

# **Exciton dephasing and multiexciton recombinations in a single carbon nanotube**

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## **ABSTRACT**

We studied the temperature and excitation intensity dependence of exciton luminescence in individual single-walled carbon nanotubes (SWNTs) using single nanotube photoluminescence (PL) spectroscopy. The linear temperature dependence of the PL linewidth in a single SWNT implies that the exciton dephasing is dominated by the interaction between the exciton and the phonon mode with very low energy under lower excitation conditions. Saturation of the PL intensity and broadening of the PL linewidth in a single SWNT occur simultaneously with an increase in excitation laser intensity. Our findings show that the rapid exciton-exciton annihilation through multiparticle Auger recombination broadens the homogenous PL linewidth.

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A single-walled carbon nanotube (SWNT) of about 1 nm diameter and length greater than several hundred nanometers is a prototypical system of one-dimensional (1D) structures [1]. The electronic and optical properties of SWNTs recently have attracted a great deal of attention [2], because the Coulomb interaction in the 1D system is strongly enhanced by the spatial carrier confinement [3]. These enhanced Coulomb interactions lead to the formation of tightly bound excitons. Following theoretical studies of the SWNTs [4-6], the experimental studies support the existence of stable excitons in SWNTs [7,8]. The characteristic optical properties of SWNTs are determined by the dynamics of 1D excitons with extremely large binding energies in the range of 200 to 400 meV [7,8]. It was also expected that the strong Coulomb interaction would enhance the many body effects of excitons in a SWNT. Indeed, the recent experimental results of pump-probe and transient photoluminescence (PL) spectroscopy have revealed the existence of strong exciton-exciton interactions under intense laser excitation [9-11].

Exciton-dephasing time and dephasing mechanisms play an important role in determining the optical properties of semiconductors and their device applications. In fact, the exciton-dephasing containing fruitful physics has been extensively studied in low-dimensional semiconductors [12,13], because these information will be necessary toward realization of quantum optical devices. However, the detailed exciton-dephasing mechanisms in 1D SWNTs are not clear. The spectral linewidth of a single SWNT contains essential information on exciton-dephasing time; therefore, photoluminescence (PL) measurements on single SWNTs [14-18] must be performed to clarify the intrinsic spectral linewidth hidden by the inhomogeneously broadened ensemble-averaged PL

spectra. In addition, the signatures of the exciton-exciton interaction and exciton dynamics appear in single SWNT PL spectra and their linewidths. The single carbon nanotube PL spectroscopy will provide new insight into the exciton-dephasing and exciton-exciton interactions in a SWNT, which cannot be obtained from ensemble-averaged PL and pump-probe spectroscopy.

In this paper, we examine the temperature and excitation power dependence of PL spectra from spatially isolated single SWNTs using single carbon-nanotube PL spectroscopy. The narrow Lorentzian PL spectrum is clearly observed, and the homogeneous PL linewidth linearly decreases with decreasing temperature. This linear temperature dependence implies that the exciton-dephasing is dominated by the interaction between the exciton and the phonon mode with very low energy. In the high excitation regime, the homogeneous linewidth broadens nonlinearly with an increase in excitation intensity. Our observation suggests that the broadening of homogeneous linewidth arises from the annihilation of excitons through a rapid Auger recombination process. The exciton-exciton interaction details concerning biexciton stability will be discussed.

The samples used in this study were spatially isolated SWNTs synthesized on Si substrates using an alcohol catalytic chemical vapor deposition (ACCVD) method [19]. The Si substrates were patterned with parallel grooves, typically  $\sim 1$   $\mu\text{m}$  wide and 500 nm deep using an electron-beam lithography technique. The isolated SWNTs grow from one side toward the opposite side of the groove or hang down toward the bottom of the

groove. Several kinds of samples were prepared by changing the growth temperature and time. We used the SWNTs sample without matrix and surfactant around the nanotubes to reduce the local environmental fluctuation effect [20]. Single nanotube PL spectroscopy was performed on a sample grown at 750 °C for 30 s with a typical average luminescent SWNT number density of  $0.1\text{-}1/\mu\text{m}^2$ .

Single SWNT PL measurements were carried out at temperatures ranging from 4.7 to 90 K using a home-built confocal microscope [18]. The SWNT samples were mounted in a cryostat and excited with either a He-Ne laser (1.96 eV) or a mode-locked Ti:Al<sub>2</sub>O<sub>3</sub> laser (1.72 eV). The pulse width and repetition rate of the Ti:Al<sub>2</sub>O<sub>3</sub> laser were 150 fs and 80 MHz, respectively. The laser light was focused on the sample surface through a microscope objective. The PL signal from SWNTs was detected by a 30-cm spectrometer equipped with a liquid-nitrogen (LN<sub>2</sub>)-cooled InGaAs photodiode array with a spectral range of 0.78 to 1.38 eV. The spectral resolution of the system was about 1.4 meV. The detector accumulation time was typically between 3 and 60 s.

