

New thermal mechanisms in sub-10nm structures

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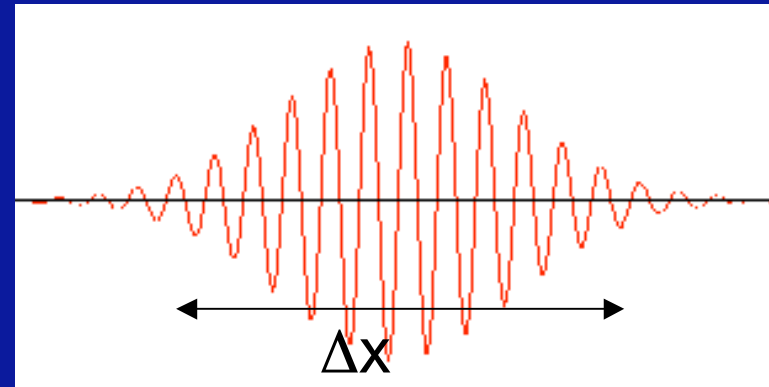
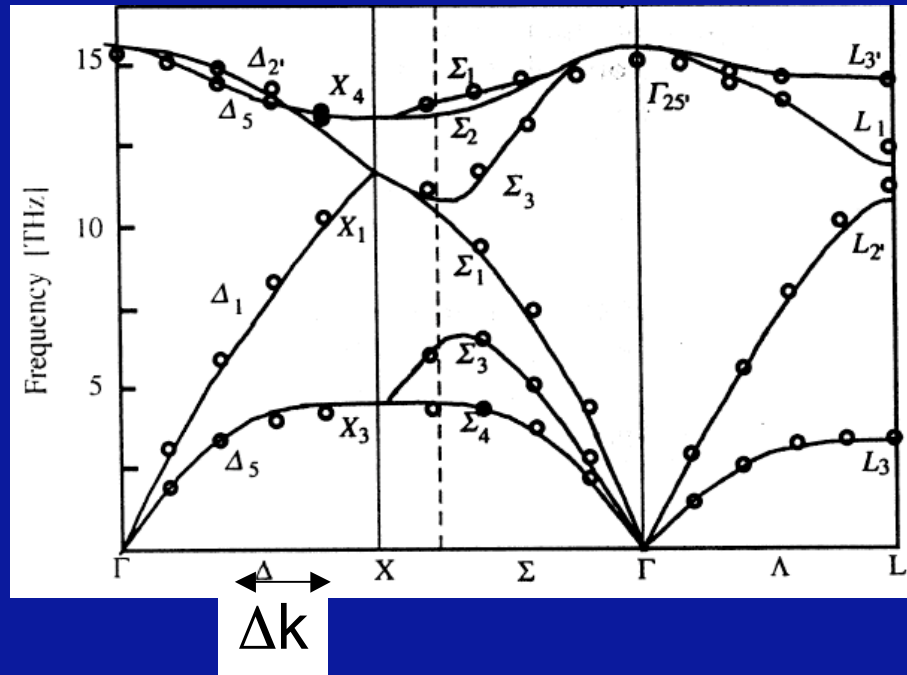
CNRS - Ecole Centrale Paris - France



Seminar at Mechanical Engineering Dpt, University of Tokyo, September 26th

PHONON PARTICLE/PHONON WAVE

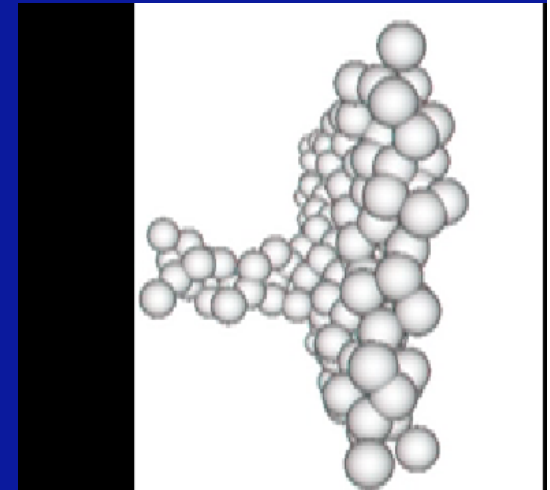
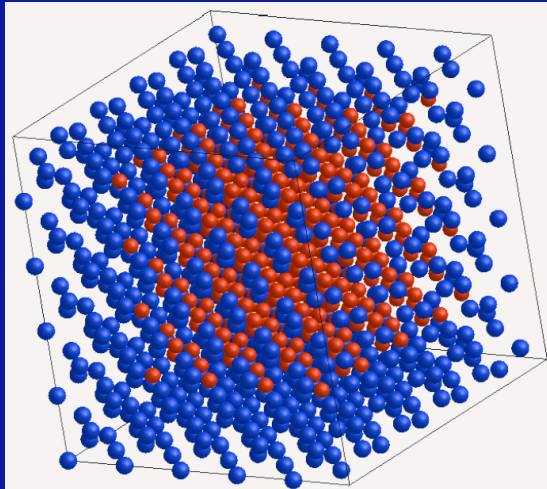
$$\Delta k \cdot \Delta x > 2\pi$$



$$\Delta k < p/a \Rightarrow \Delta x > 10a = 5\text{nm}$$

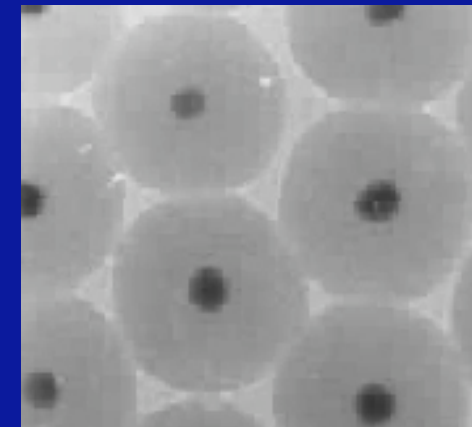
New thermal mechanisms in sub-10nm structures

1. Heat Conduction in Nanowires

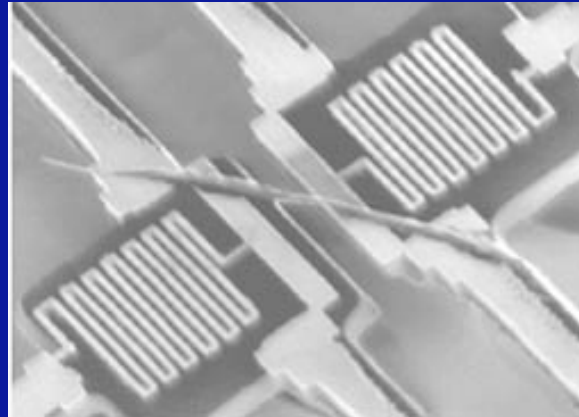
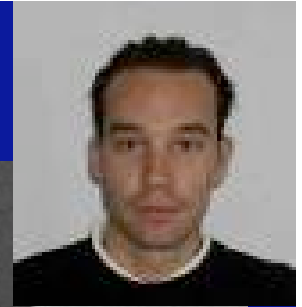


2. 3D atomic scale phononic crystals

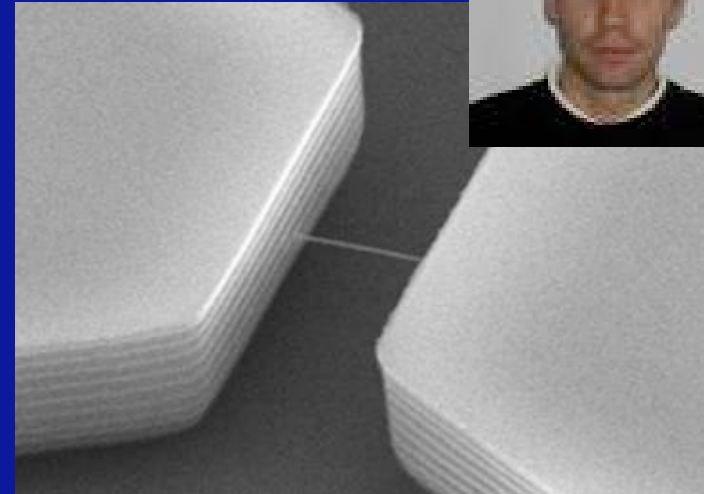
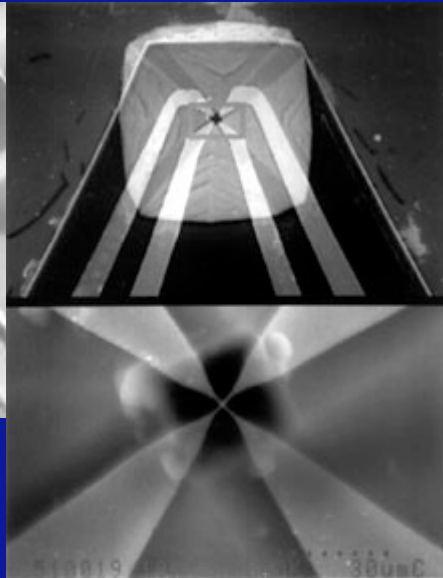
3. Transient Non Fourier Heat Conduction



