S. Maruyama

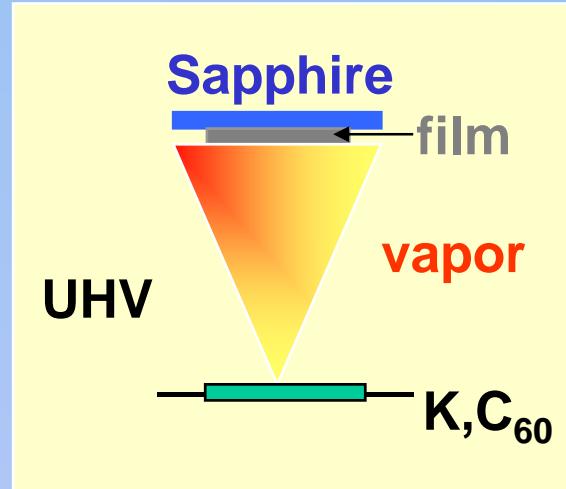
H. Shiozawa et al.
Adv. Mat. (2008); PRB (2008)

Combinational doping: Intercalation of peapods

Peapods: e.g. C₆₀ filled SWCNT

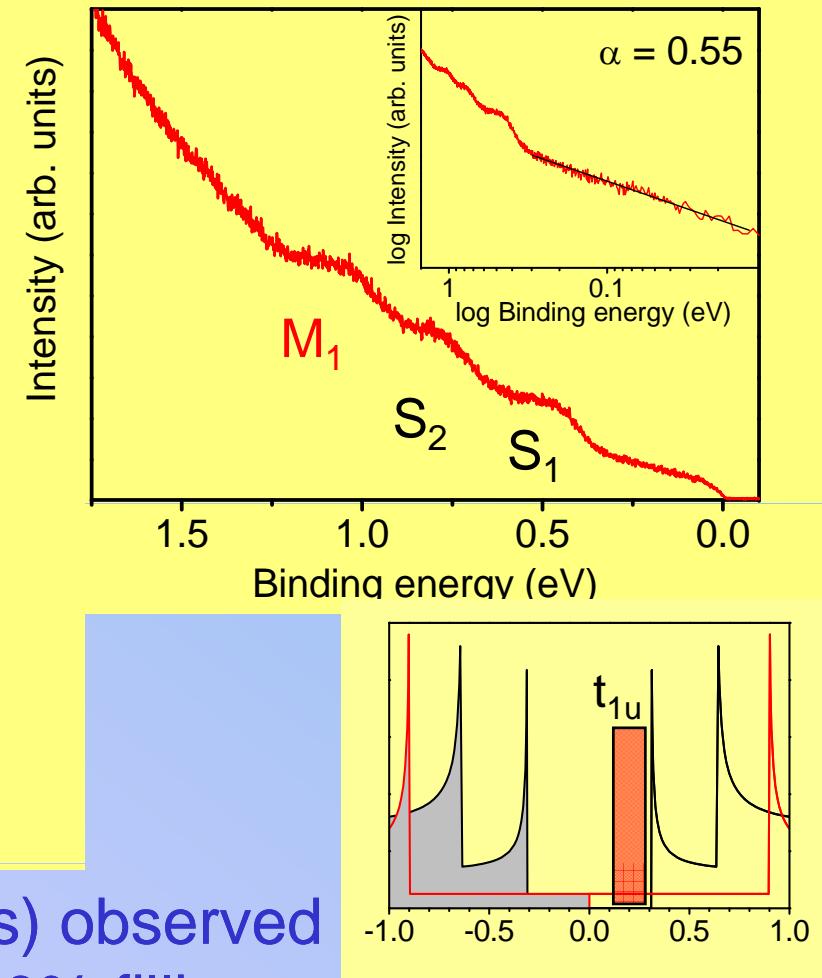
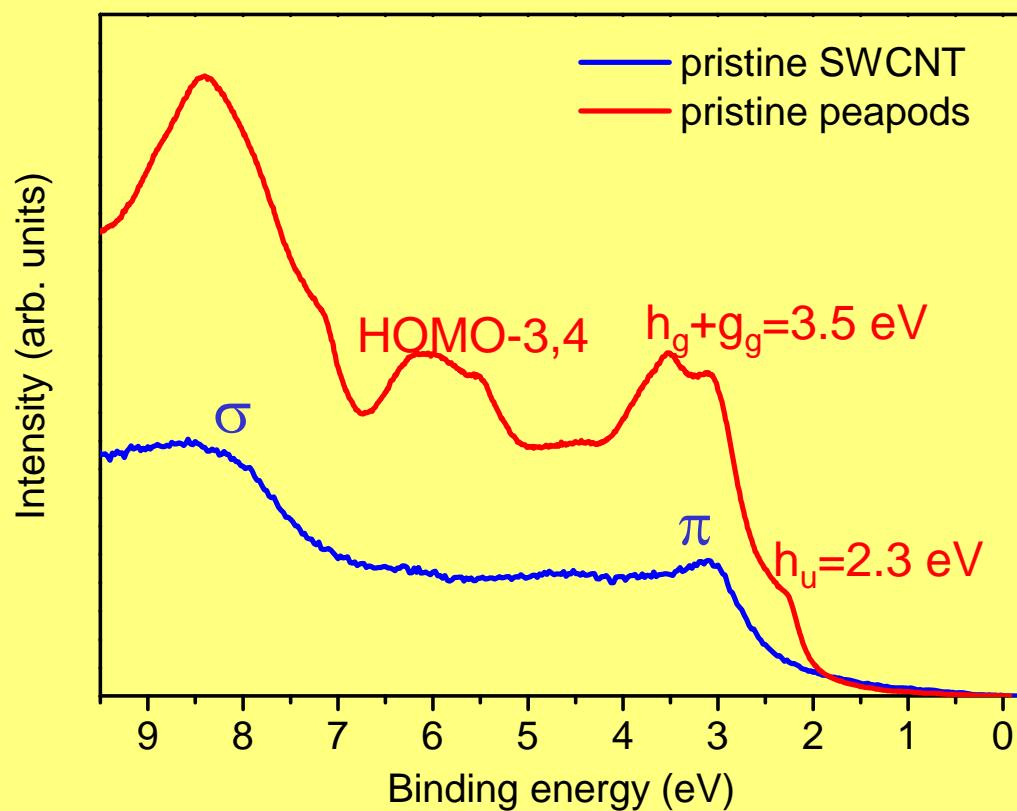


Filling/Intercalation



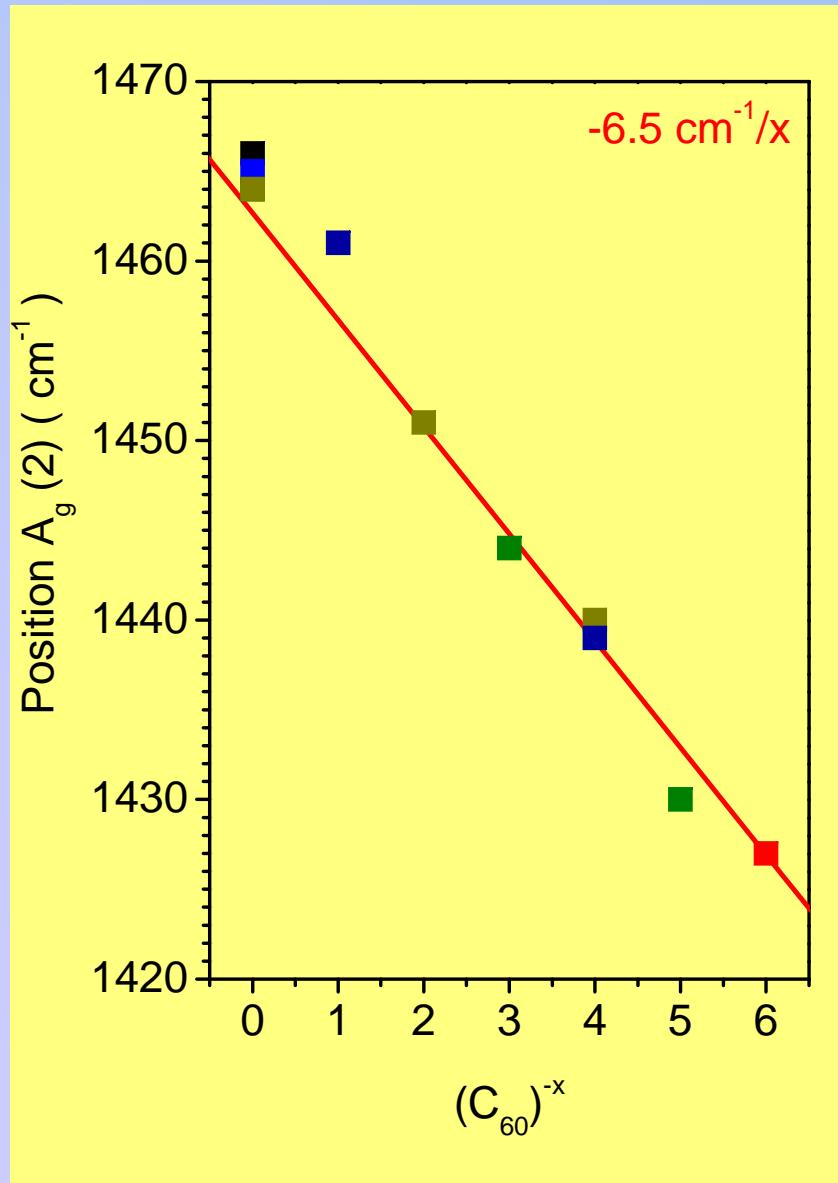
in-situ filling:
UHV evaporation
(5x10⁻¹⁰ mbar)
K, C₆₀ at 120/400 C
Equilibration 12 h
at RT/350 C
Annealing at
T=650 C

Electronic structure of C_{60} peapods



C_{60} derived molecular orbitals (Mos) observed
Binding energies as in bulk C_{60} , 100% filling
Power law $\alpha=0.48-0.55 \rightarrow$ basically no increase !!!

