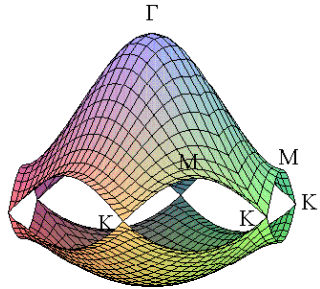


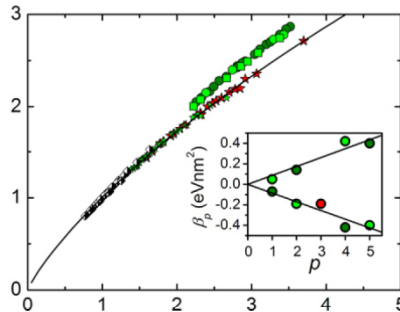
Resonance Raman Studies of Exciton Behavior In Single-Walled Carbon Nanotubes

Stephen K. Doorn
Los Alamos National Laboratory

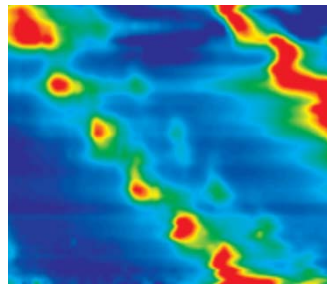
Presentation Overview



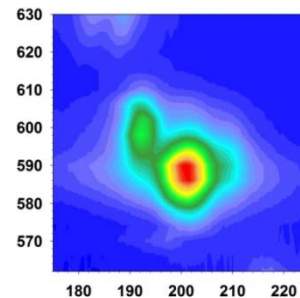
1. SWNT Background



2. Excitonic Behavior of Semiconducting Nanotubes



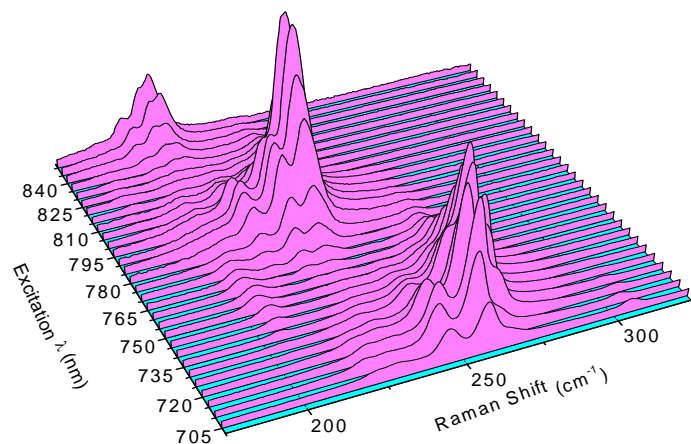
3. Electronic Behavior of Metallic Nanotubes