Figure 1(a) shows the temperature dependence of the PL spectra of an isolated SWNT excited with a He-Ne laser. The spectra were measured under lower excitation conditions (less than  $\sim 100\ \mu\text{W}$ ) from 4.7 to 90 K. The SWNT has a chiral index of (10,3), based on the emission energies found in Refs. 21 and 22. The PL spectra clearly narrow when the temperature decreases [23]. As indicated by the circles in the figure, the spectral lineshapes are well fitted by Lorentzian functions. The linewidth (full-width at half-maximum, FWHM) of a single SWNT, corresponding to a homogeneous linewidth is

plotted as a function of temperature in Fig. 1(b). The temperature dependence of the linewidth shows almost linear behavior (as indicated by a solid line) in the range of 5 to 90 K. The inset of Fig. 1(b) shows the temperature dependence of the PL peak energy. The energy change from 5 to 90 K ( $\sim 1$  meV) is much smaller than that in other compound semiconductor quantum wire (12 meV from 5 to 90 K in GaAs wires) in the weak exciton-phonon coupling regime [24]. This weak temperature dependence is consistent with the experimental fact of narrow PL spectra with Lorentzian lineshape in Fig. 1(a).

Both the dephasing time and the lifetime of the excitons contribute to homogeneous linewidth broadening. The measured PL lifetimes of excitons are considerably longer ( $\geq 20$  ps) from 10 to 300 K [20], and this contribution to the linewidth is negligibly small ( $\leq 0.01$  meV) under lower excitation conditions. Thus, the temperature-dependent linewidths are determined approximately by the dephasing time of the exciton due to the exciton-phonon interactions. This linear temperature dependence implies that the very low energy phonon modes ( $\ll k_B T$ ) dominate the exciton-dephasing. The contributed low energy phonon to the exciton-dephasing is assigned as the TW (twisting) mode in SWNTs [18,25]. Based on the homogeneous linewidth, the exciton-dephasing time is evaluated from 350 fs at 90 K to more than 940 fs at 5 K [26].

Figure 2(a) shows PL spectra obtained from a typical single SWNT [assigned chiral index: (11,4)] at 30 K and at various excitation intensities of 1.72-eV and 150-fs laser pulses. Each spectrum has a single peak located at 0.941 eV. Even in the high-

excitation region above about 15 pJ per pulse, the PL spectrum shows a single peak, without a change in the peak energy. Furthermore, additional spectral structures are not observed in any lower energy regions. Similarly, the PL bands due to biexcitons (M-line) and inelastic exciton scattering (P-line) [3] are not observed in this spectral region.

Spectrally integrated PL intensities are plotted as a function of excitation laser intensity in Fig. 2(b). In a low excitation region below 10 pJ, the PL intensity grows almost linearly with excitation intensity (as indicated by a dotted line). Conversely, in the higher excitation intensity region ( $> 20$  pJ), saturation of the PL intensity is clearly apparent.

We show normalized PL spectra for a single SWNT, excited with different intensities, on an expanded energy scale in Fig. 3(a). The spectral linewidth broadens with increasing excitation intensity. The homogeneous linewidth is plotted in Fig. 3(b) as a function of excitation intensity, with the FWHM of the linewidth broadening nonlinearly. This broadening cannot be explained by laser heating effects because the nonlinear broadening behavior contradicts the linear temperature dependence of the linewidth broadening [see Fig. 1(b)]. This excitation intensity-dependent linewidth broadening indicates that the exciton dynamics are strongly affected by the multiple excitons present in a SWNT.

Under higher excitation conditions, when additional homogeneous linewidth broadening and saturation of the PL intensity occur, more than one exciton is created in a