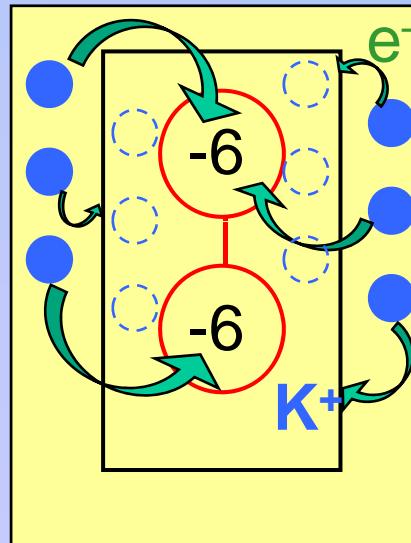
Combinational doping: Intercalation of C_{60} peapods



Competitive charge transfer to SWCNT
and C_{60} peas !

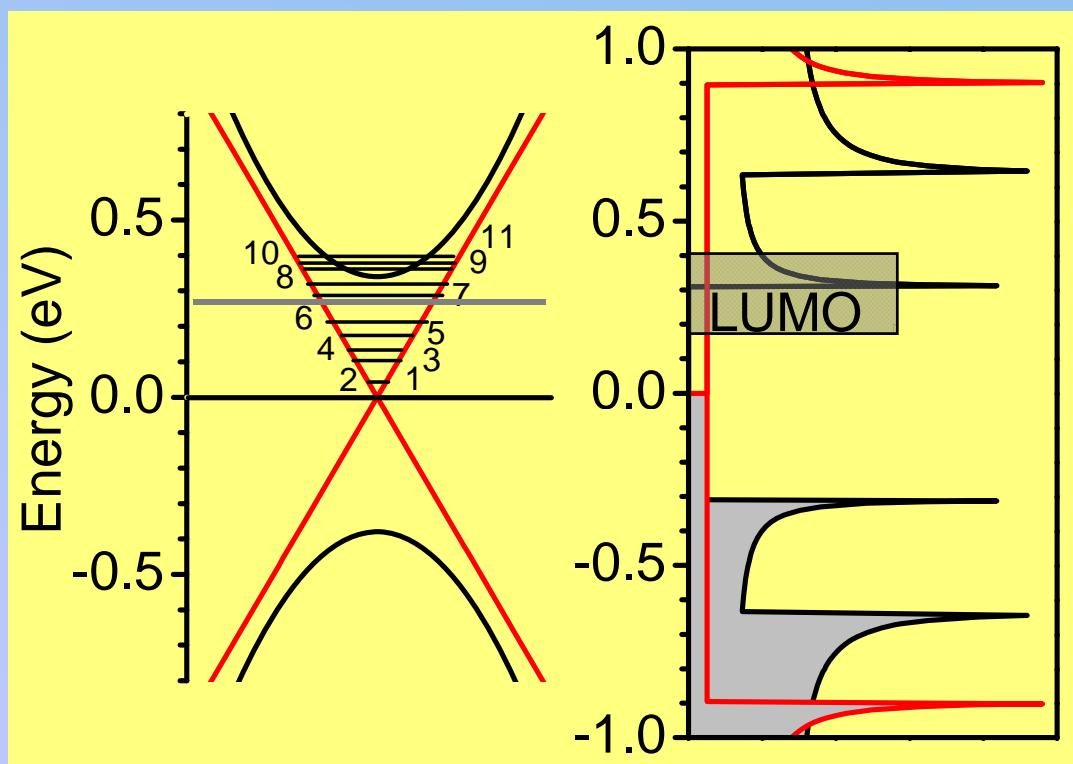
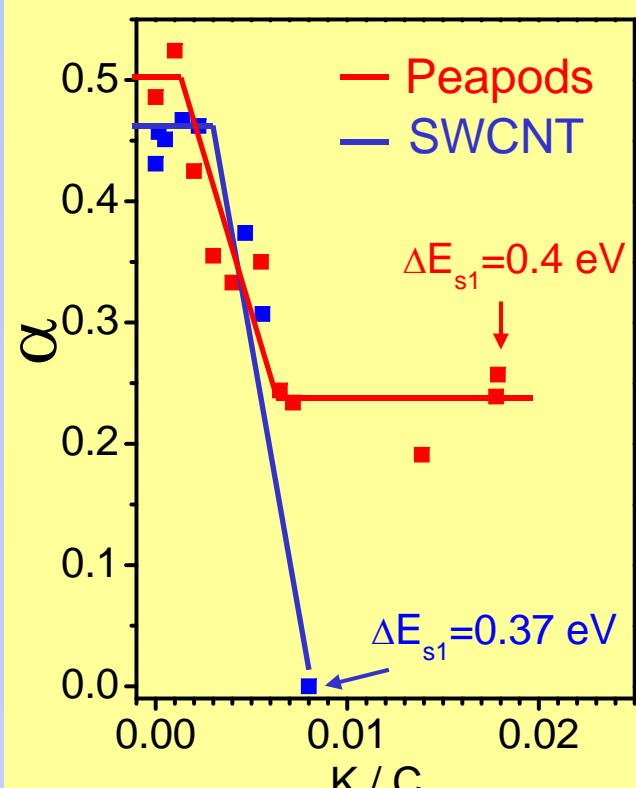
Raman: $A_g(2)$ of C_{60}^{-x} :

- charge transfer up to C_{60}^{-6}
- no line phases:
 $0.8 < x < 5.3$ average charge transfer
- charge transfer induced polymerization



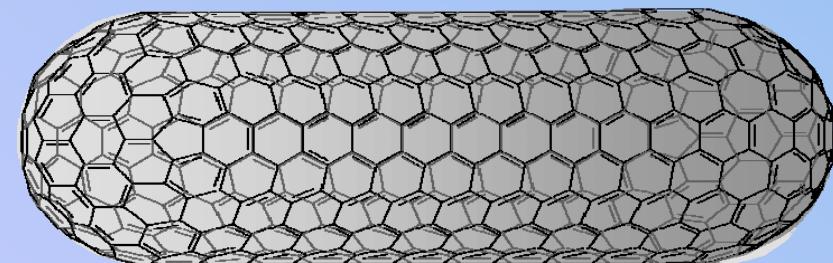
T. Pichler et al.,
PRL 87, 267401 (2001),
PRB 67, 125416 (2003).