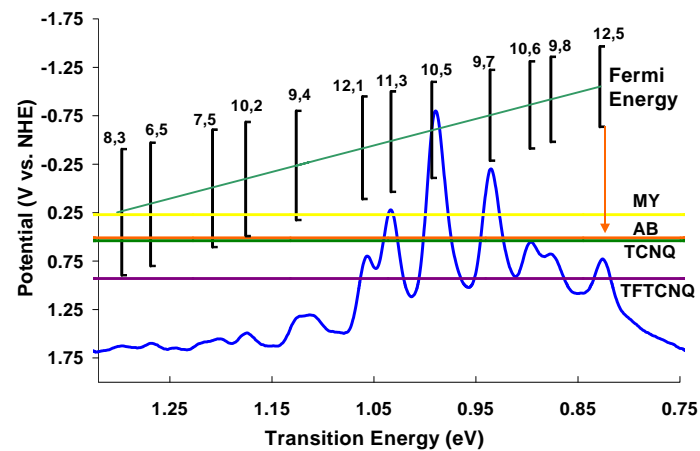
4. Raman of Enriched Metallic Samples

LANL Nanotube Effort

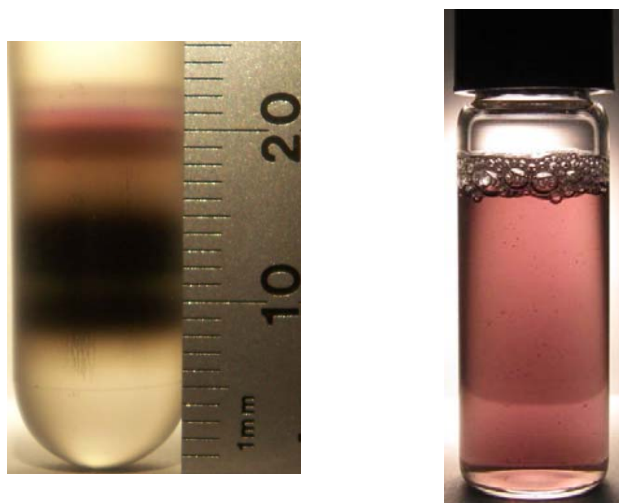
Fundamental Spectroscopy



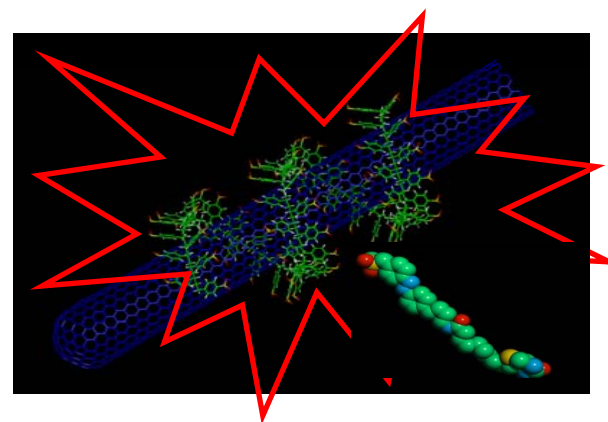
Nanotube Redox Chemistry



Nanotube Separations

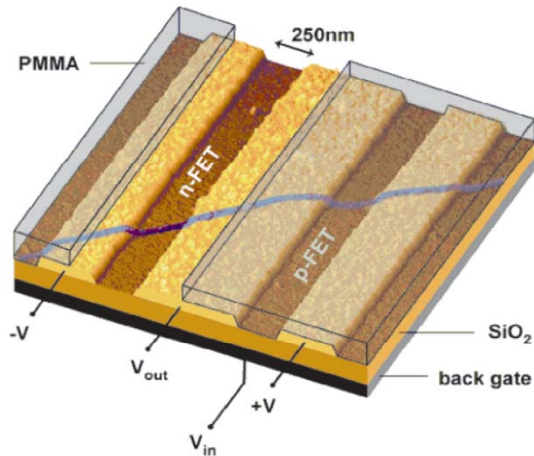


Nanotube-Based Sensing



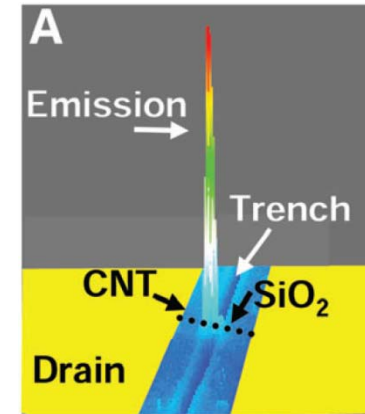
Example Applications

Nanoscale Electronics



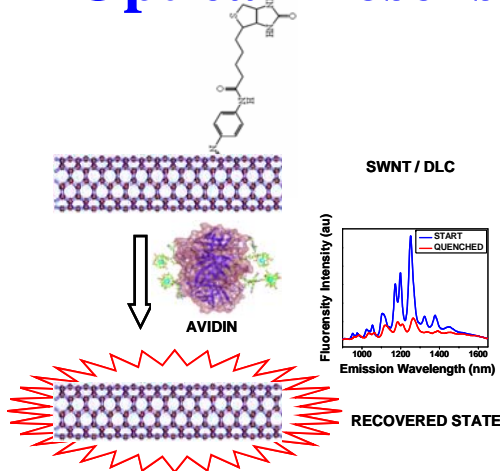
Derycke, et. al., *Nano Lett.* **1**, 453 (2001).

Opto-Electronics



Chen, et. al., *Science.* **310**, 1171 (2005).

Optical Biosensing

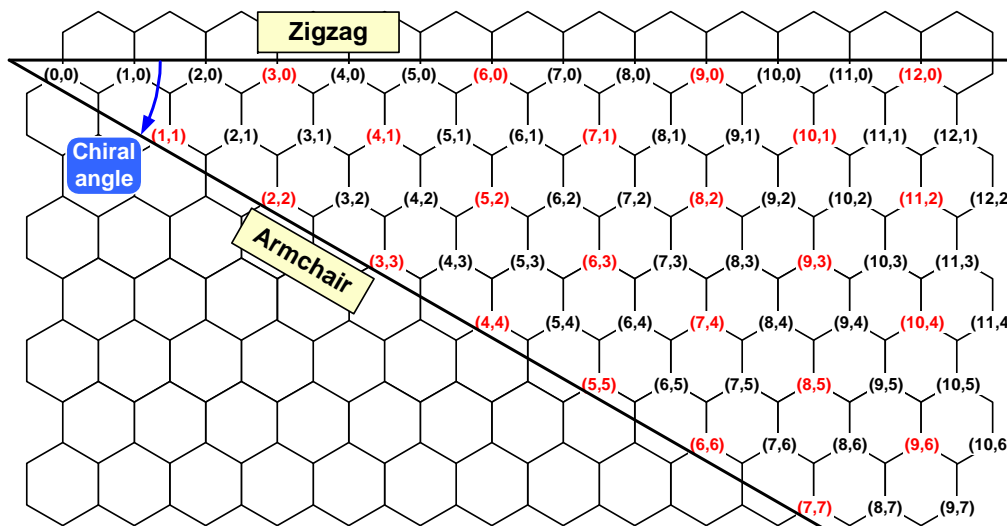


Satishkumar, et. al., *Nature Nanotech.* **2**, 560 (2007).

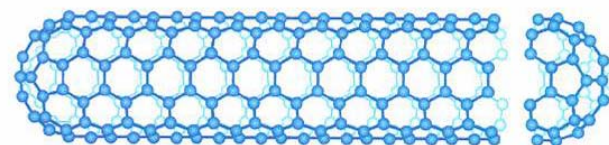
Fundamental Issues:

- what is the electronic structure?
- what is the nature of the optically excited state?
- how do electrons/excitons couple to phonon structure?

Construction of Nanotubes from a Graphene Sheet

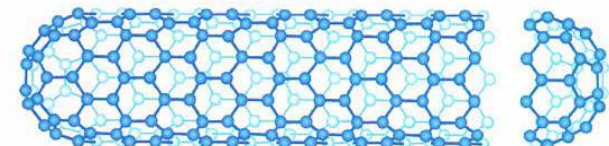


Roll-up Indices (n,m)
Reveal Structural and
Electronic Character



(n,m) = (5,5)

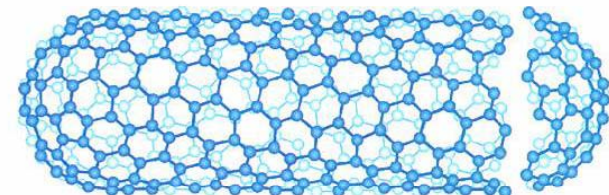
$(n-m, \text{mod}3) = 0$ metallic



(n,m) = (9,0)

$(n-m, \text{mod}3) = 1$

semiconductor



(n,m) = (10,5)

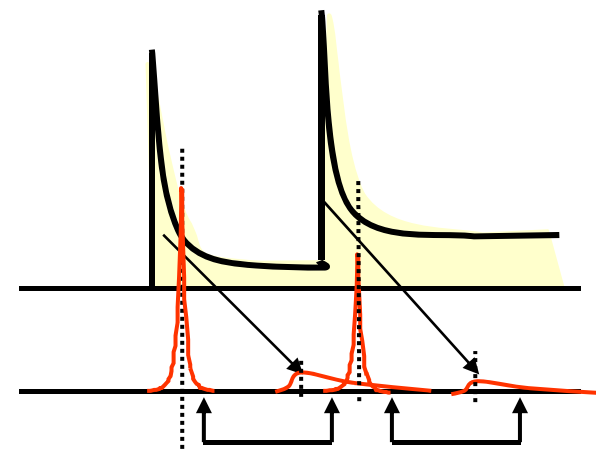
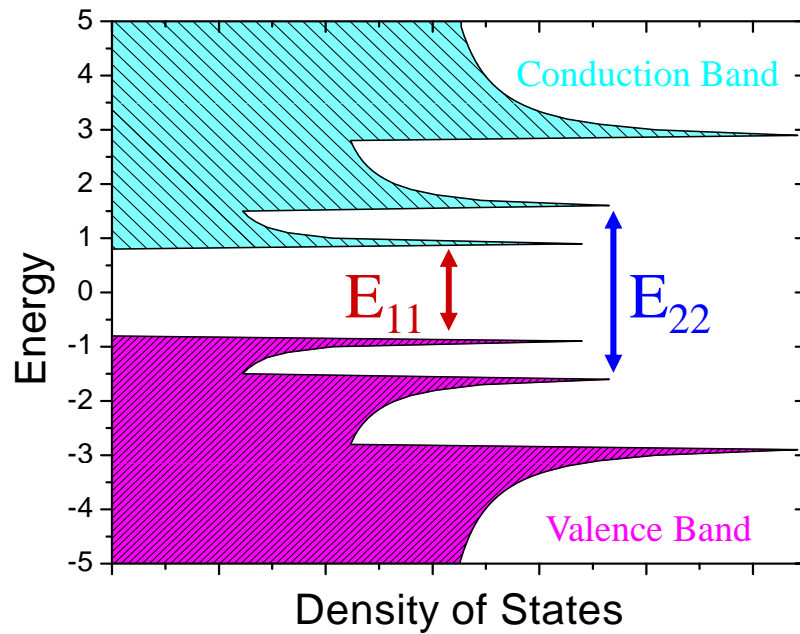
$(n-m, \text{mod}3) = 2$