SWNT, leading to a remarkable scattering process between excitons. Exciton-exciton annihilation occurs due to the rapid Auger nonradiative recombination process through strong Coulomb interactions [27-29]. The exciton annihilation rate ( $1/\tau_{\text{ex-ex}}$ ) is given by [10]

$$1/\tau_{\text{ex-ex}} = A \cdot L \cdot n \cdot (n-1), \quad (1)$$

where  $A$  is the exciton annihilation (Auger) constant,  $L$  ( $\sim 1 \mu\text{m}$ ) the length of the SWNT, and  $n$  the time-averaged exciton number. The exciton number should be a discrete value. However, the exciton number fluctuates at each single excitation event. Then, we might consider that the time-averaged exciton number  $n$  changes continuously in the analysis of the time-averaged PL data. Taking the radiative lifetime ( $\tau_{\text{rad}} > 1 \text{ ns}$ ) [30] and the nonradiative lifetime, due to the exciton trapping of the nonradiative centers ( $\tau_{\text{non-rad}} \sim 40 \text{ ps}$ ) [20] into consideration, the PL intensity  $I_{\text{PL}}$  at the excitation power  $P$  is described by,

$$I_{\text{PL}} \propto \tau_{\text{rad}}^{-1} \cdot P / (\tau_{\text{rad}}^{-1} + \tau_{\text{non-rad}}^{-1} + \tau_{\text{ex-ex}}^{-1}). \quad (2)$$

The PL intensity saturation at higher excitation regions, as shown in Fig. 2, can be explained by the opening of the nonradiative relaxation path due to the exciton-exciton annihilation.

The homogeneous linewidth  $\Gamma$  in Fig. 3(b) is determined by both the excitation power independent exciton-phonon interaction term  $\Gamma_{\text{ex-ph}}$  and the excitation power dependent Auger term  $\Gamma_{\text{ex-ex}}$  ( $\propto 1/\tau_{\text{ex-ex}}$ ),

$$\Gamma = \Gamma_{\text{ex-ph}} + \Gamma_{\text{ex-ex}}. \quad (3)$$

The homogeneous linewidth broadening at higher excitation regions, as shown in Fig. 3, can be explained by the shortening of exciton lifetime due to the Auger process. The

nonlinear behaviors of the PL intensity and the homogeneous linewidth broadening can be reproduced using Eq (1)-(3) as indicated by the solid lines in Fig. 2(b) and 3(b). The time-averaged exciton number  $n$  can be translated from the excitation intensity  $P$  at which the PL intensity saturates. The best fitted curves in Fig. 2(b) and Fig. 3(b) with varying the fitting parameter  $A$  reproduced the experimental results well. Then, the derived exciton-exciton annihilation constant  $A$  of  $1.6 \text{ ps}^{-1}\mu\text{m}$  is almost consistent with the theoretically calculated value based on perturbation theory [10]. From this result, it was found that the multiparticle Auger process occurs very efficiently with an Auger recombination time estimated at 800 fs for  $\sim 1 \mu\text{m}$ -long SWNT when two excitons are present in a SWNT. This very short, in comparison to bulk semiconductors, time constant is characteristics of low-dimensional systems, including SWNTs and is comparable to the exciton-dephasing time at low temperature. These results indicate that the exciton dephasing is limited under high excitation conditions, not by the exciton-phonon interactions, but by the exciton-exciton interactions.

Under the highest excitation conditions in this study, shown in Fig. 2(a), no spectral structures arising from exciton complexes, such as a biexciton, were observed in a lower energy region of the exciton PL line. However, Pedersen *et al.* [31] have discussed the exciton-exciton interaction and the stability of the biexciton in a SWNT based on the effective-mass theory. The biexciton binding energy is estimated at 37 meV in a semiconducting SWNT with a 1.07 nm diameter [corresponding chirality: (11,4)]. This theoretical calculation predicts that the two excitons energetically form the bound biexciton when they are present in the SWNT. However, PL signals related to the

biexciton emission were not experimentally observed, even under high excitation and low temperature conditions. In other pump-probe spectroscopy experiments, no transient absorption signal related to the biexciton has been observed [11]. This inconsistency between theoretical prediction and the experimental results is attributable to the dynamic behavior of the excitons. As discussed above, the nonradiative Auger recombination process occurs efficiently at about 1 ps order. As a result, one of the two excitons nonradiatively relaxes in the ground state when two excitons are present in a SWNT. The exciton-exciton scattering rate in a SWNT is much higher than in other 1D semiconductors, such as GaAs quantum wire, in which biexciton emission has been observed experimentally [32]. The biexciton formation from two uncorrelated excitons and the exciton-exciton scattering process are compensated in a SWNT. Thus, the biexciton formation in a SWNT is obstructed because of the very efficient Auger recombination process resulting from the strong Coulomb interaction in 1D systems.

In conclusion, we performed a single carbon nanotube PL spectroscopy on single isolated SWNTs by changing the temperature and the excitation intensity. The homogeneous linewidths decreased linearly with a decrease in temperature. The temperature dependence implies that the exciton-dephasing time was dominated by the interaction between the exciton and the phonon mode with very low energy. In the high-excitation regime, the homogeneous linewidth broadened nonlinearly with an increase in excitation intensity. Our observation suggests that the exciton-exciton annihilation through the rapid Auger recombination process contributes to the homogeneous

broadening in a SWNT. It is pointed out that the instability of a biexciton in a carbon nanotube is due to the efficient exciton annihilation process.