C_{60} peapods at low doping close to E_F



→ Second conduction channel in doped C_{60} LUMO,
a bundle of metallic peapods is a Fermi liquid

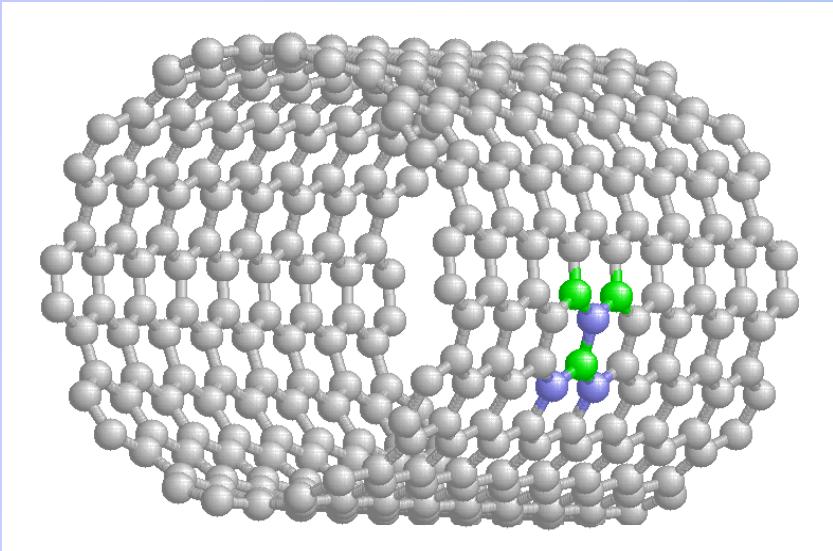
Substitution: Heteronanotubes



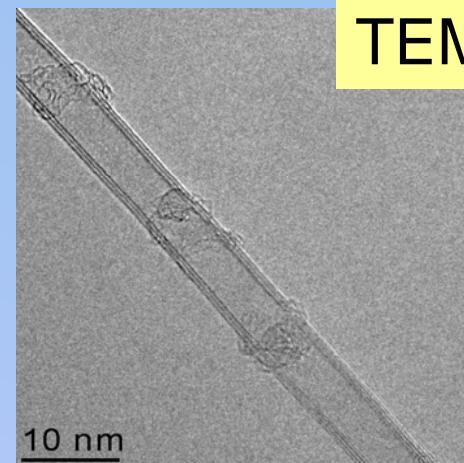
SWCNT



B,N Substitution

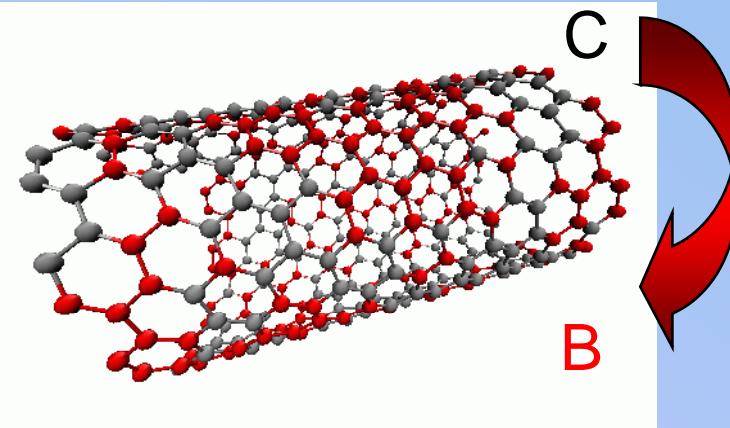


BCN-SWNT, MWBNNT



TEM of DWBN-NT

B-SWNT, BC₃NT



E. Borowiak-Palen et al. Chem. Comm. 1, 82 (2003), CPL, 387 , 516 (2003)
G. Fuentes et al. PRB 67, 35429 (2003),
PRB 69, 245403 (2004)

Summary Part 3

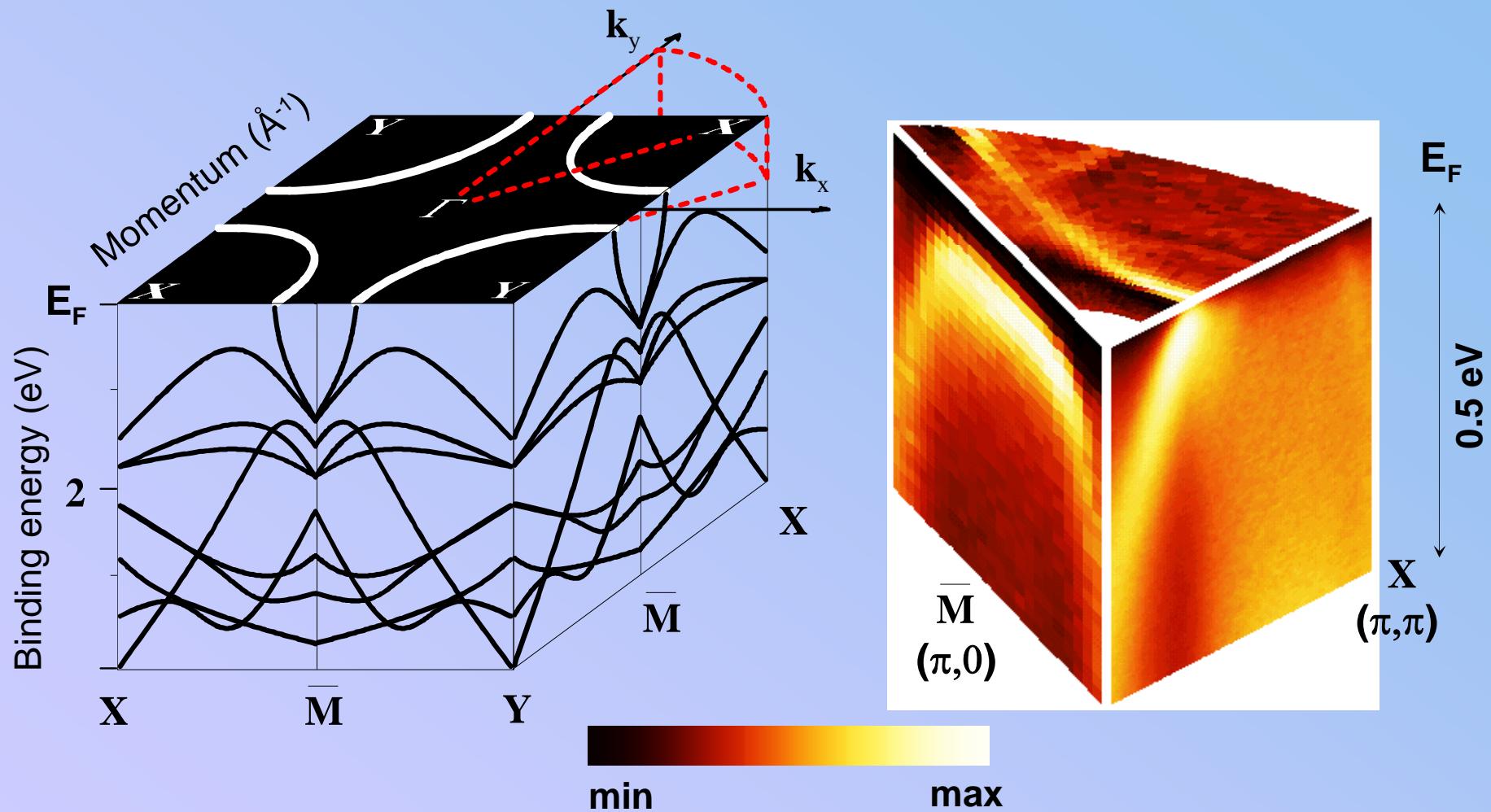
Tuning the electronic properties of SWCNT by functionalisation:

- p-, n- type intercalation of SWCNT bundles:
 - lattice expansion upon intercalation
 - direct measure of the doping level and charge carrier plasmon
 - Fermi level shift upon functionalization
 - interplay between charge transfer and hybridisation
 - Transition 1D-3D metal
- filled SWCNT, e.g. peapods:
 - bulk filling ratios, new chemistry inside a SWCNT nanoreactor
 - competitive charge transfer in intercalated peapods
 - Transition 1D-3D at higher doping compared to pristine SWCNT
- substituted SWCNT:
 - bonding environment
 - no ridged band shift upon doping
 - new BC₃ nanotubes

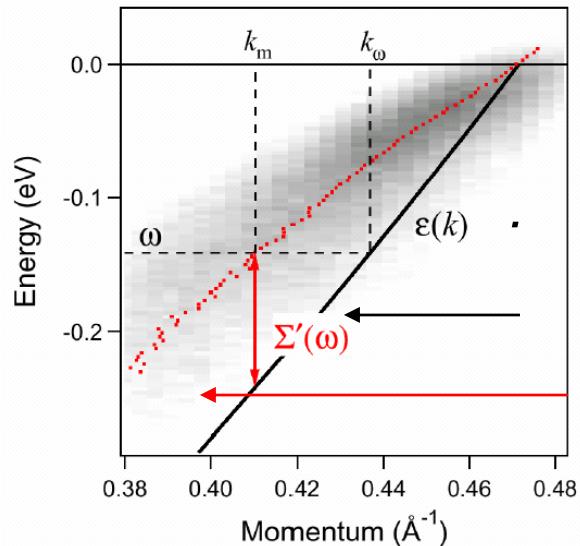
Outline

- Introduction/Motivation/Experimental
- Part 1: Angle resolved dielectric response of SWCNT
 - a) SWCNT bundles vs. individual SWCNT
 - b) Implications on excitation spectrum in 1D,2D
 - c) 1D „Drude“ plasmon in intercalated SWCNT?
- Part 2: Electronic structure of metallicity selected SWCNT
 - a) Textbook example for intrinsic XPS line shapes
 - b) Unravelling the 1D DOS in PES and XAS
 - c) Nature of metallic ground state in metallic SWCNT
- Part 3: Functionalized SWCNT:
 - a) Examples for intercalation, filling and substitution
 - b) Nanochemical reactions inside SWCNT
- Part 4: Electronic and phononic structure of graphene systems
 - a) Electron dispersion, phonon dispersion and EPC
 - b) GIC KC_8 revisited: A key to graphene
- Summary and outlook