Nature of Electronic Excited States

Band Model

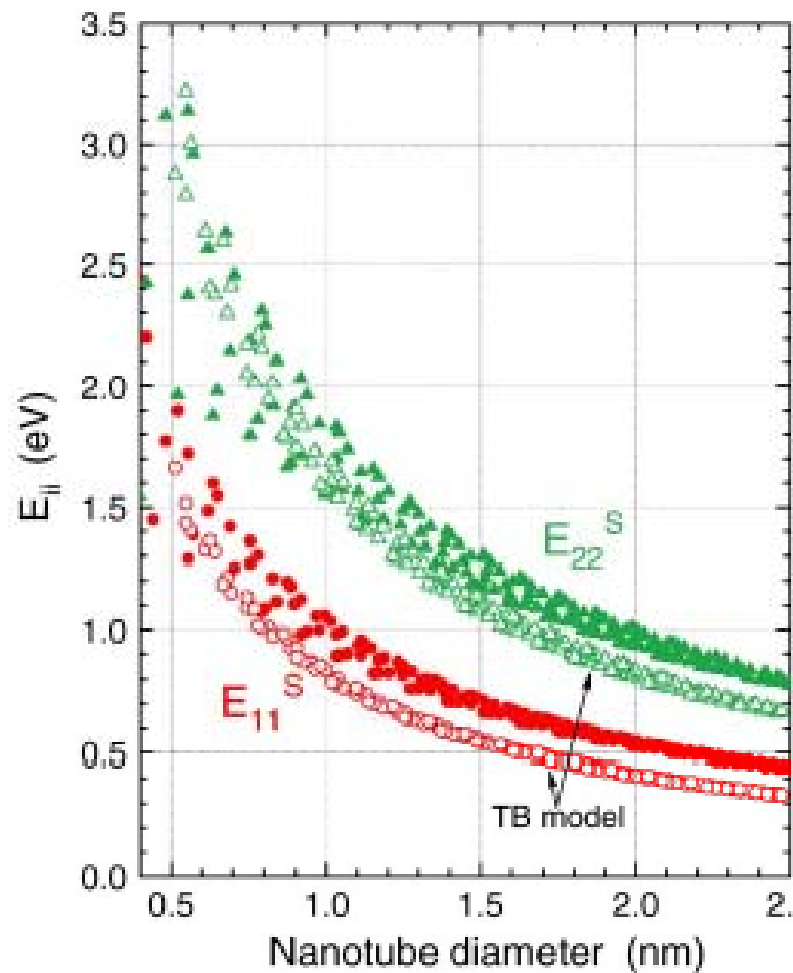
vs.

Excitonic Description

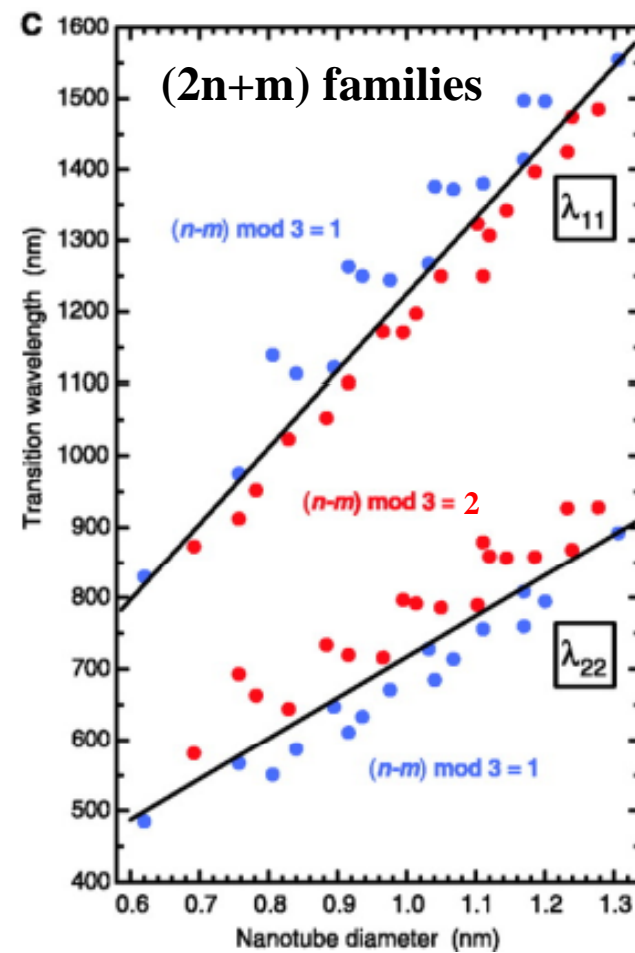


T. Ando J. Phys. Soc Japan
66, 1066 (1997)

Experimental Deviation from Tight-Binding Description



R.B. Weisman, S.M. Bachilo, D.Tsyboulski,
Appl. Phys. A, **78**, 1111 (2004).

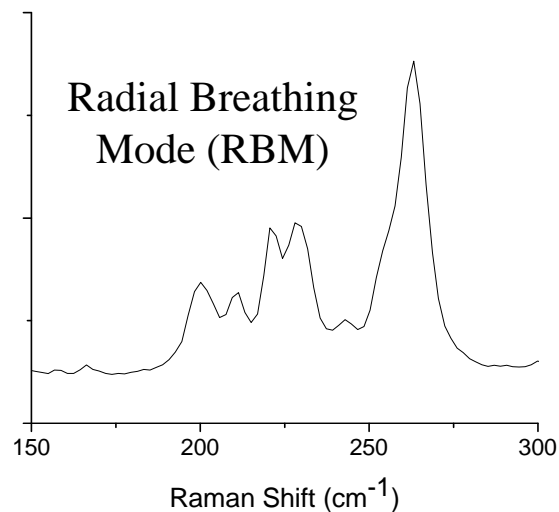


S.M. Bachilo et. al., *Science*,
298, 2361 (2002).