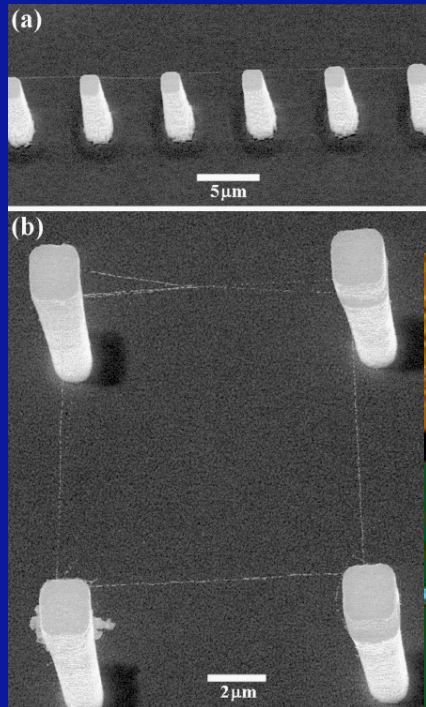
1-Nanowires and Thermal Contact Resistances



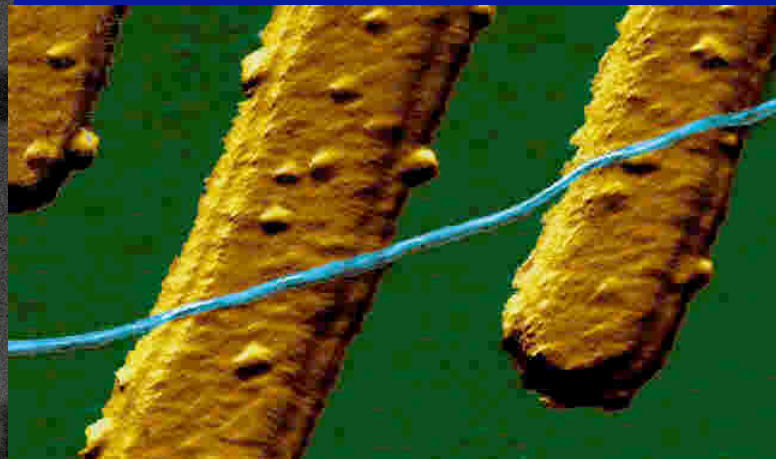
Metrology



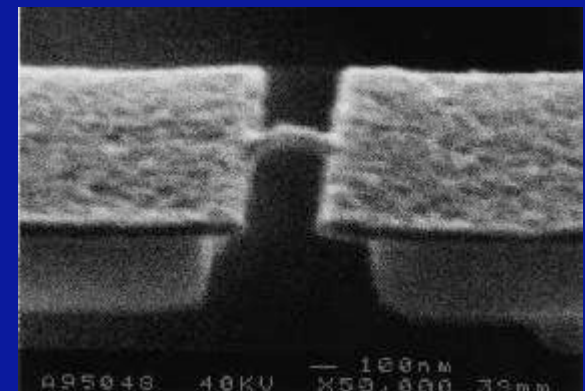
Spectrometry



NanoElectronics



Molecular Junction



Macro-FOURIER

$$\Delta T^+ = 0$$

Thermal Bath T_1

$$\nabla T^+$$

Thermal Bath T_0

$$\nabla T^+_{\infty} \frac{1}{\lambda DR_F}$$

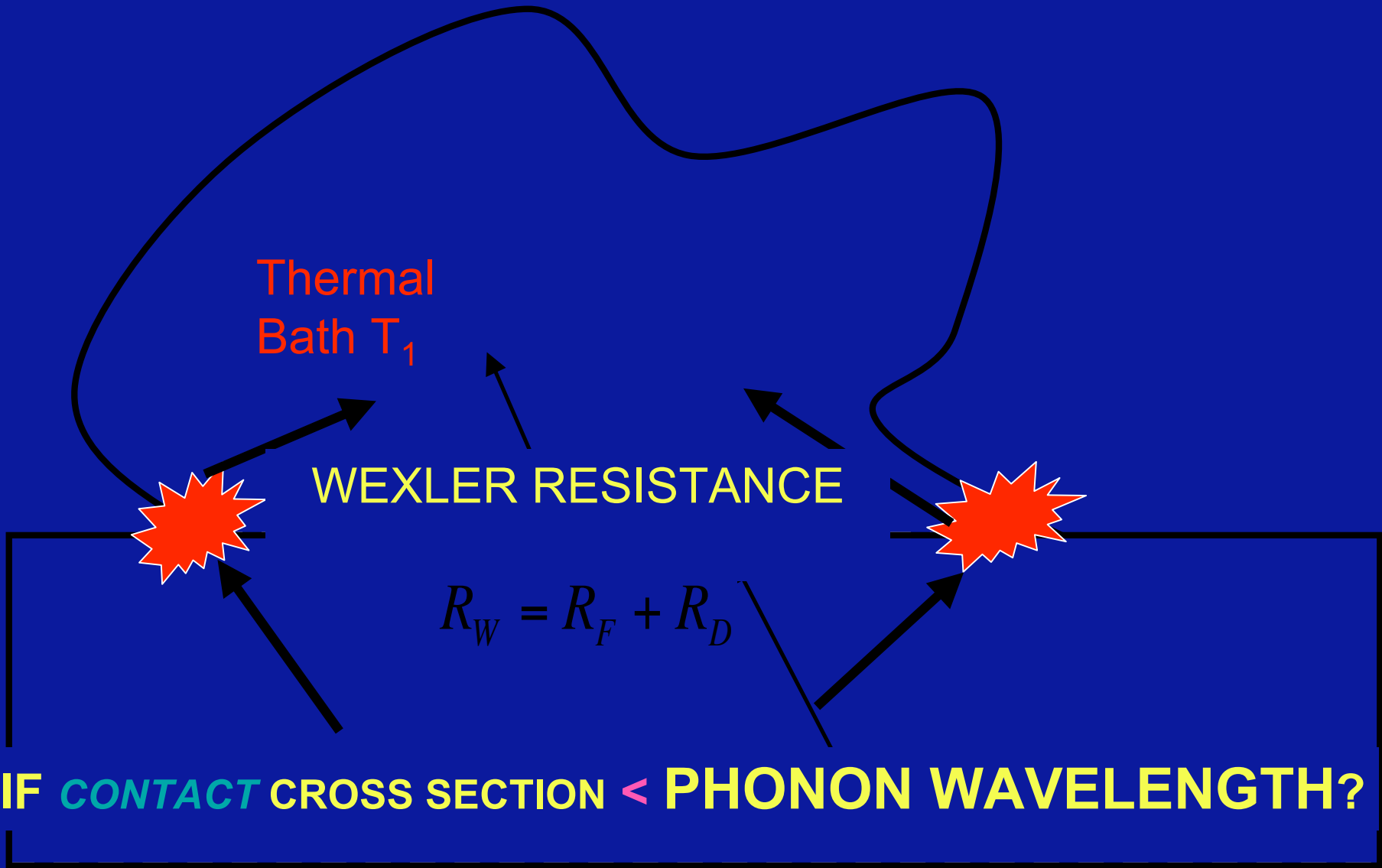
NanoContact Kn > ≈ 1

Thermal
Bath T₁

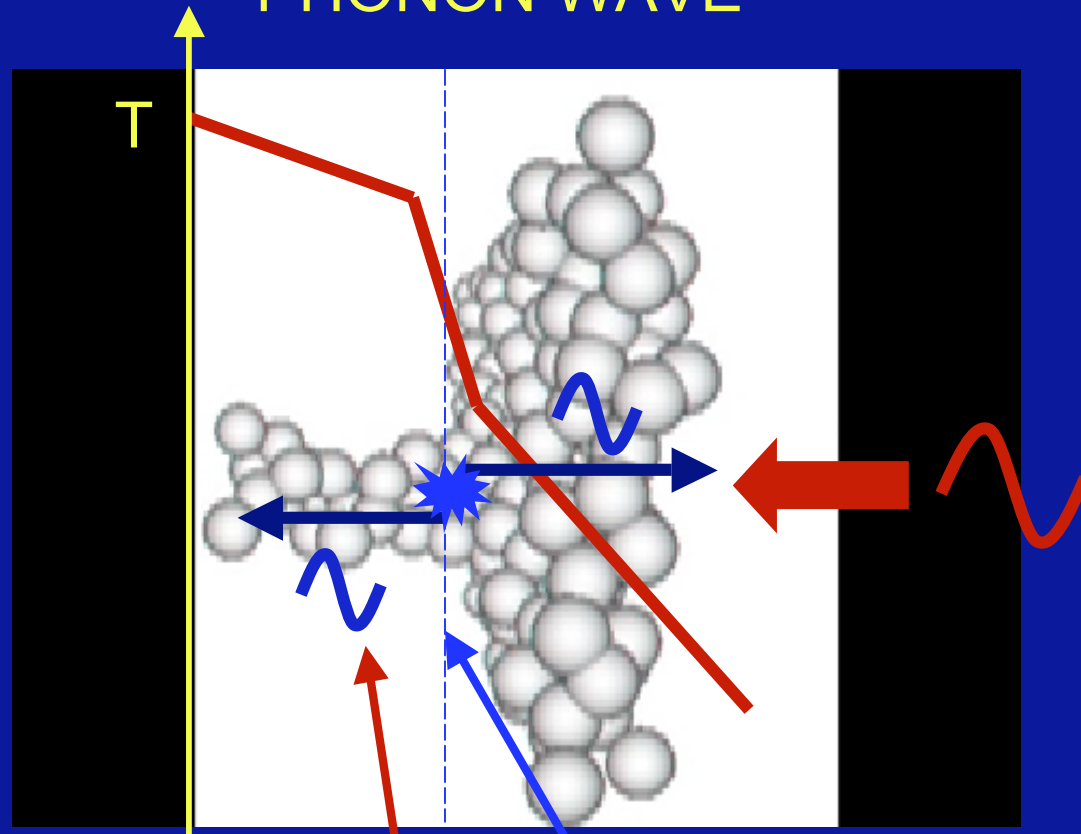
WEXLER RESISTANCE

$$R_W = R_F + R_D$$

IF CONTACT CROSS SECTION < PHONON WAVELENGTH?