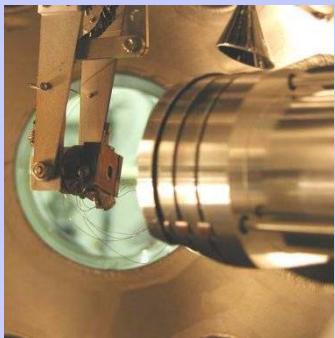
ARPES: Experimental momentum-energy space „cake“



ARPES measures quasiparticle dispersion



Bare band dispersion
Quasiparticle dispersion



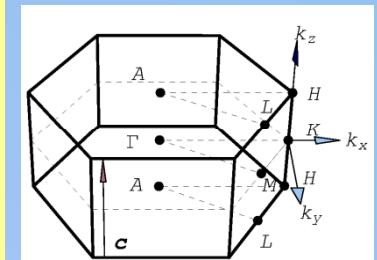
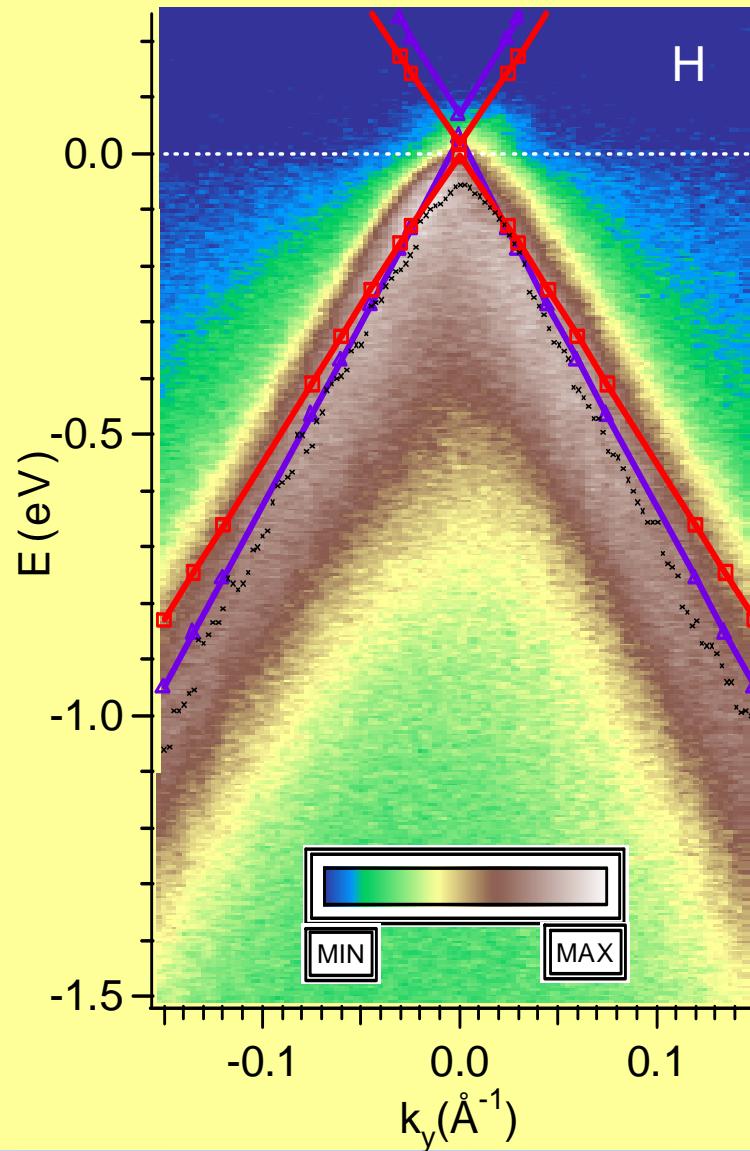
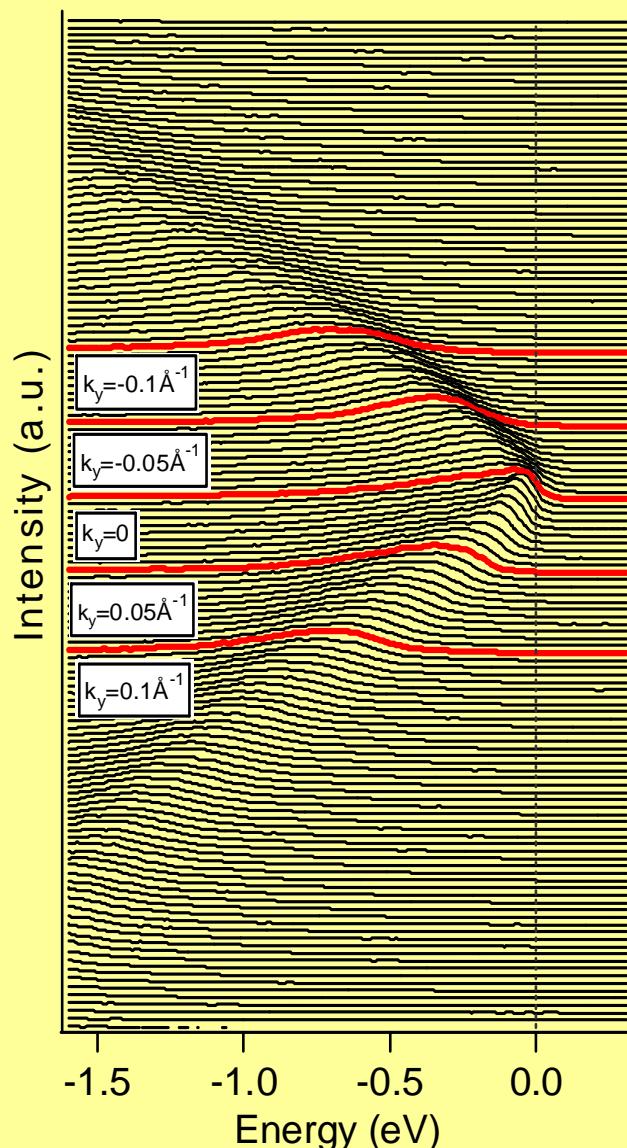
Single particle spectral function:

$$A(\mathbf{k}, \omega) = \frac{|\text{Im } \Sigma(\mathbf{k}, \omega)|}{(\omega - \omega_b(\mathbf{k}) - \text{Re } \Sigma(\mathbf{k}, \omega))^2 + (\text{Im } \Sigma(\mathbf{k}, \omega))^2}$$

Self energy correction Σ depends on:

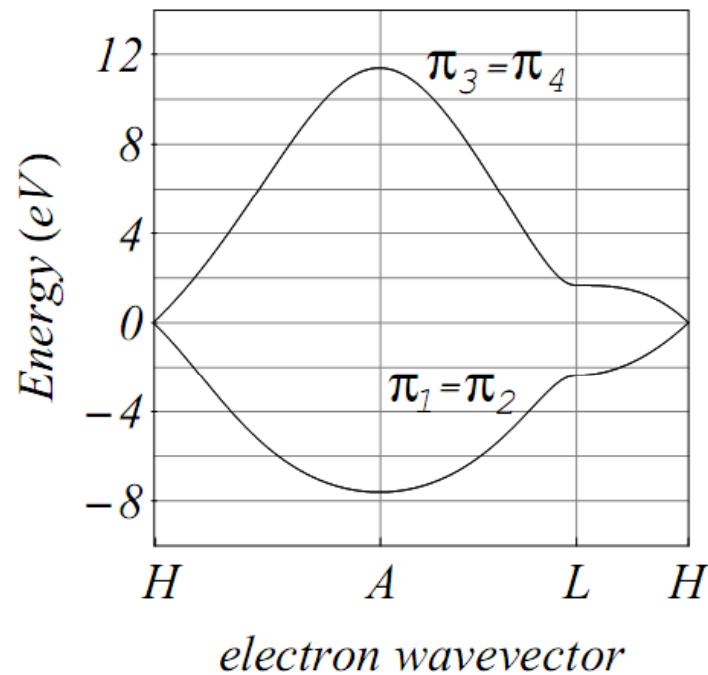
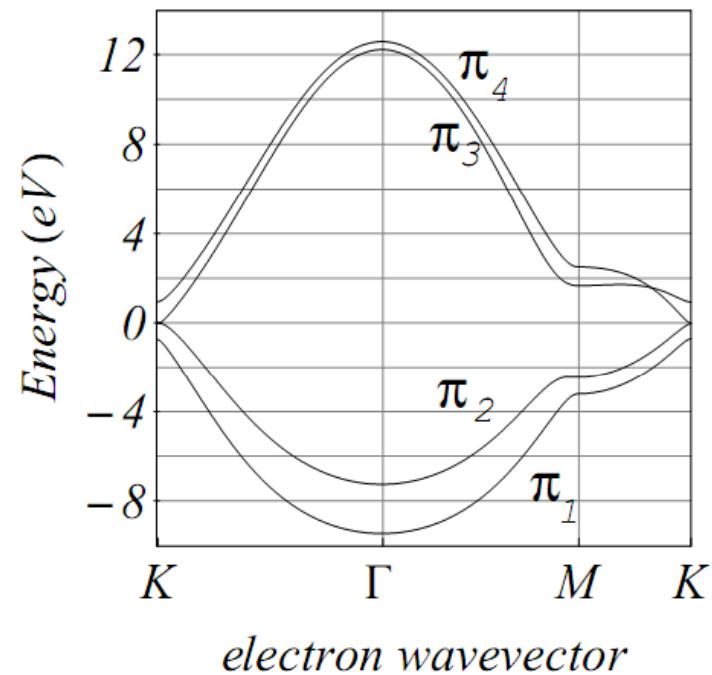
Electron-phonon, electron-plasmon,
electron-electron interactions,..