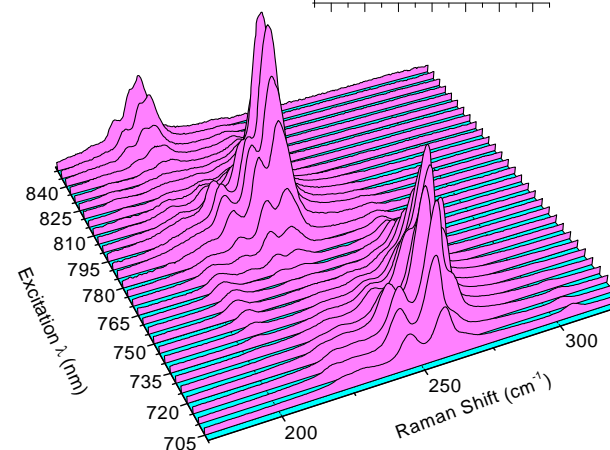
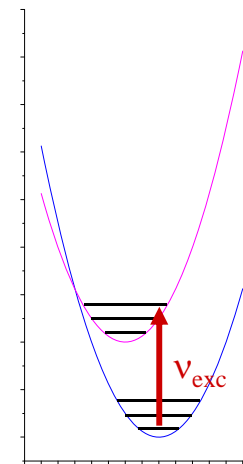
Raman Spectroscopy Is....

Vibrational:

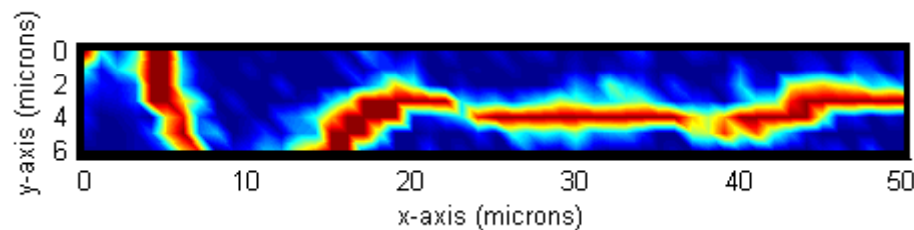
$$\omega_{\text{RBM}} = \frac{C_1}{d_{\text{SWNT}}} + C_2$$



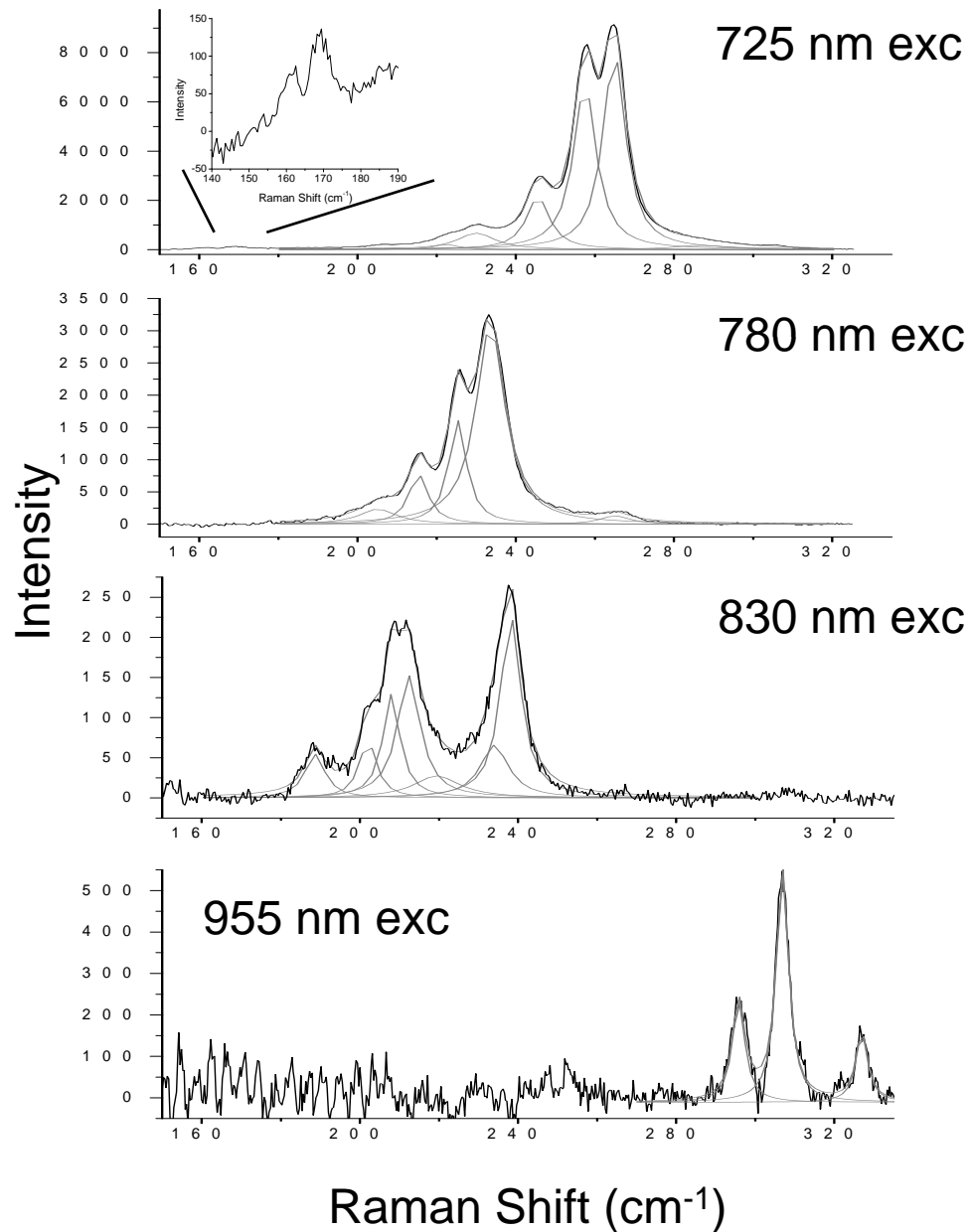
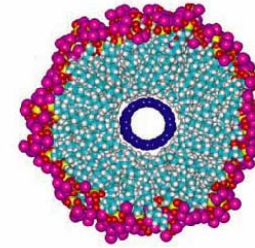
Electronic:



Imaging:

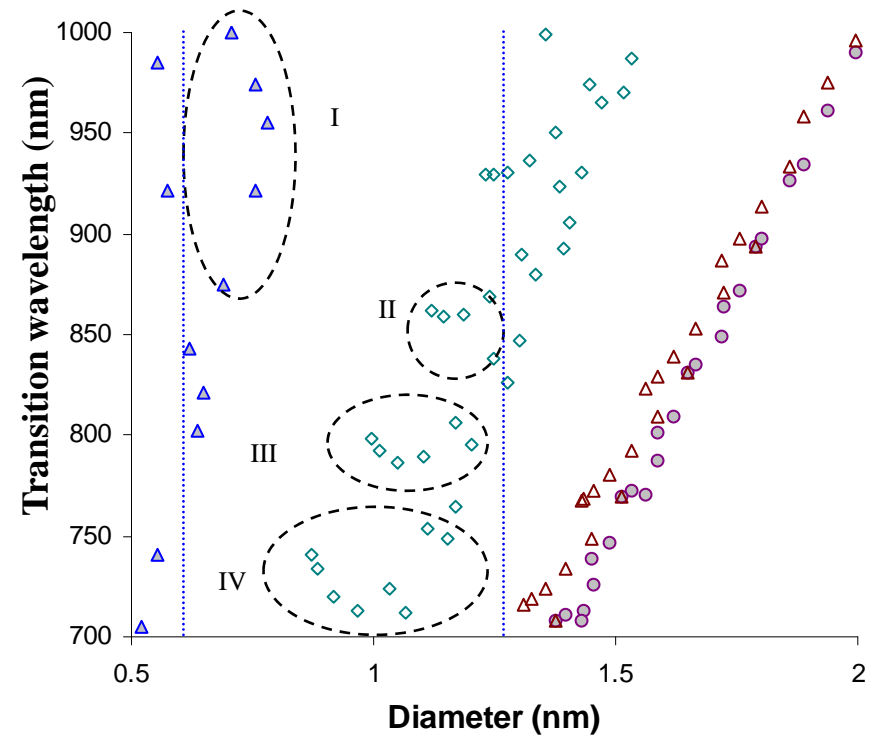
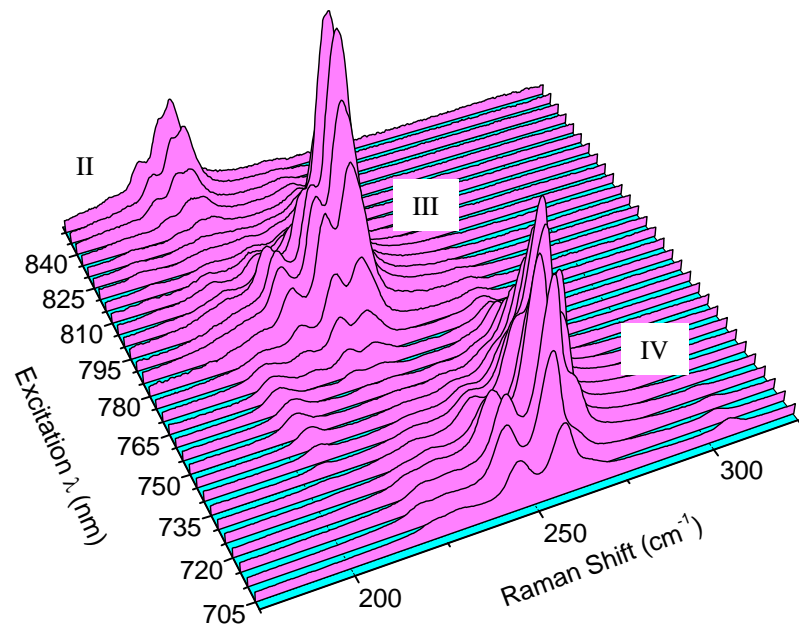
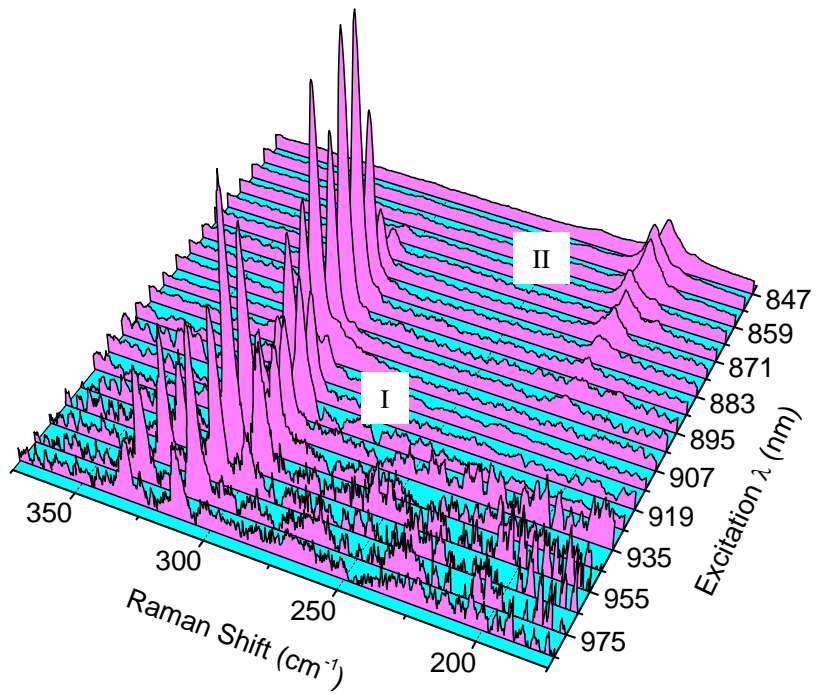


Raman on SDS Solubilized HiPco NTs



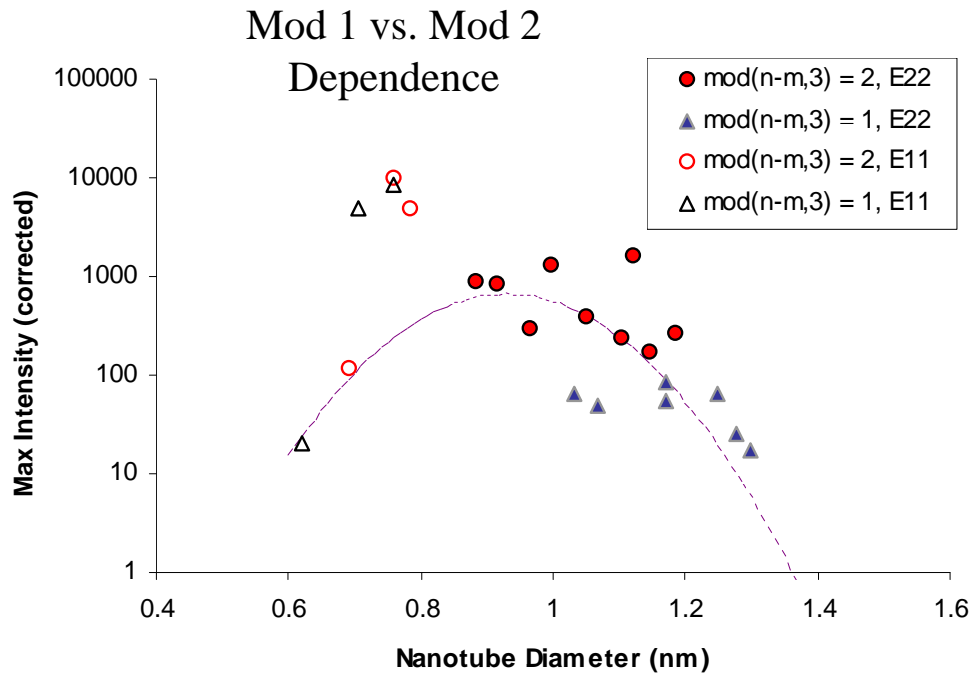
Individualized HiPco nanotubes
with observed diameters ranging
from 0.6 to 1.6 nm.