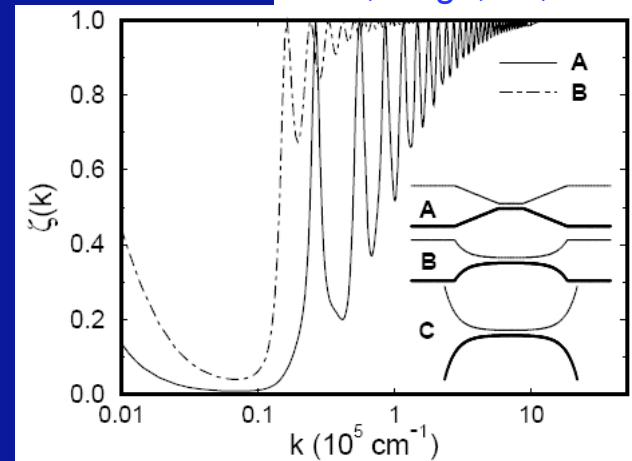


PHONON WAVE

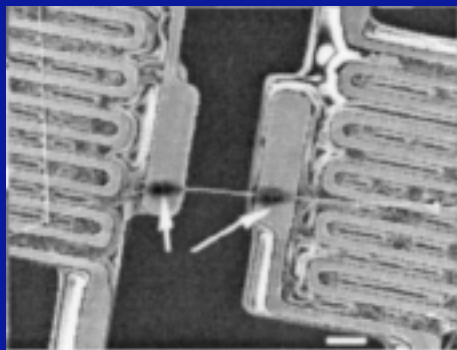
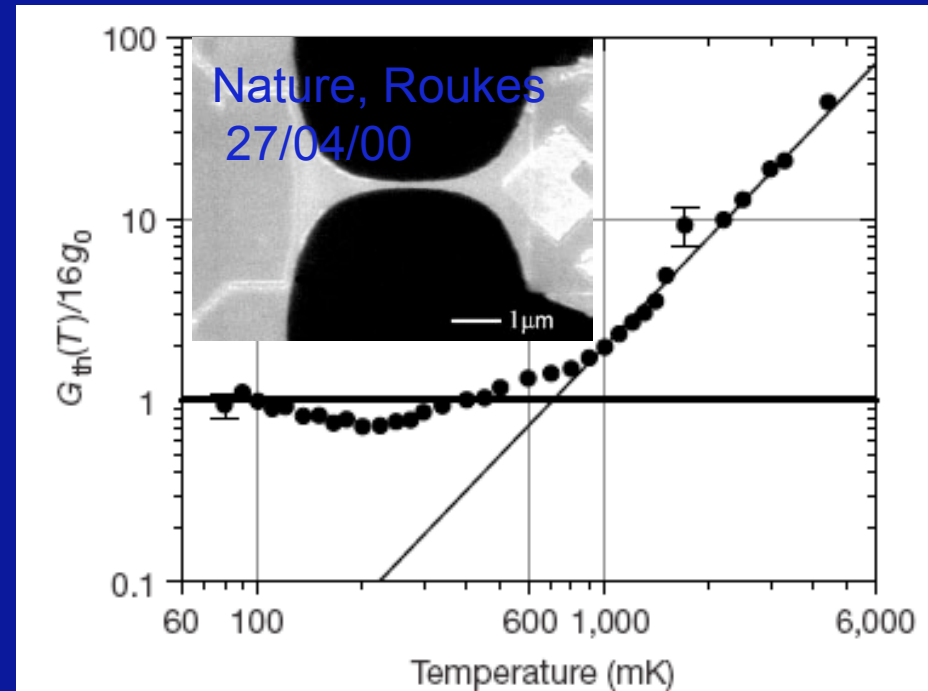
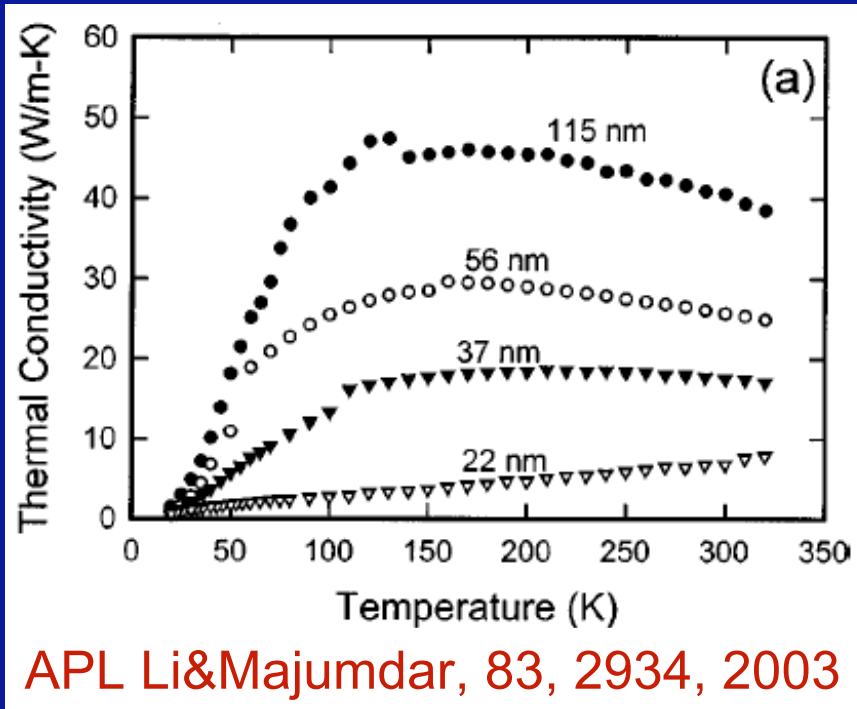


PRL, Rego, 81, 232

$$\dot{Q} = \sum_{\alpha} \int_0^{\infty} \frac{dk}{2\pi} \hbar \omega_{\alpha}(k) v_{\alpha}(k) (\eta_R - \eta_L) \zeta_{\alpha}(k)$$



The Mysterious Story of Nanowires Thermal Conductivity



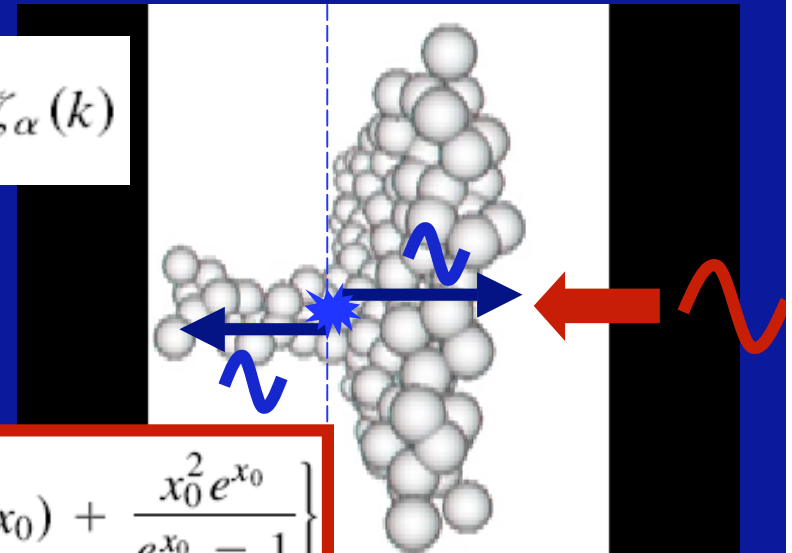
M.R. says:

not obtained for long wires

Wire Thermal Conductance

$$\dot{Q} = \sum_{\alpha} \int_0^{\infty} \frac{dk}{2\pi} \hbar \omega_{\alpha}(k) v_{\alpha}(k) (\eta_R - \eta_L) \zeta_{\alpha}(k)$$

1D in k-space

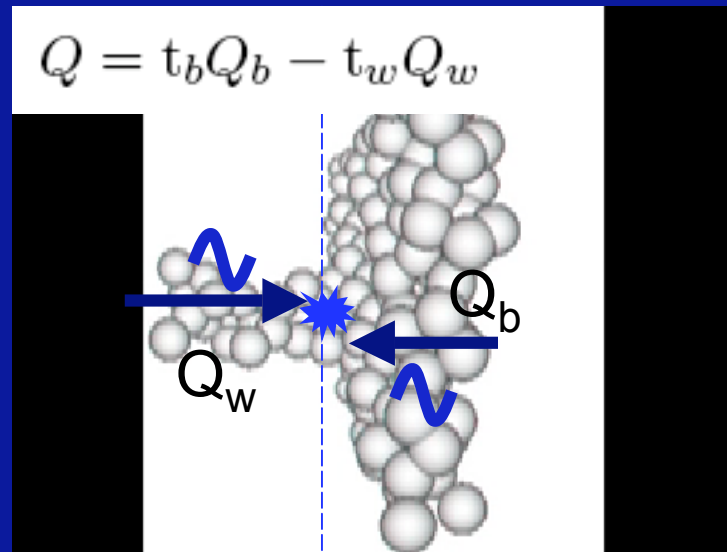


$$\kappa = \frac{k_B^2 \pi^2}{3h} T N_{\alpha} + \frac{k_B^2}{h} T \sum_{\alpha'}^{N_{\alpha'}} \left[\frac{\pi^2}{3} + f(x_0) + \frac{x_0^2 e^{x_0}}{e^{x_0} - 1} \right]$$

Quantum of Conductance

- Conductance of **1 phonon BRANCH** /NOT/ 1 Quantum
- Temperature Dependent
- Predominant CONDUCTANCE?