Electron-electron correlation in graphite



RED=LDA
BLUE=GW
BLACK DOTS
=MAX

K,H Point: New 3NN neighbor TB fit to GW



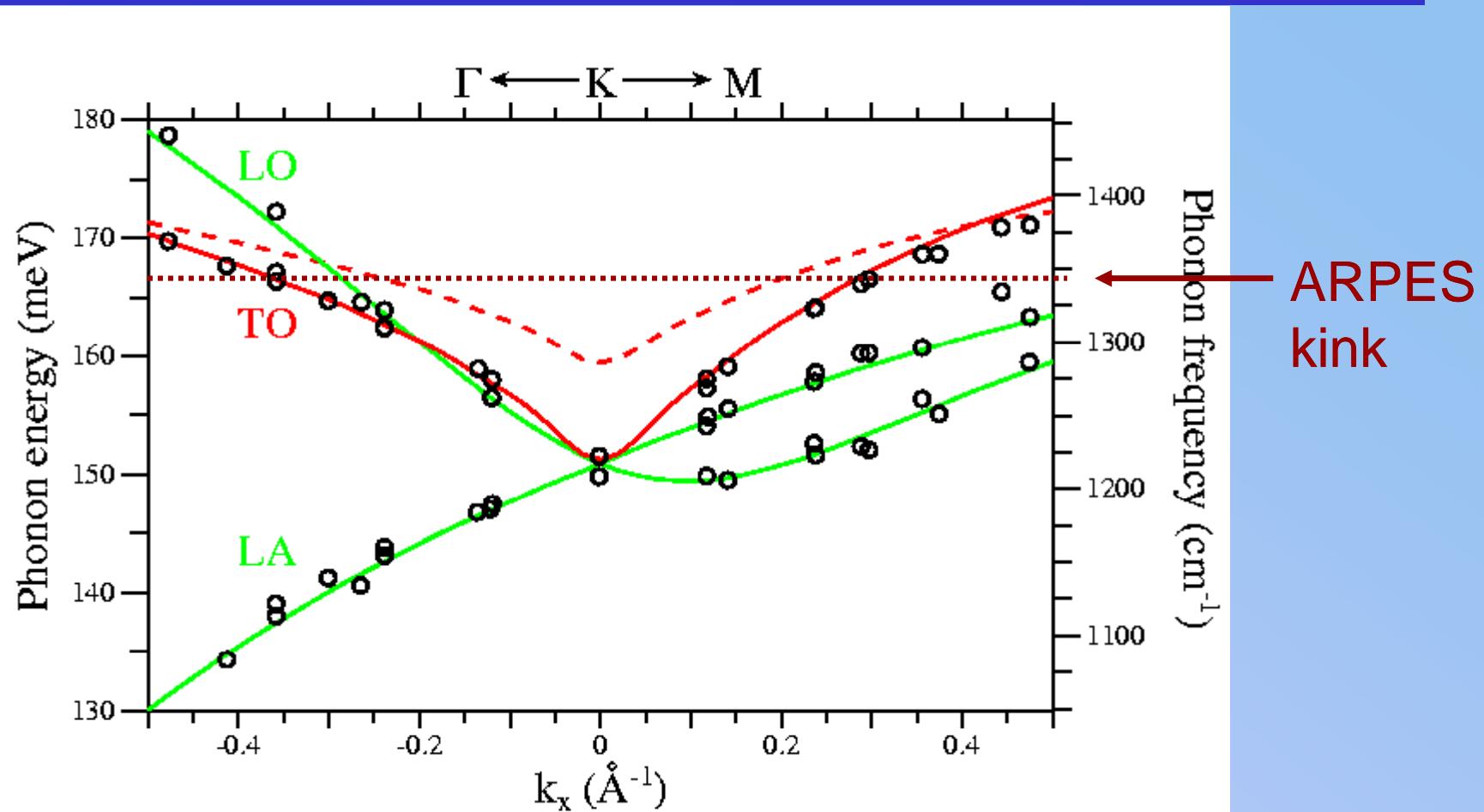
Method	γ'_0	γ'_1	γ'_2	γ'_3	γ'_4	γ'_5	E'_0	Δ'
TB-GW ^a	3.053	0.403	-0.025	0.274	0.143	0.030	-0.025	-0.005 ^b
TB-LDA ^a	2.553	0.343	-0.018	0.180	0.173	0.018	-0.022	-0.018
EXP ^c	3.16	0.39	-0.02	0.315	0.044	0.038	-0.024	-0.008
LDA ^d	2.598	0.364	-0.014	0.319	0.177	0.036	-0.026	-0.013
EXP ^e	2.9	0.3	-	0.1	0.12	-	-	
KKR ^f	2.92	0.27	-0.022	0.15	0.10	0.0063	0.0079	-0.027

Point	Method	π_1	π_2	π_3	π_4
Γ	GW	-9.458	-7.257	12.176	12.541
	TB-GW	-9.457	-7.258	12.184	12.540
M	GW	-3.232	-2.441	1.655	2.491
	TB-GW	-3.216	-2.457	1.656	2.495
K	GW	-0.736	-0.025	-0.025	0.917
	TB-GW	-0.728	-0.024	-0.024	0.909
H	GW	0.020	0.020	0.025	0.025
	TB-GW	0.020	0.020	0.025	0.025

Very good agreement of the GW QP band structure to experiment

A. Grüneis et. al., PRB 78, 205425 (2008)

Assignment of the coupling phonons by IXS

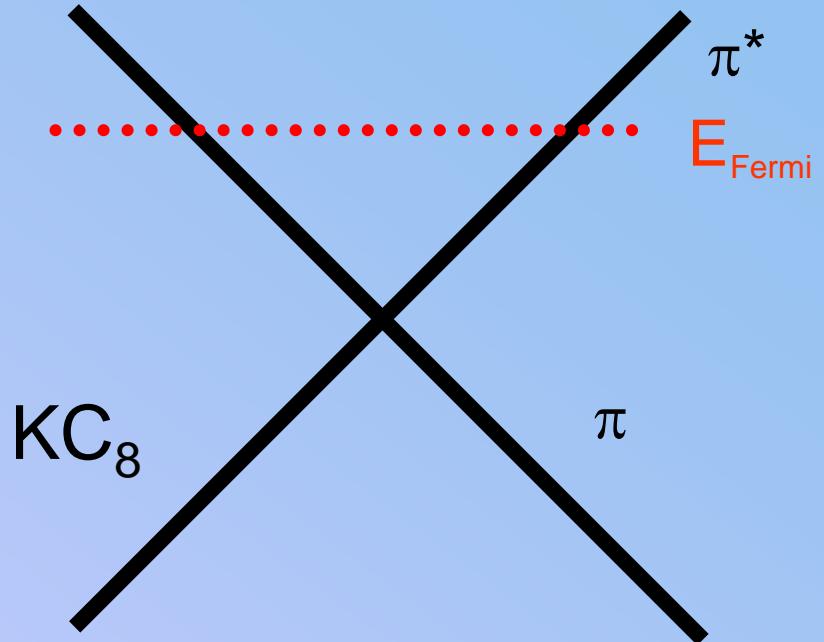
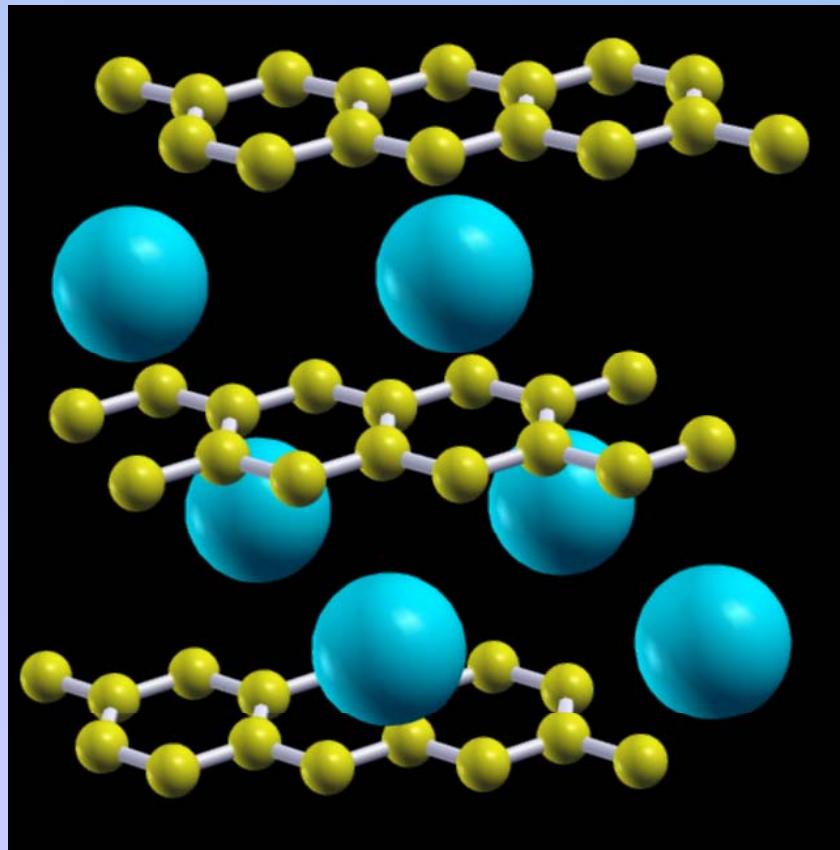