Variation in $E_{v,H}$ with chirality
results in sampling of different
chiral groupings as Raman
excitation is tuned.



S.K. Doorn, et. al., *Appl. Phys. A*, **78**, 1147 (2004).

Chirality Dependence of Raman Scattering Intensity



S.K. Doorn, et. al., *Appl. Phys. A*, **78**, 1147 (2004).

Chiral dependence reverses on going from E_{11} to E_{22} excitation.

Explains weakness of (6,4) and (8,4) chiralities with E_{11} excitation. Are strong with E_{22} .

$$\alpha = \sum_{i,j} \frac{M_{ee}^{g,i} M_{ee}^{i,g} M_{ep}^{i,j}}{(E_{\text{laser}} - E - i\Gamma_r) (E_S - E - i\Gamma_r)}$$

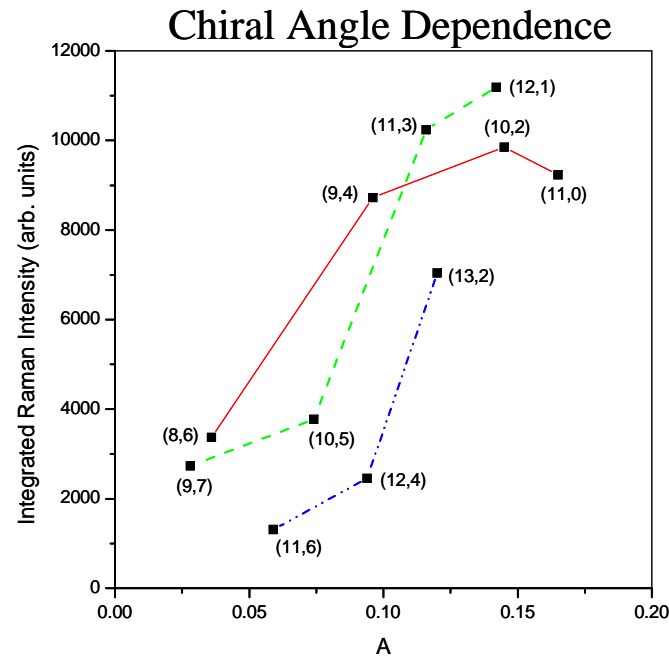
Originates in exciton-phonon coupling.

$$\hat{V}_{\text{exc-ph}}^{\text{RBM}} \propto \left[\frac{\partial \gamma_0(\tau_l)}{\partial \tau_l} - 3 \delta_0(\tau_l) \frac{\partial \sigma_0(\tau_l)}{\partial \tau_l} \right] \mathbf{X}$$

$$\left\{ \text{sign}(n_e + \nu/3) \cos 3\theta + \frac{5}{4\sqrt{3}} \frac{a}{R} |n_e + \nu/3| \left(1 + \frac{\cos^2 3\theta}{5} \right) \right\}$$

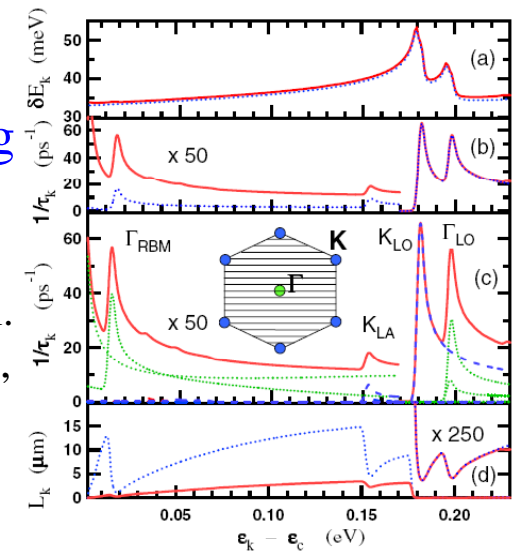
S.V. Goupalov, *Phys. Rev. B*, **71**, 153404 (2005).

S.V. Goupalov, B.C. Satishkumar, S.K. Doorn,
Phys. Rev. B, **73**, 115401 (2006).



Phonon Scattering And Transport

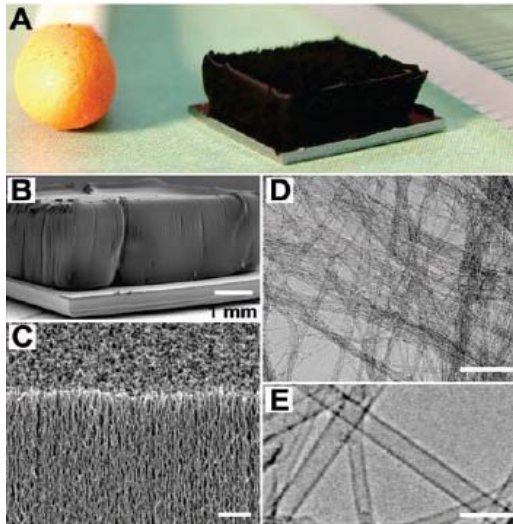
V. Perebeinos, et.al.
Phys. Rev. Lett., **94**,
086802 (2005)