CONTACT CONDUCTANCE



$$Q = G \cdot (T_w - T_b)$$

$$G = \frac{1}{2} \{t_b C_b v_b + t_w C_w v_w\}$$

$$t_w = \frac{C_b v_b}{C_b v_b + C_w v_w}$$

Diffuse Transmission: Phonons lose memory at interface
 => **MAXIMUM FLUX**

$$\frac{1}{G} = \frac{1}{C_w v_w} + \frac{1}{C_b v_b}$$

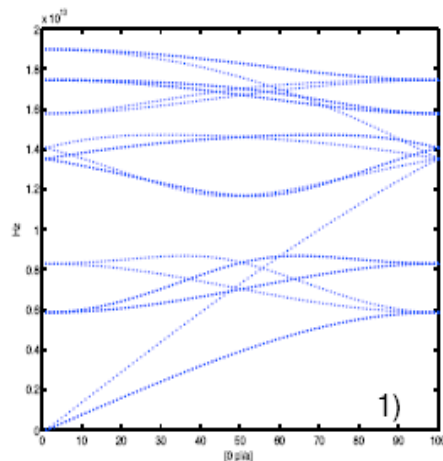
$$C_b v_b = \frac{\partial Q_b}{\partial T} = \pi g_{3D} \sum_m \int_{\|\mathbf{k}\|} \hbar \omega_{m,\mathbf{k}} \frac{\partial \tilde{n}}{\partial T} v_{m,\mathbf{k}} k^2 d\|\mathbf{k}\|$$

$$C_w v_w = \frac{\partial Q_w}{\partial T} = g_{1D} \sum_m \int_{\|\mathbf{k}\|} \hbar \omega_{m,\mathbf{k}} \frac{\partial \tilde{n}}{\partial T} v_{m,\mathbf{k}} d\|\mathbf{k}\|$$

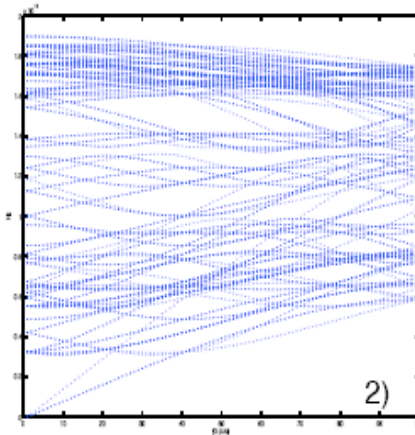
DISPERSION CURVES

Lattice dynamics

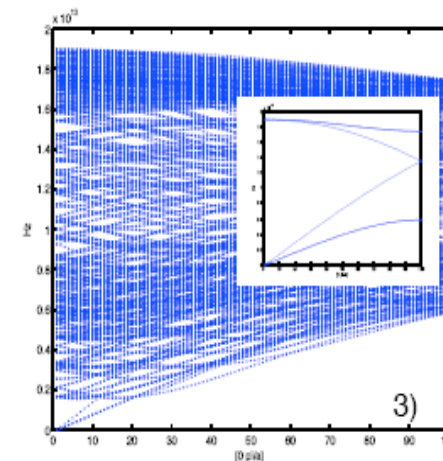
2x2cells in section



5x5



10x10



Motion Equation

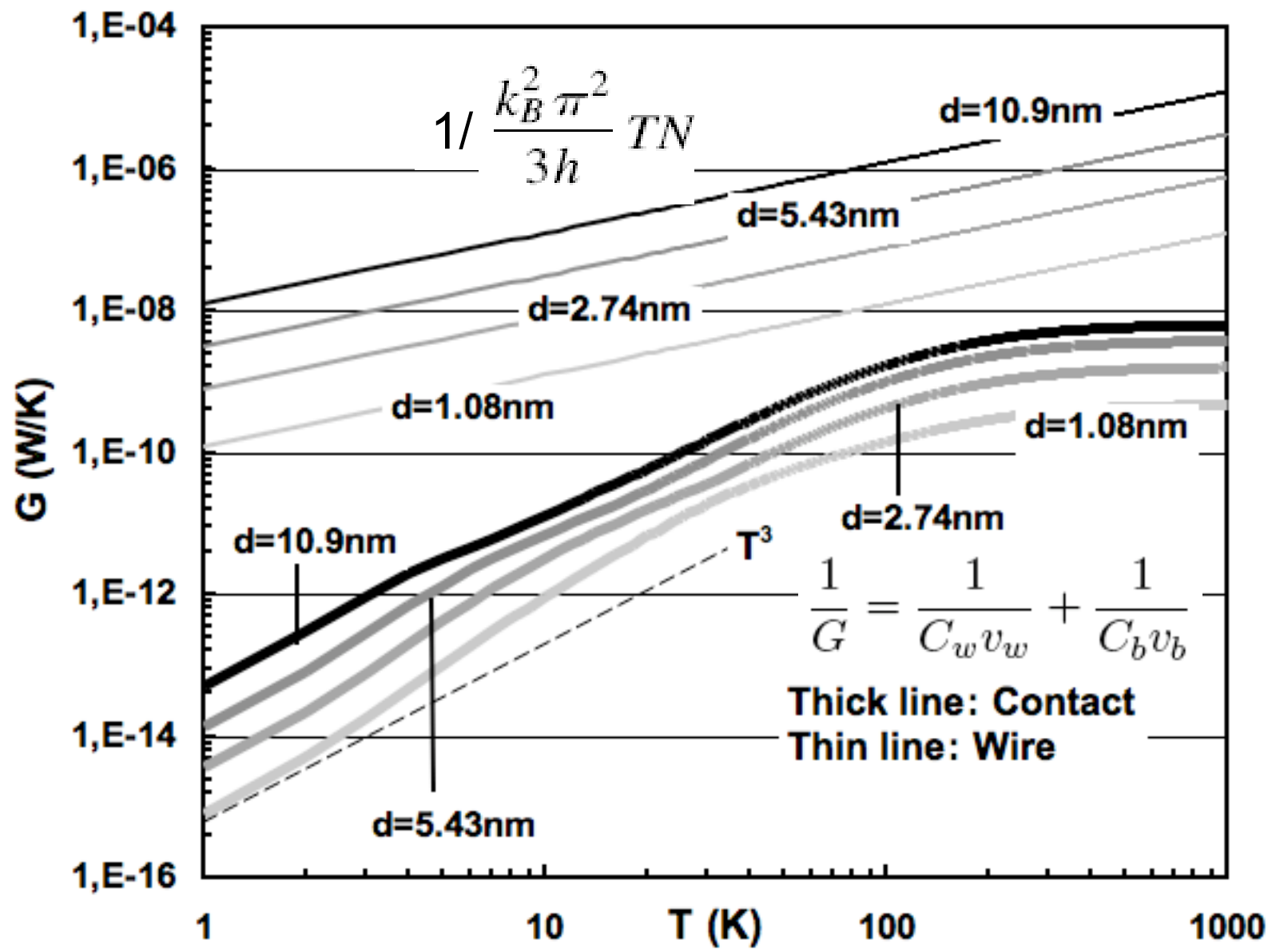
$$m_k \ddot{u}_{k\mu} = \sum_{l'k',\nu} \phi_{l'k',l'k',\nu} u_{l'k',\nu}$$

Solution:

Plane Monochromatic Wave

$$\omega^2(\mathbf{q}) e_{k\mu} = \sum_{k',\nu} D_{\mu\nu}(\mathbf{q}_{kk'}) e_{k'\nu}$$