First unambiguous assignment of the TO/LO/LA branches,
Kohn anomaly at K using IXS.

First purely experimental EPC $166 (\text{eV}/\text{A})^2$

A. Grüneis et al. cond. mat. arXiv:0904.3205

KC₈ revisited: Interlayer decoupling by intercalation



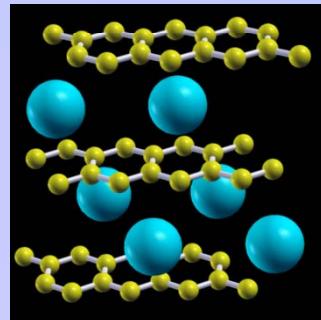
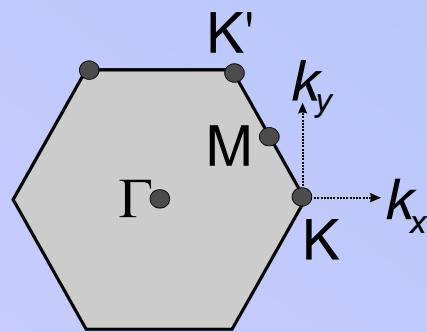
- 1) **Interlayer distance increases** from 0.335 nm to 0.535 nm.
- 2) Charge transfer K 4s to C π
- 3) First description by 2x2 Hamiltonian from Blinkowski.

(J. Blinowski, N.H Nguyen, Rigaux, C., Vieren, J.-P., J. Phys. 41, 47 (1980))

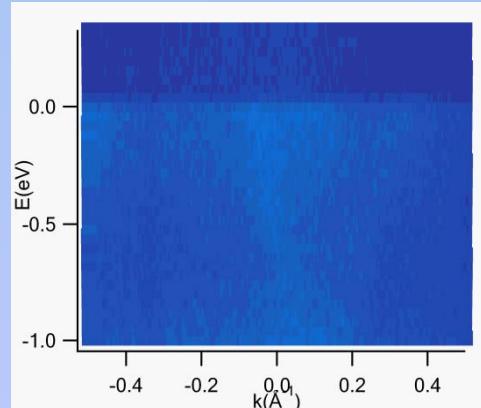
ARPES of KC_8 revisited



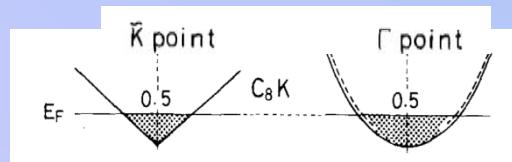
Golden KC_8 .