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## FIGURE CAPTIONS

**Fig. 1** (color online) (a) Temperature dependence of PL spectra of a single SWNT from 4.7 to 90 K excited with He-Ne laser light. Solid circles denote the Lorentzian function fits. (b) Spectral linewidths (full-width at half-maximum, FWHM) as a function of temperature. The dotted line represents the spectral resolution. The solid line indicates almost linear temperature dependence. Inset shows the temperature dependence of the PL peak energy.

**Fig. 2** (color online) (a) Low-temperature PL spectra obtained from a typical single SWNT [assigned chiral index: (11,4)] at various excitation intensities for 1.72-eV and 150-fs pulses. (b) Integrated PL intensity as a function of the excitation laser intensity. As indicated by the dotted line, under weak excitation conditions, the PL intensities show almost linear power dependence. The solid line corresponds to the fitted curve.

**Fig. 3** (color online) (a) Normalized PL spectra of a single SWNT on an expanded energy scale excited with different intensities. (b) Homogeneous linewidths as a function of excitation intensity. The solid line corresponds to the fitted curves considering exciton annihilation through the exciton-exciton scattering process.

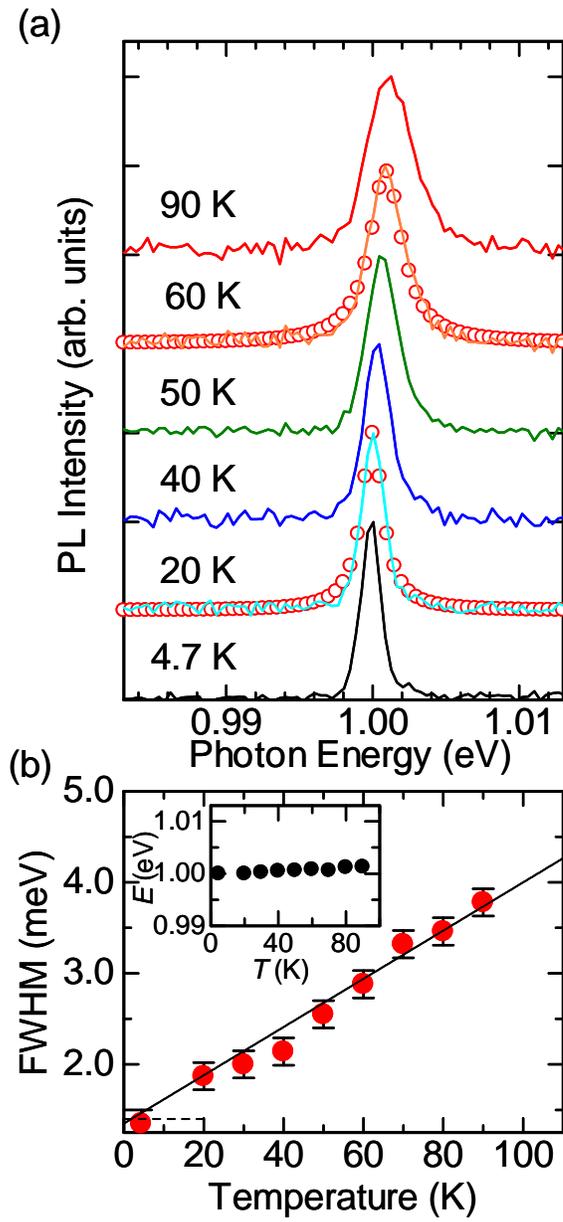


Fig. 1 K. Matsuda et al.