WIRE vs CONTACT



Submitted PRB

CONTACT RESISTANCE PREDOMINANT 2 ODM

CONCLUSION Nanowire and Contact

Heat flux in of SUB-10nm NANOWIRES
is dominated
by *CONTACT RESISTANCE*

QUANTUM THERMAL CONDUCTANCE
CAN NOT BE *MEASURED* IN THOSE WIRES

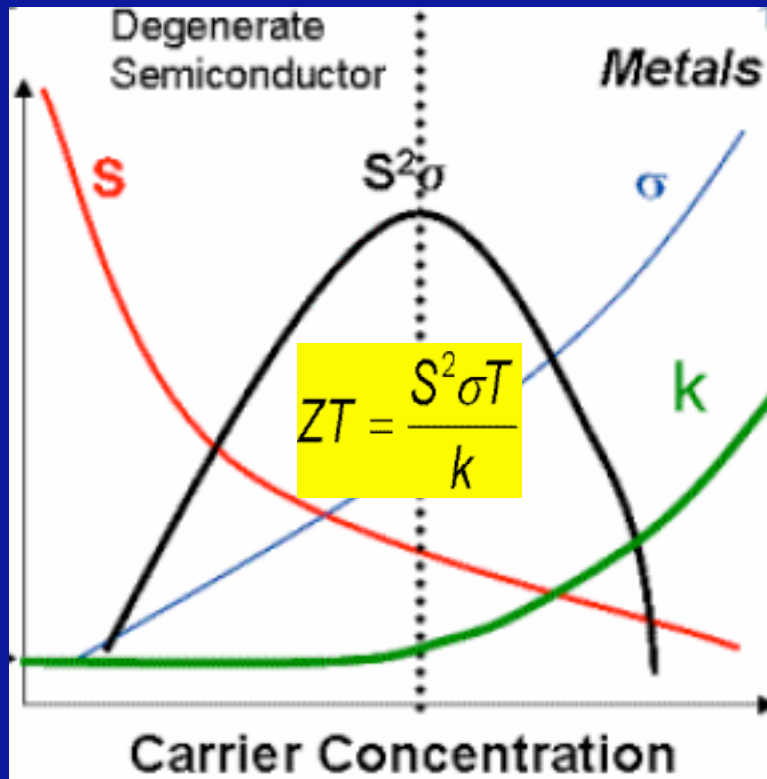
2 to 5 Orders of Magnitude Difference between Resistances

Thermal resistance of wave effects much larger
than Fourier and *Wexler* Resistances

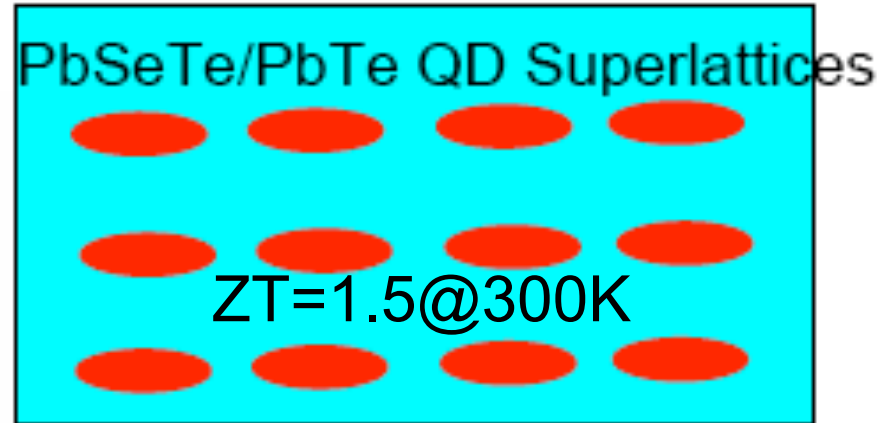
Might be interesting for thermoelectrics or insulating materials

The shape of the contact has significant impact

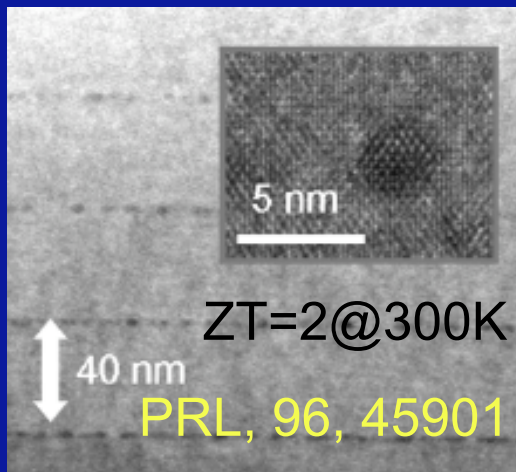
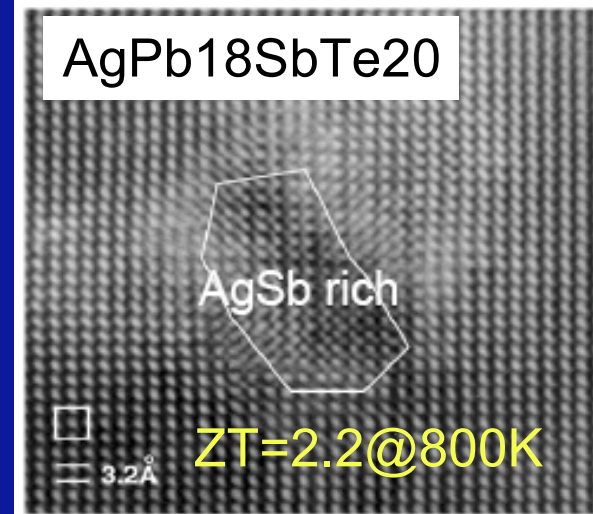
2- 3D Atomic Scale Phononic Crystals



Harman et al., *Science* **297**, 2229

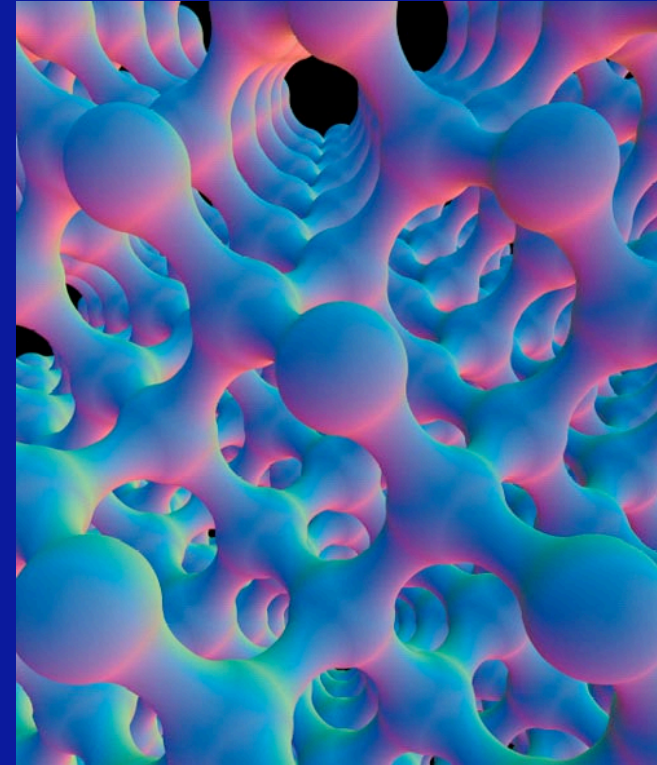
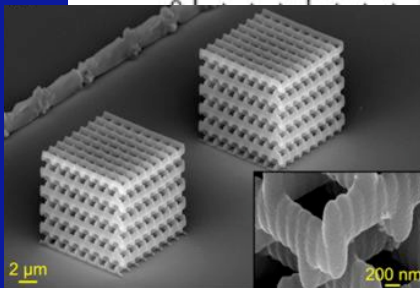
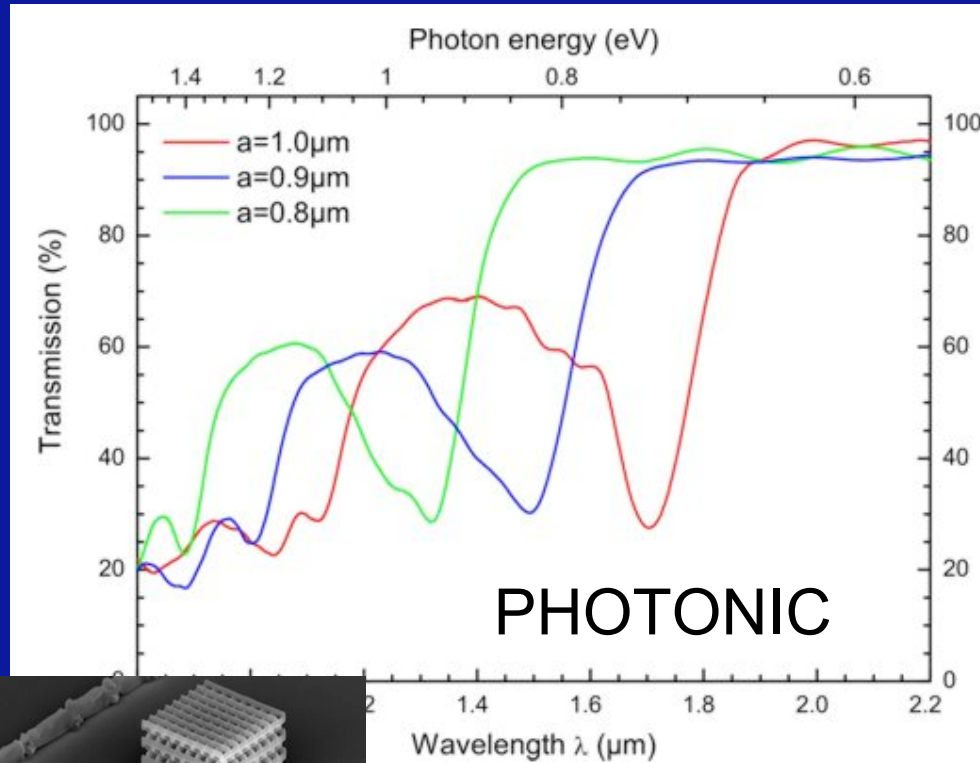