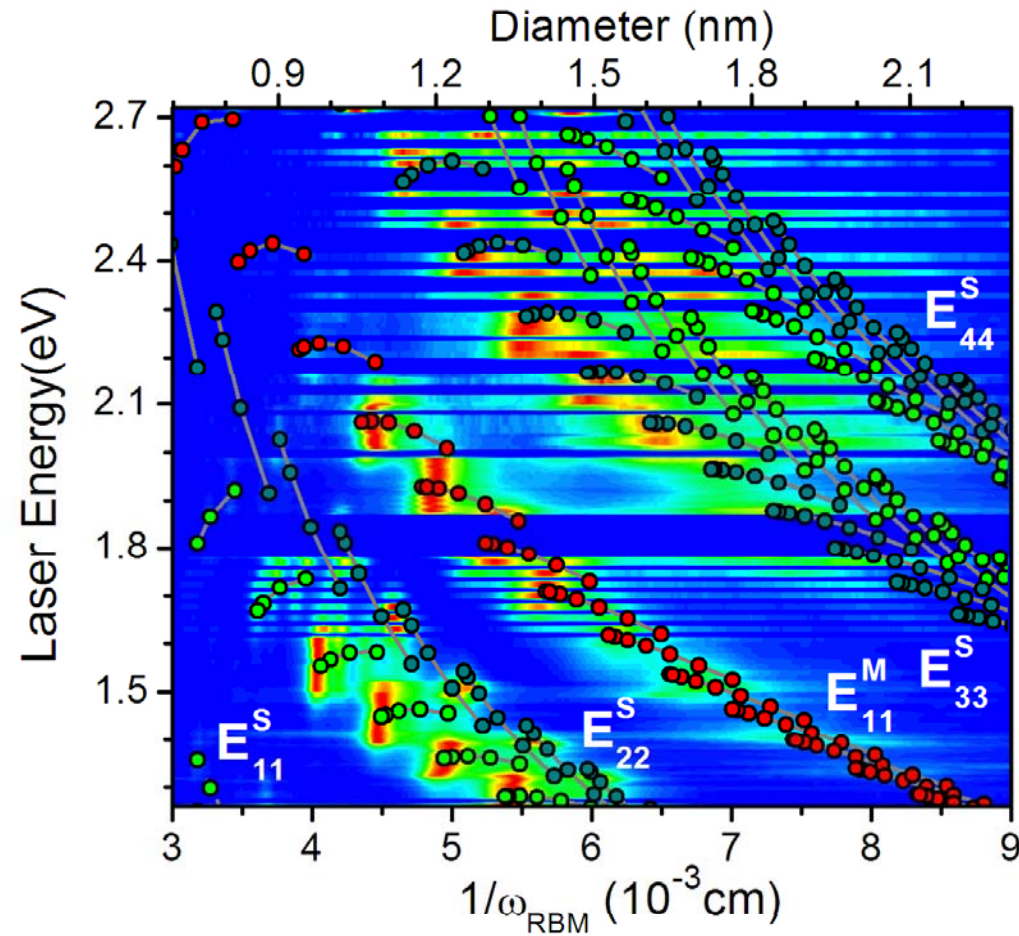
Electronic Structure at High Excitation Energies

E_{11} and E_{22} Excitations Produce Bound Excitons
What About E_{33} and E_{44} ?

Raman on
Vertical Arrays

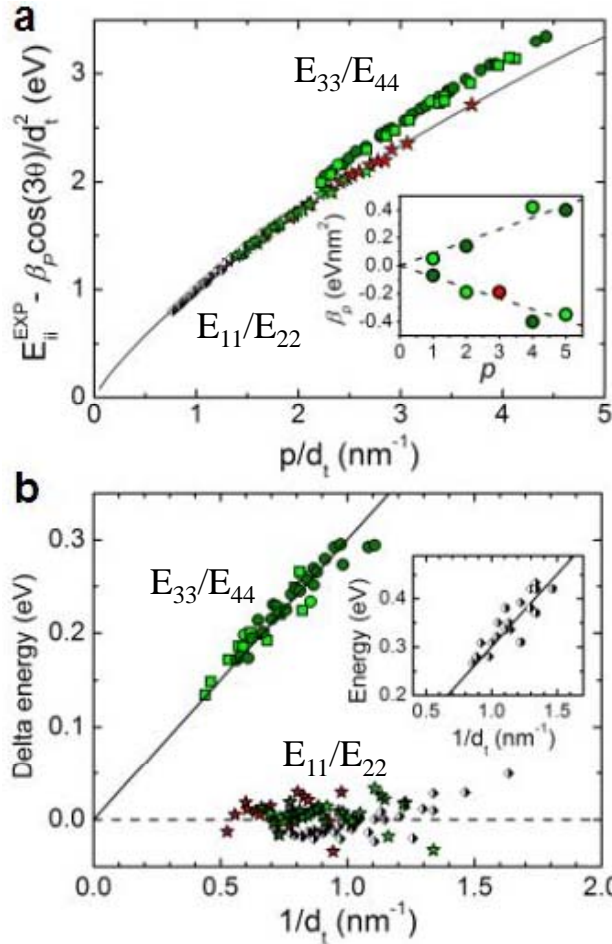


K. Hata et al. Science 306, 1362 (2004)
Murakami et. al. Carbon, 43 2664 (2005)



E33/E44 Energy Behavior Departs From E11/E22

P.T. Araujo, et. al., *Phys. Rev. Lett.*
98, 067401 (2007).



Nonlinear Scaling Analysis Of Transition Energies

$$E_{ii}(p, d_t) - \beta_p \cos 3\theta / d_t^2 = a \frac{p}{d_t} \left[1 + b \log \frac{c}{p/d_t} \right]$$

Chirality dependence Single Particle Energy Coulomb Contribution

For E_{33} , E_{44} :

Add Additional factor of $0.305/d_t$

Implicates change in
excitonic behavior.

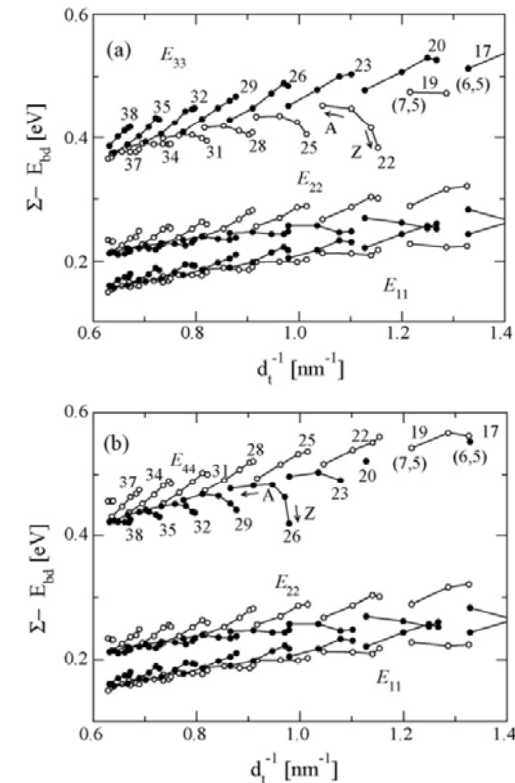
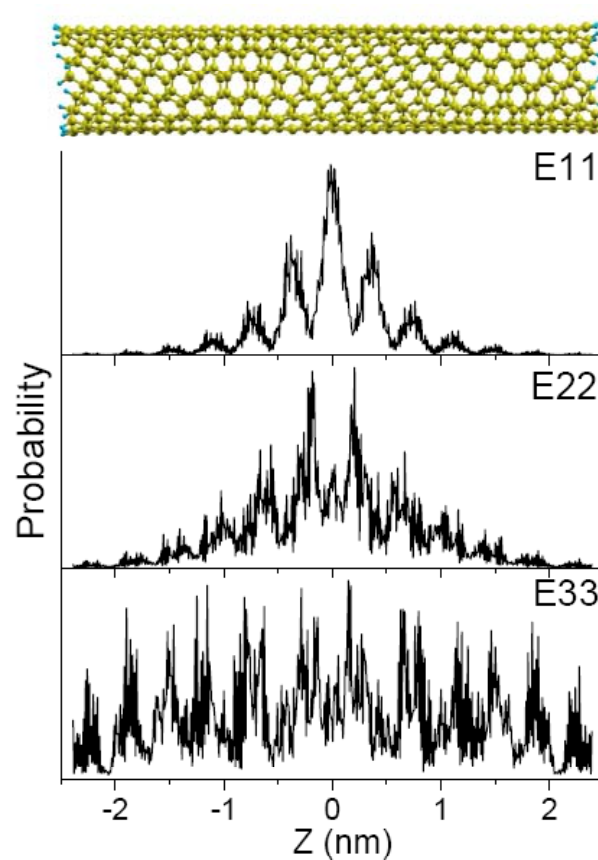
Reflection of Decrease In Binding Energy?