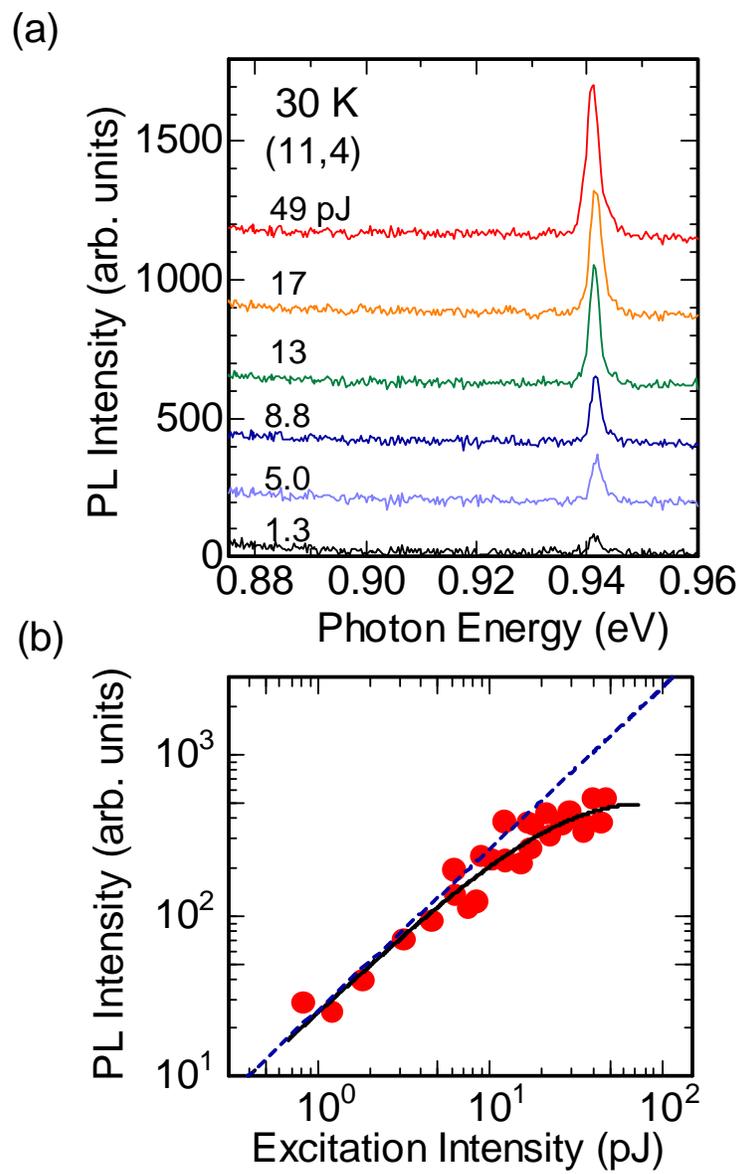


Fig. 2 K. Matsuda et al.

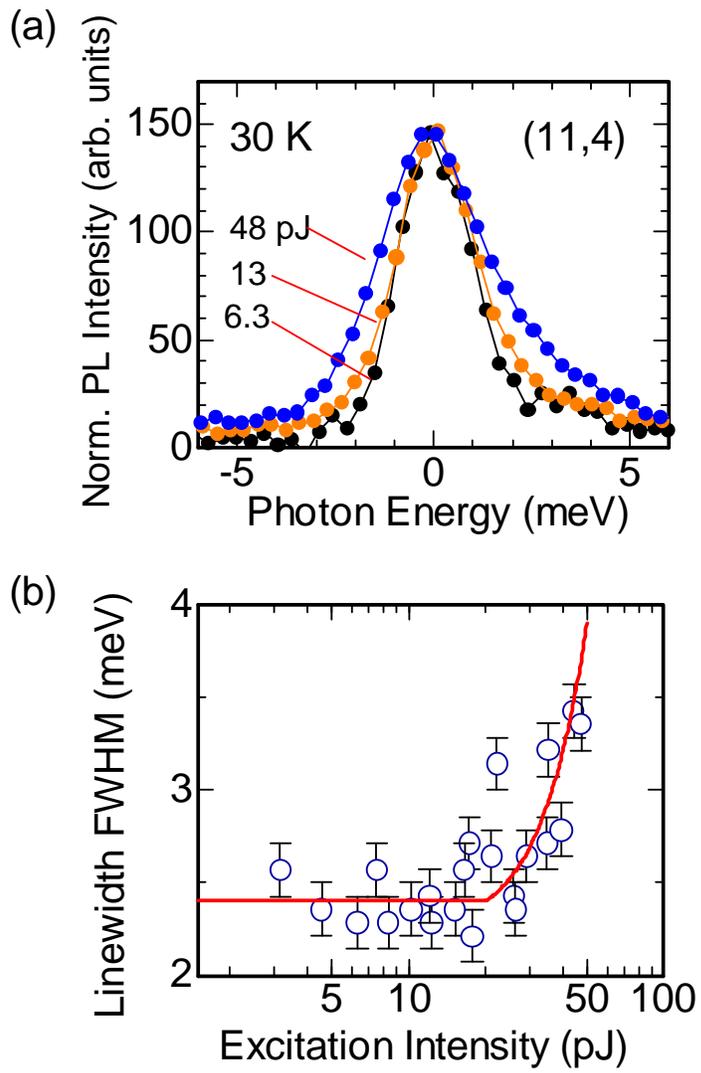


Fig. 3 K. Matsuda et al.