Hsu et al., *Science* **303**, 818 (2004)



THERMAL CONDUCTIVITY

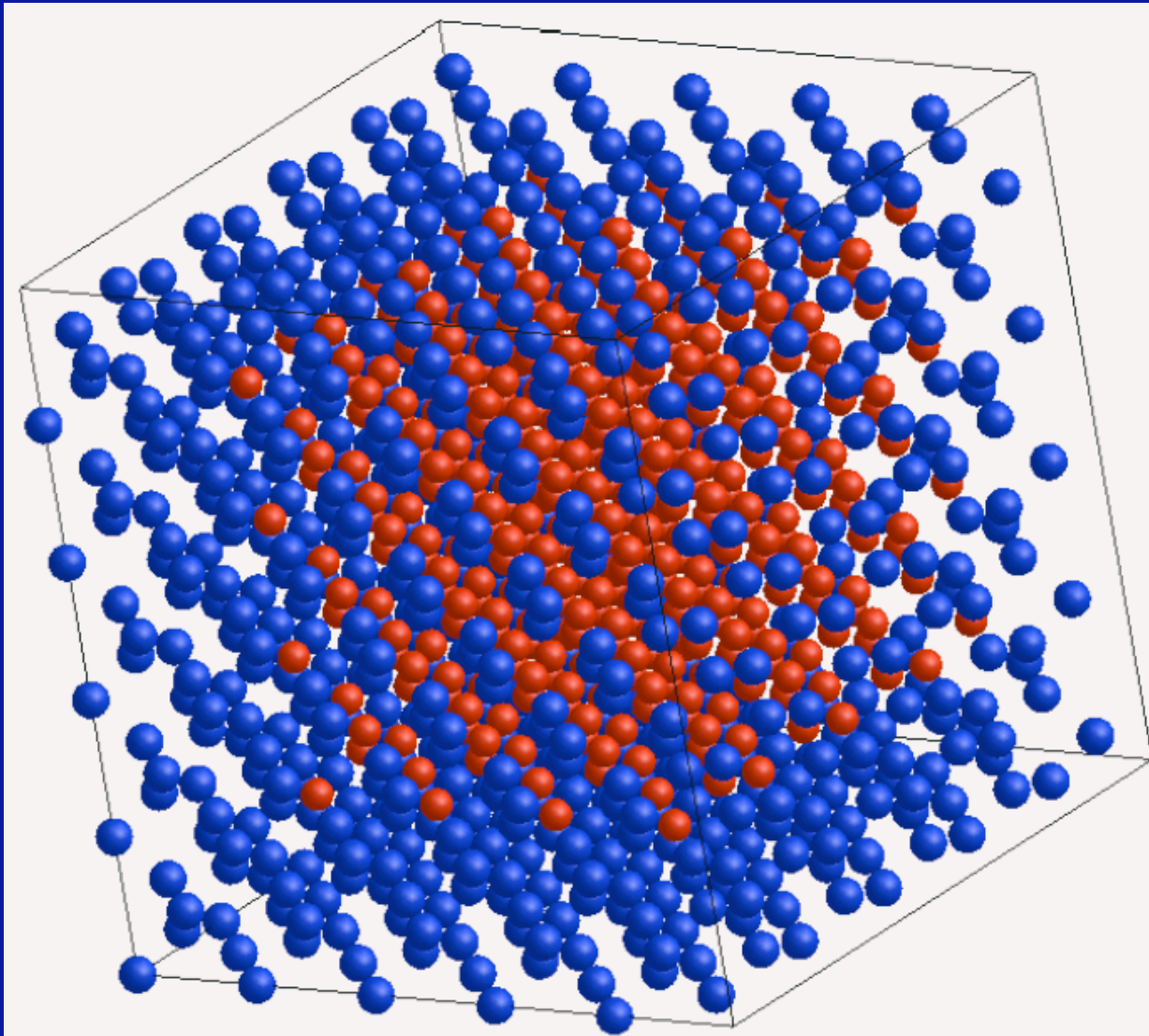
3D-PHONONIC CRYSTALS



3D PHONONIC CRYSTALS
TO GENERATE GAPS
IN PHONON SPECTRA?

THERMAL PHONONS $\lambda \approx 1 \text{ nm} \Rightarrow$ ATOMIC SCALE PERIOD

STRUCTURE



Ge in Si

5 Lattice Constants

Cubic NanoParticles

SW potential

Periodic Boundaries

Lattice Dynamics

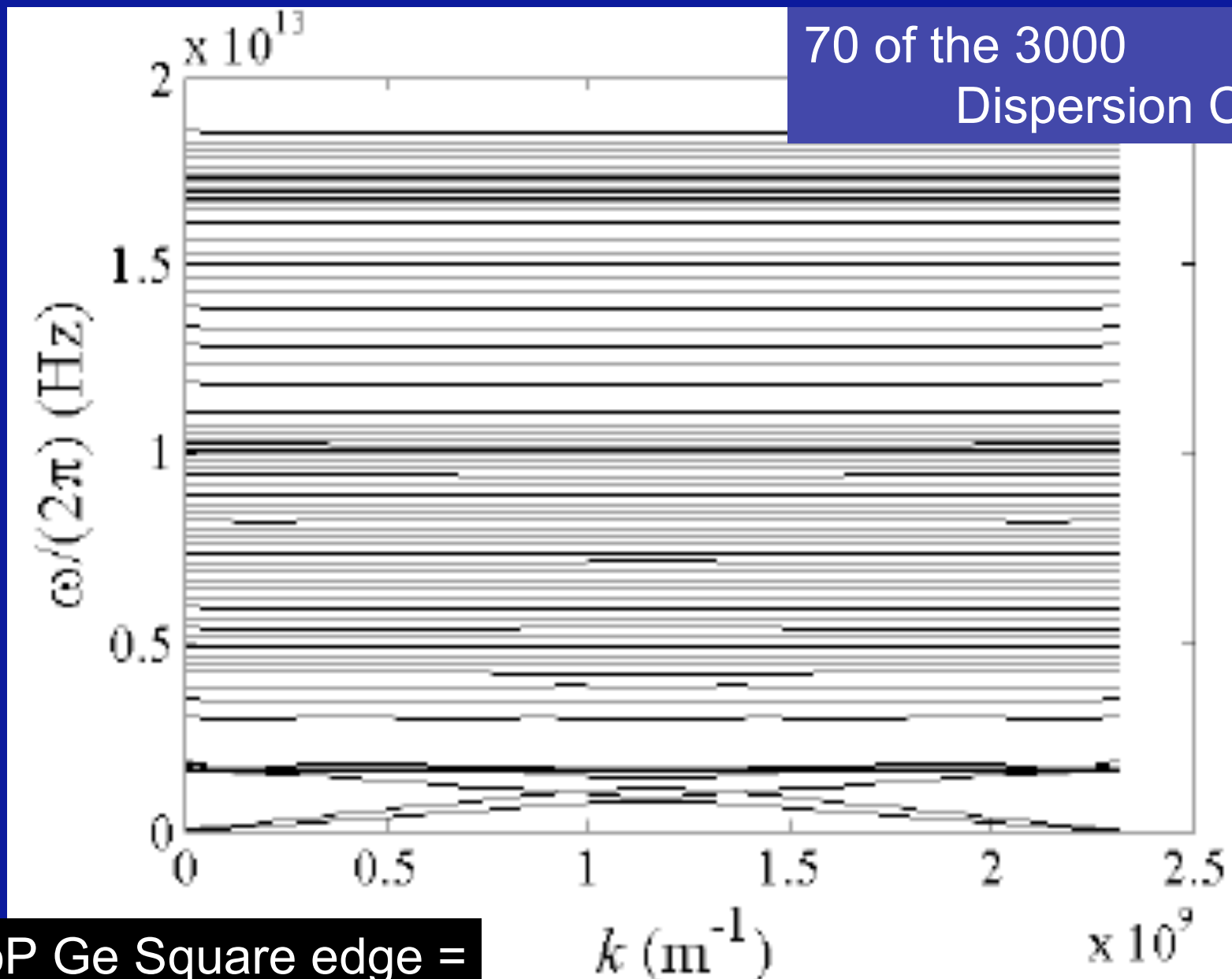
THERMAL CONDUCTIVITY

$$\lambda = \frac{1}{3} \frac{\hbar}{2\pi^2} \sum_m \int_0^{2\pi/d} \Lambda_{k,m} k^2 \omega_{k,m} \frac{\partial n_{k,m}^{(0)}}{\partial T} |v_{k,m}| dk$$