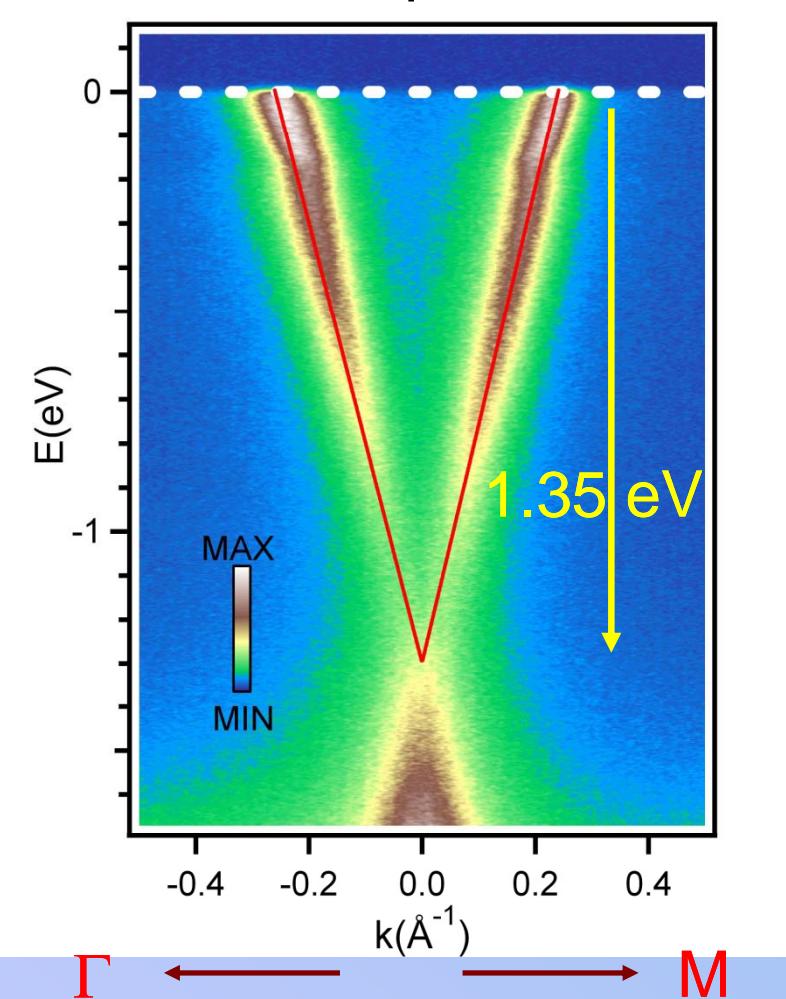
Γ point



Ionic charge transfer:
No 4s-like band

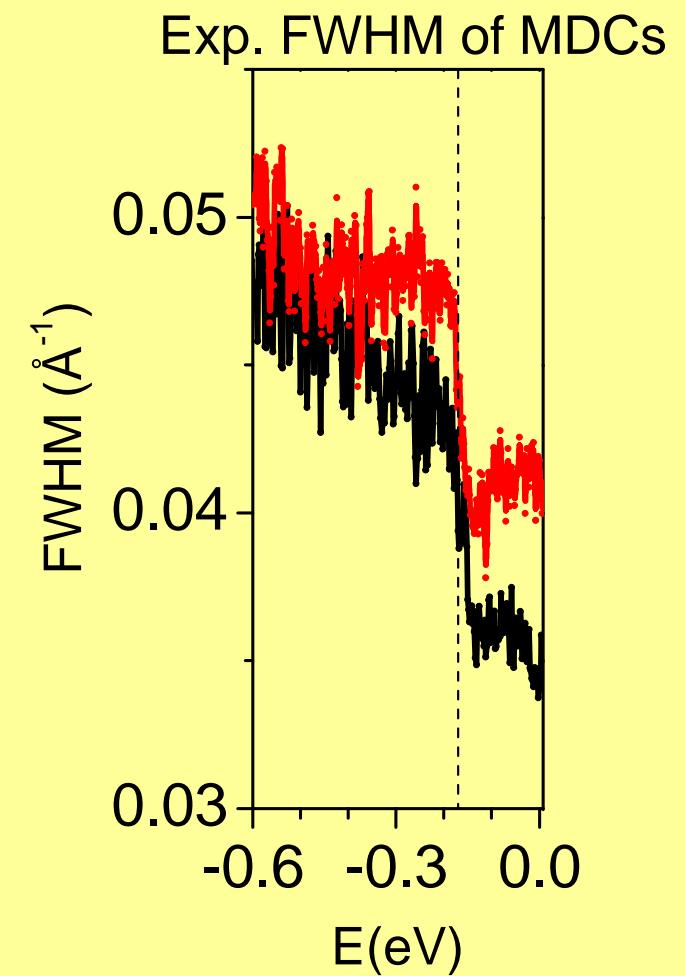
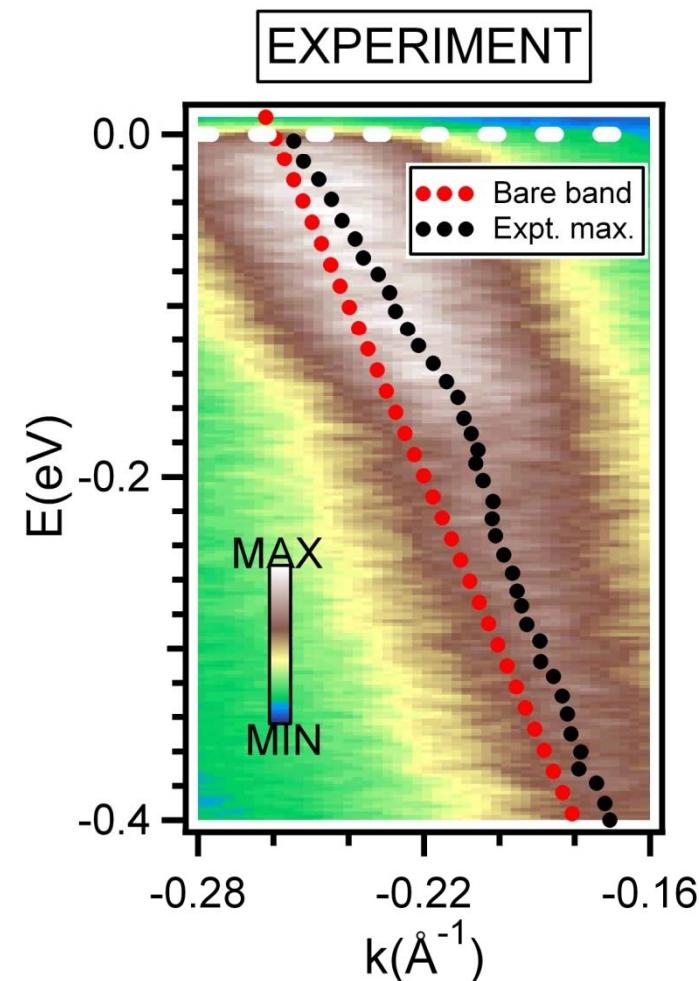
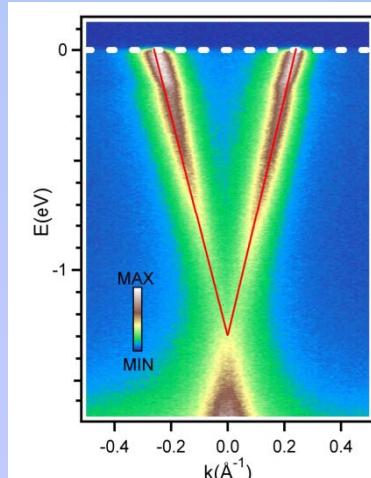


K point



- Kink in QP dispersion at 166 meV
→ coupling to a bosonic mode
- Dirac point shifts by 1.35 eV: linear dispersion of Dirac Fermions

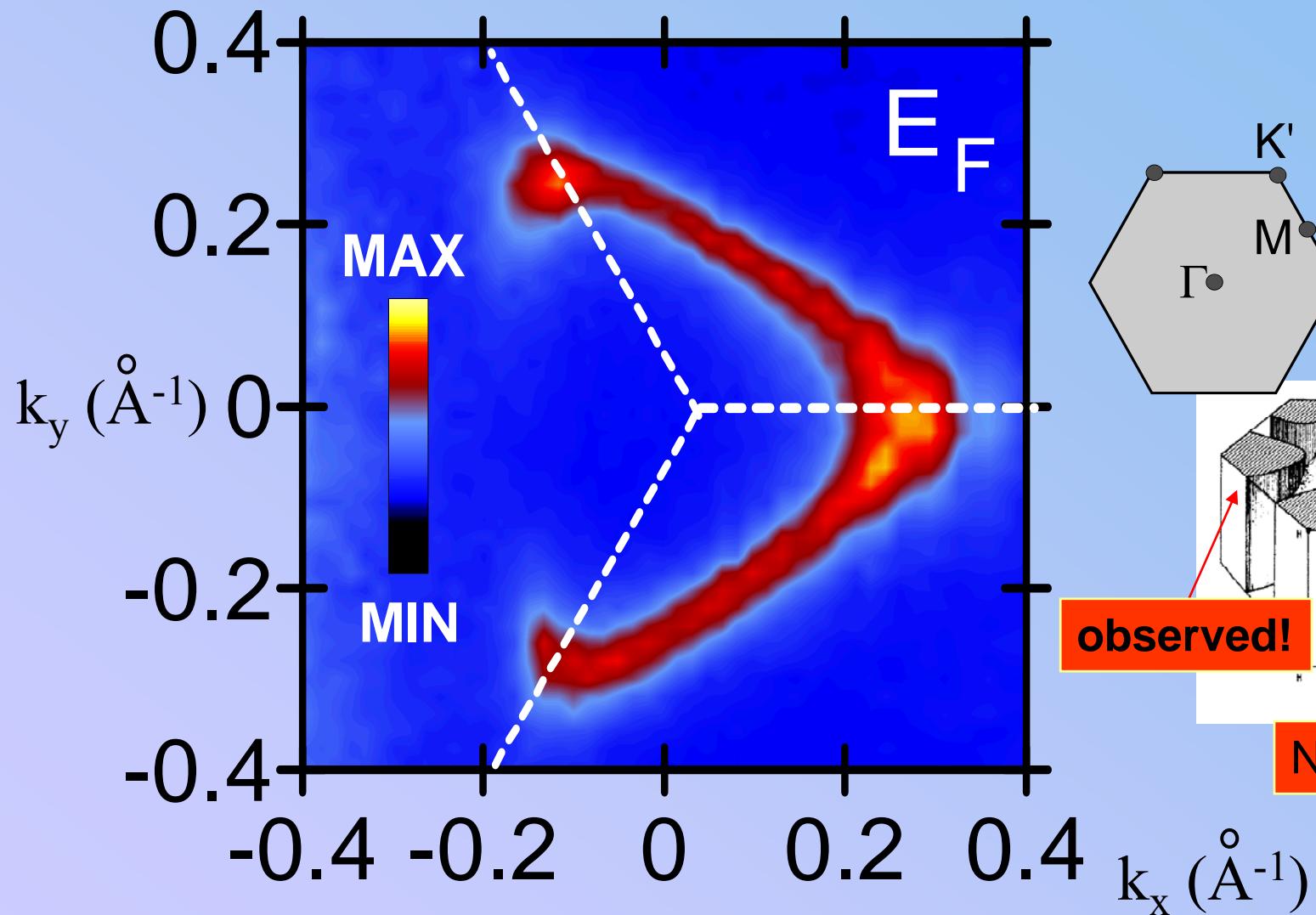
Self consistent self-energy analysis of the „kink“



- 1) Strong el-ph. coupling to TO phonon at K point at 166 meV
- 2) Direction dependent el-ph coupling constant 0.3-0.6.