$$\lambda \approx \frac{1}{3} \langle \Lambda \rangle \langle C_p v \rangle$$

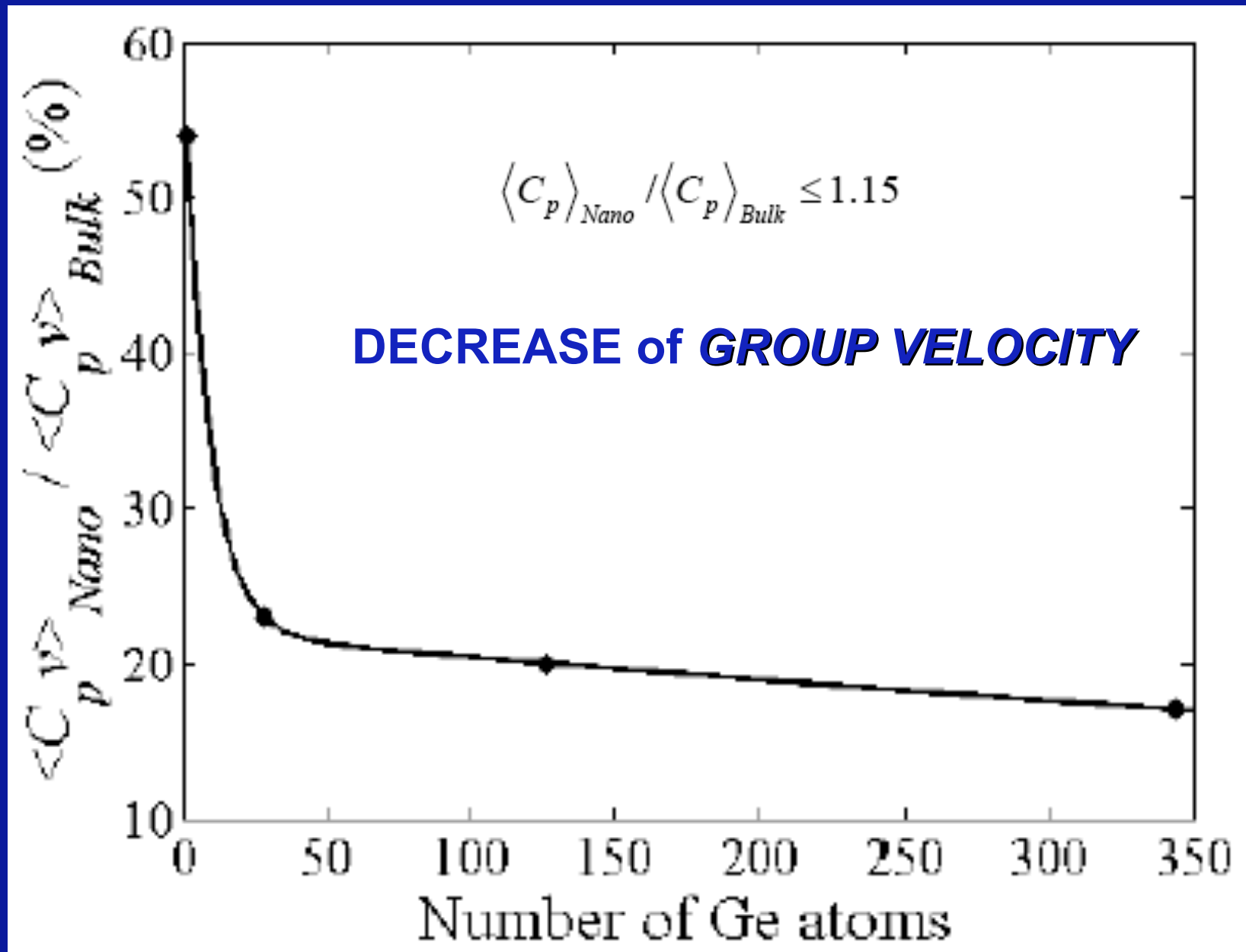
$$\langle C_p v \rangle = \frac{\hbar}{2\pi^2} \sum_{m=1}^{N_m} \int_0^{2\pi/d} k^2 \omega_{k,m} \frac{\partial n_{k,m}^{(0)}}{\partial T} |v_{k,m}| dk$$
$$\langle \Lambda \rangle = \frac{d}{2\pi N_m} \sum_{m=1}^{N_m} \int_0^{2\pi/d} \Lambda_{k,m} dk,$$

DISPERSION CURVES



NanoP Ge Square edge =
3 Lattice constants

$\langle C_{p.v} \rangle$



Submitted PRB

CONCLUSION 3D Phononic Crystals

NO PHONON GAP but
1 ORDER OF MAGNITUDE Decrease
Of GROUP VELOCITY

PRESUMPTION OF MFP DECREASE
BY 1-2 ORDER OF MAGNITUDE

POSSIBLE CONDUCTIVITY DECREASE by FACTOR 1000 = 0.1 W/mK!

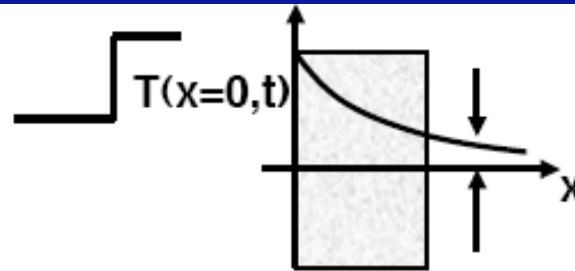
LESS THAN **ALLOY**, GLASS and WATER

3- TRANSIENT NON FOURIER CONDUCTION



Velocity of Heat ?

$$\frac{T(x,t) - T_0}{T_{Bound} - T_0} = \operatorname{erfc}\left(\frac{x}{\sqrt{4\alpha t}}\right)$$

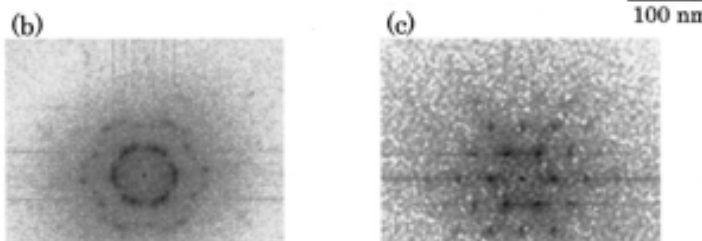
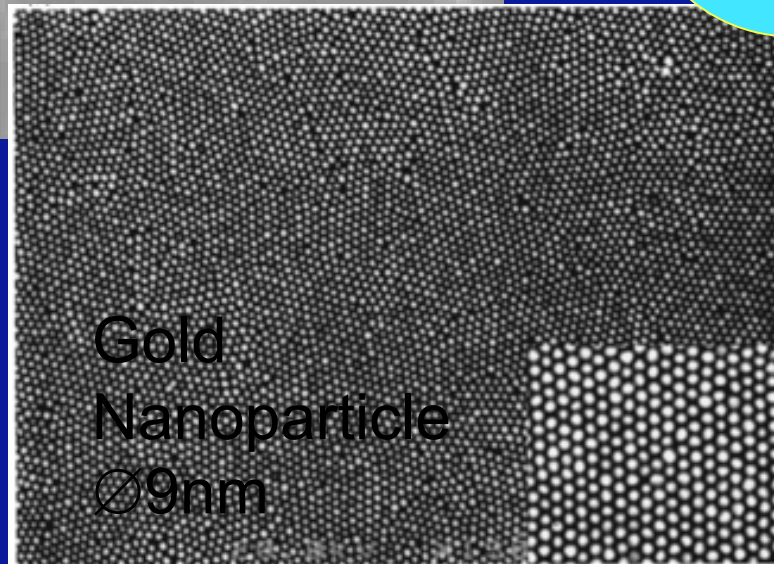
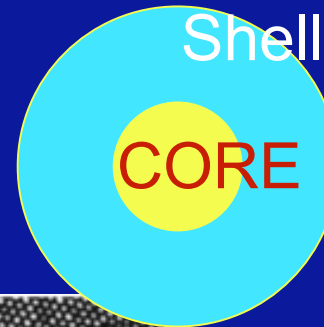


Cattaneo-Vernotte Model: Inertia

$$\tau \cdot dq/dt + q = -\lambda \cdot \operatorname{grad}T$$

NEVER PROVED EXPERIMENTALLY

EXPERIMENT



Pump-Probe Femtolaser

Core Electron absorbing

Relaxing on Core Phonons

Conduction to Shell

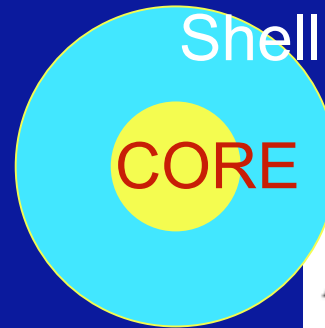
Shell Heat conduction:

Fourier or Non-Fourier

2-STEP+BALLISTIC-DIFFUSIVE MODELS

$$C_e \frac{\partial T_e}{\partial t} = -G(T_e - T_l) + P_{vol}(t)$$

$$V_p C_l \frac{\partial T_l}{\partial t} = V_p G(T_e - T_l) - H(t)$$



$$H(t) = \int_{S_p} \mathbf{q}(\mathbf{r}, t) \cdot \mathbf{n} ds$$

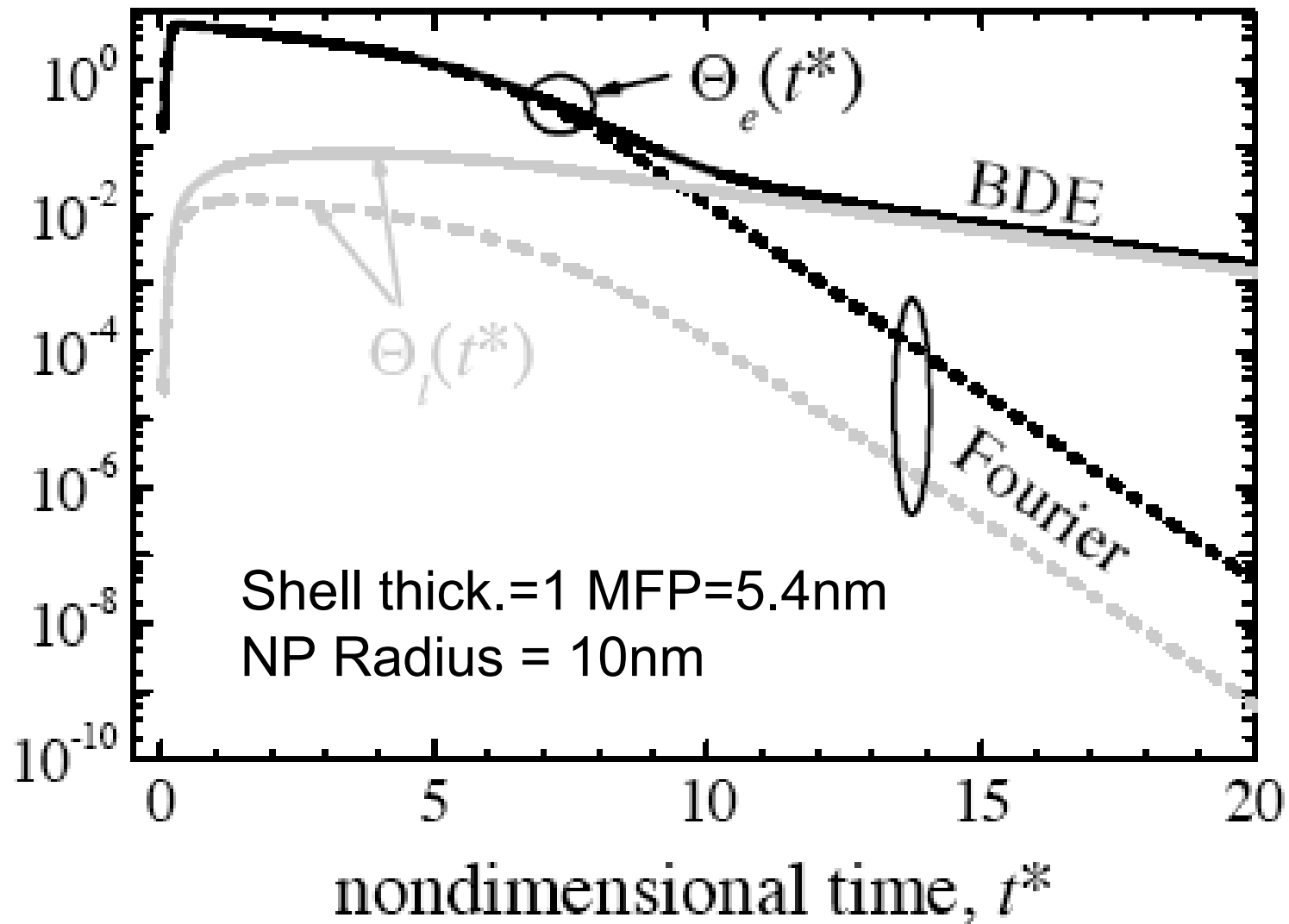
BALLISTIC-DIFFUSIVE EQUATION

$$C \left(\tau \frac{\partial^2 T_m}{\partial t^2} + \frac{\partial T_m}{\partial t} \right) = \nabla(k \nabla T_m) - \nabla \cdot \mathbf{q}_b + \left(\dot{q}_e + \tau \frac{\partial \dot{q}_e}{\partial t} \right),$$

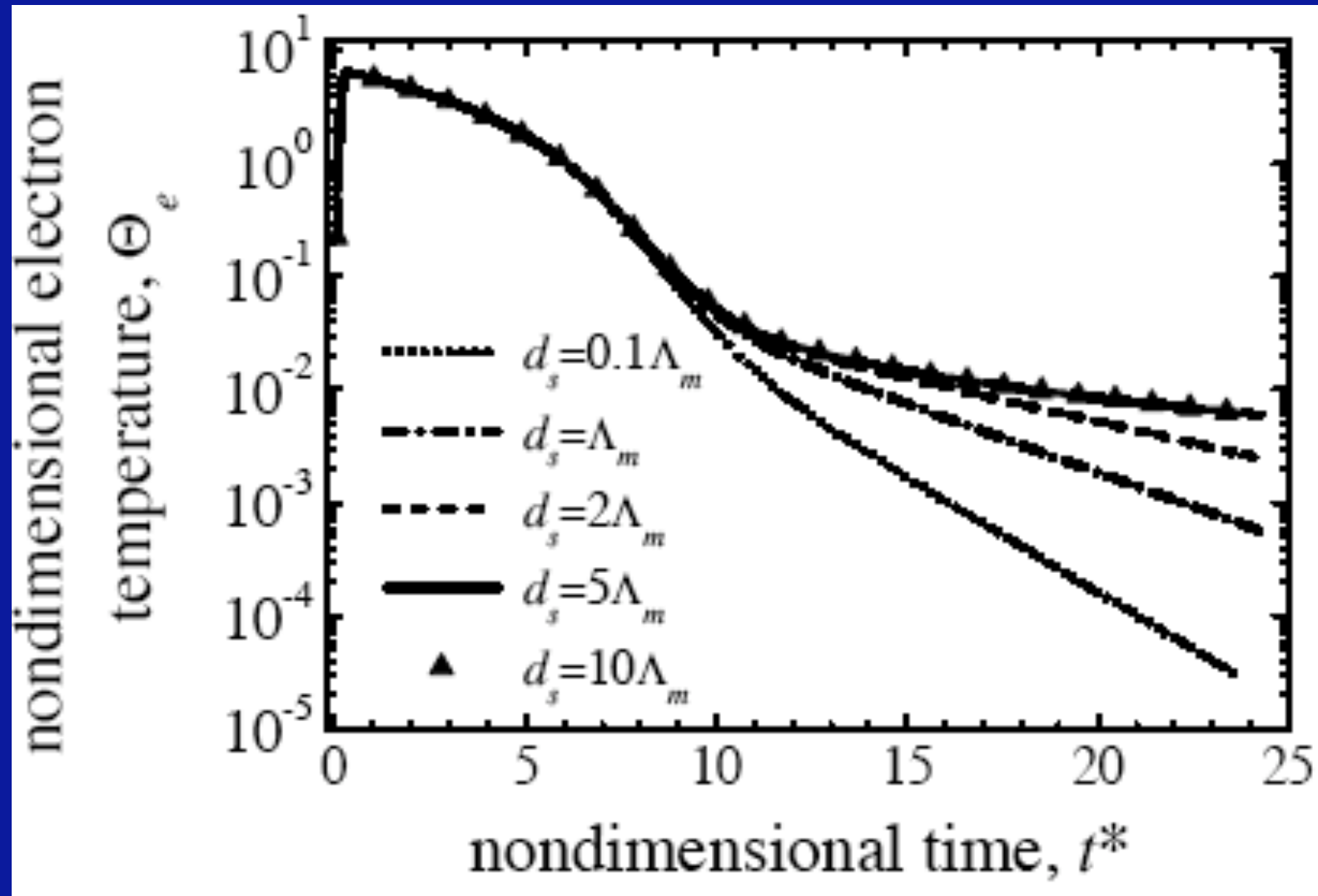
$$\mathbf{q}_b(t, \mathbf{r}) = \frac{1}{4\pi} \int \left[\int |\mathbf{v}| \hbar \omega D(\omega) f_w[t - (s - s_0)/|\mathbf{v}|, \mathbf{r} - (s - s_0)\hat{\Omega}] \exp\left(-\int_{s_0}^s \frac{ds}{|\mathbf{v}| \tau}\right) \cos\theta d\Omega \right] d\omega$$

BDE/FOURIER

non-dimensional temperature



Shell Thickness



STRONG DEPENDENCE TO SHELL THICKNESS

CONCLUSION on TRANSIENT NON FOURIER

EXPERIMENTAL ACCESS TO TRANSIENT NON-FOURIER HEAT CONDUCTION

By USING CORE-SHELL STRUCTURES

WAS NEVER SHOWN AT AMBIENT TEMPERATURE

NEED A SHELL WITH LONG MFP

Collaborators

Post-Doc Students



M. Rachidi



J.-N. Gillet



E. Rousseau

Ph.D. Students



Y. Chalopin



P.-O. Chapuis



C. Bera

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N. Mingo, CEA-LITEN
M. Plissonnier, CEA-LITEN

107 | Topics in Applied Physics

THANK YOU !

S. Volz (Ed.)

**Microscale
and Nanoscale
Heat Transfer**

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Stefan Dilhaire, Séverine Gomez,
Nathalie Trannoy, and Gilles Tessier

 Springer