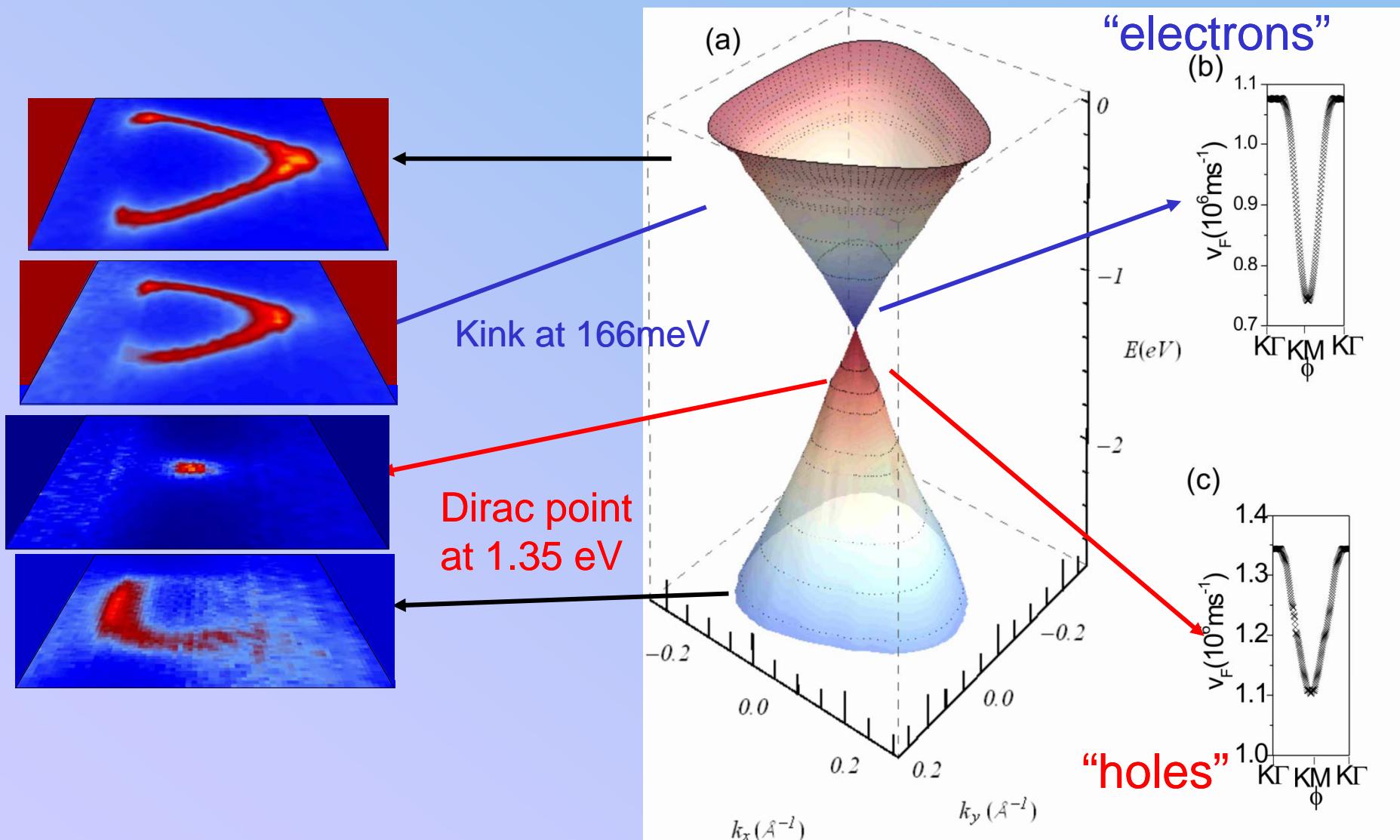
A. Grüneis et al. PRB 79, 205106 (2009)

Fermi surface of KC_8



- 1) Fermi surface only around K point.
- 2) $n = 7 \times 10^{21}$ electrons per cm^3 .

Experimental Dirac cone and Fermi velocities of doped graphene



A. Grüneis et. al., cond-mat. arXiv:0808.161; cond-matt. arXiv:0808.1467

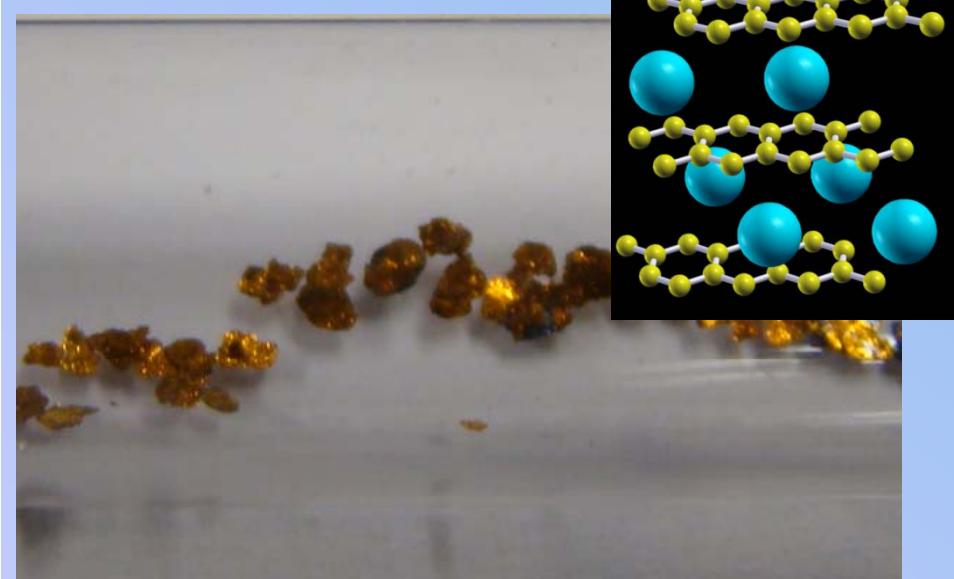
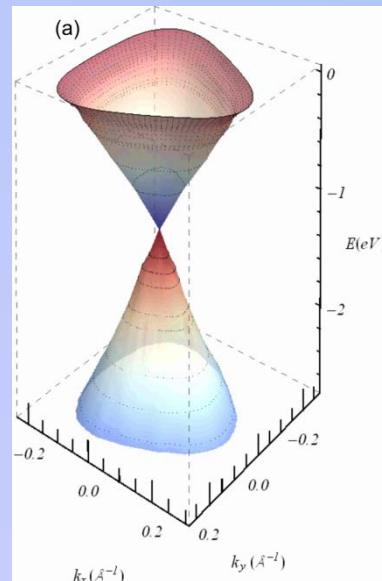
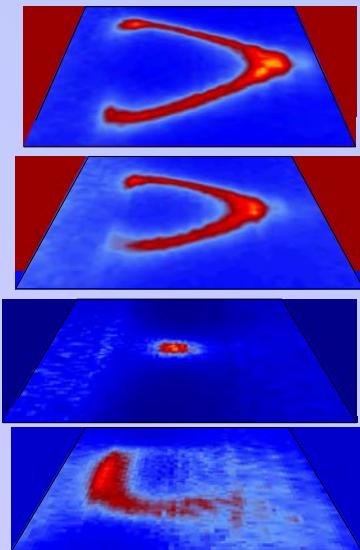
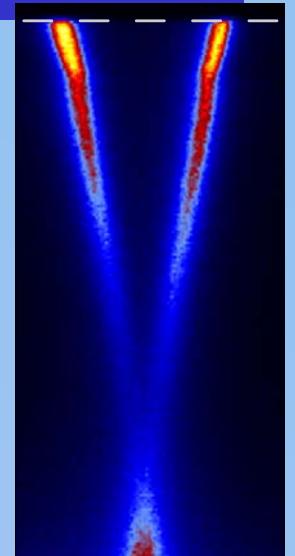
Summary Part 3:

- **Gain from GIC KC_8**

1. Charge transfer is completely ionic AA stacking and decoupled layers
2. K point TO phonon at 166 meV mediates superconductivity.
3. **Purely experimental Dirac cone!**

- **Gain from Graphite IXS, ARPES:**

1. el.el. Correlations important
2. optical K point phonons disentangled
3. experimental $EPC = 166 \text{ (eV/A)}^2$



A. Grüneis et. al., cond-mat. arXiv:0808.161; cond-mat. arXiv:0808.1467
C. Attaccalite et al. cond-mat. arXiv:0808.0786, PRB 78, 205425 (2009)

Outlook some open problems

- Low dimensional carbon systems:
- The role of instabilities in 0D, 1D, 2D systems:
 - Mechanical stability of low D systems
 - Growth of low D systems
 - Chemistry inside 1D and 2D confinement
- The role of basic correlation effects:
 - Electron-phonon (e.g. superconductivity)
 - Electron-electron (e.g. quasiparticle dispersion, TLL)
 - Electron-hole (e.g. Exciton, size, life time, binding energy)
 - Spin-spin (e.g. spin coherence, spin lifetimes in 1D)
- Tunable electronic and optical properties:
 - Key to engineer new molecular solids with novel electronic transport, optical, and magnetic properties
 - Interplay charge transfer vs. hybridization in functionalized metallicity selected SWCNT
 - Nanochemistry inside metallicity selected SWCNT.....

New research group
in Vienna:

LOW DIMENSIONAL QUANTUM SOLIDS



The beauty of instabilities...!!!

WANTED: *PhD student, postdoc eager to
unravel the equilibrium ground state
dressed by correlations.....!*

Summary

- SWCNT are truly 1D regarding optical response:
 - 1D on axis Plasmons: linear dispersion
 - Nondispersive surface plasmon at circumference
 - Charge carrier Plasmon: Non-dispersive, not acoustic??
Implications for all metallic 1D,2D systems
- Metallicity selected SWCNT:
 - Truly 1D, vHs, Tomonaga-Luttinger liquid
 - Very good agreement to diameter cumulative DOS
 - Key to understand interactions in functionalized SWCNT
Interplay between charge transfer and hybridization
- Electronic and vibronic structure of graphene systems:
 - Electron-Electron interaction important in graphene
 - Disentangled quasidegenerate TO, LA, LO at K point.
 - Purely experimental EPC = 166 (eV/A)^2
 - Fully experimental Dirac cone!
 - TO branch couples in GIC! Implications for sc?