-OR-

Reflection of Assymmetric Increase In Self-Energy?

Quantum chemistry calculation
shows delocalized electron wavefunction

Self-energy increases
faster than binding energy.

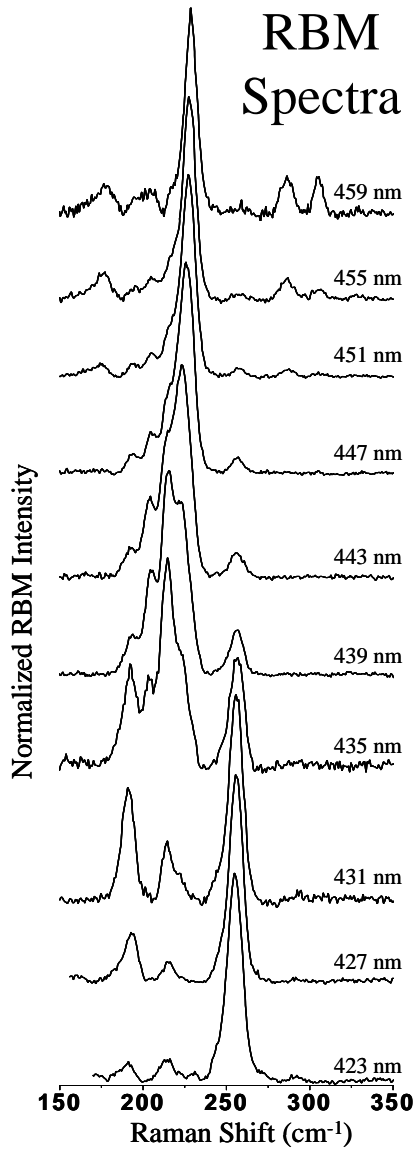


K. Sato et al. Vib. Spect. 45, 89 (2007)

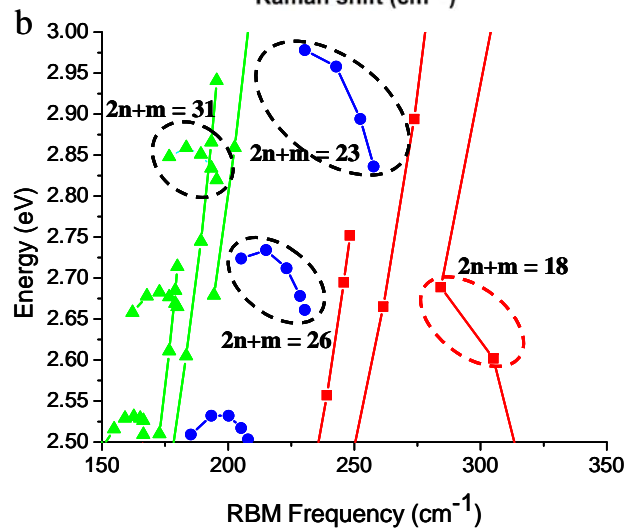
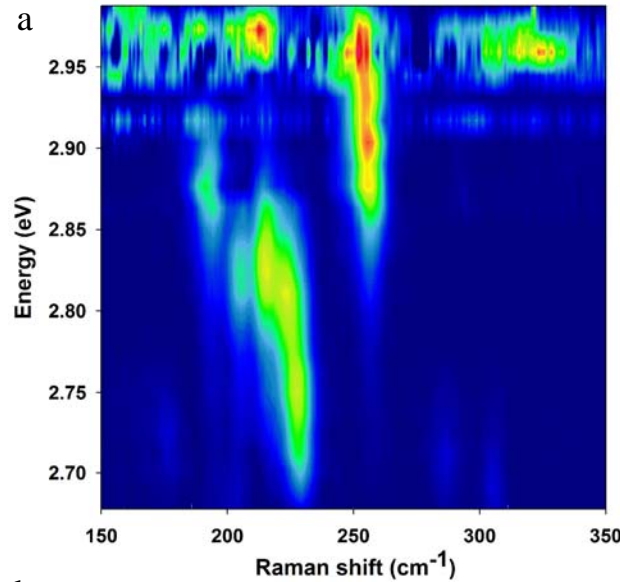
P. T. Araujo et al. PRL 98, 067401 (2007)



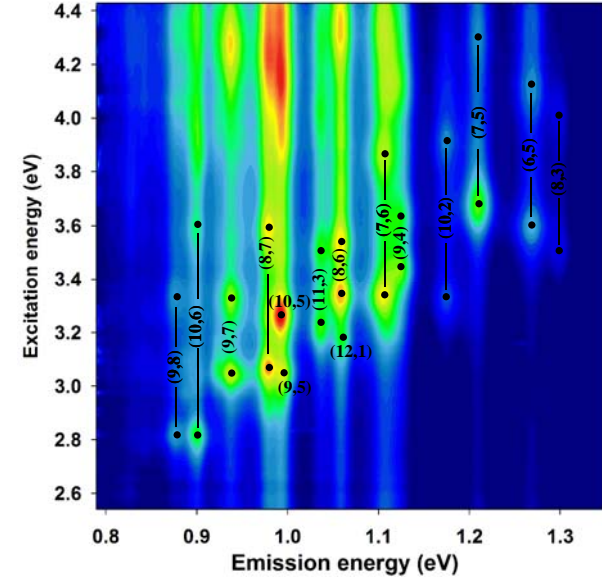
Curvature Effects on E33 and E44 Exciton Energies



Raman Excitation Map



Photoluminescence Excitation Map



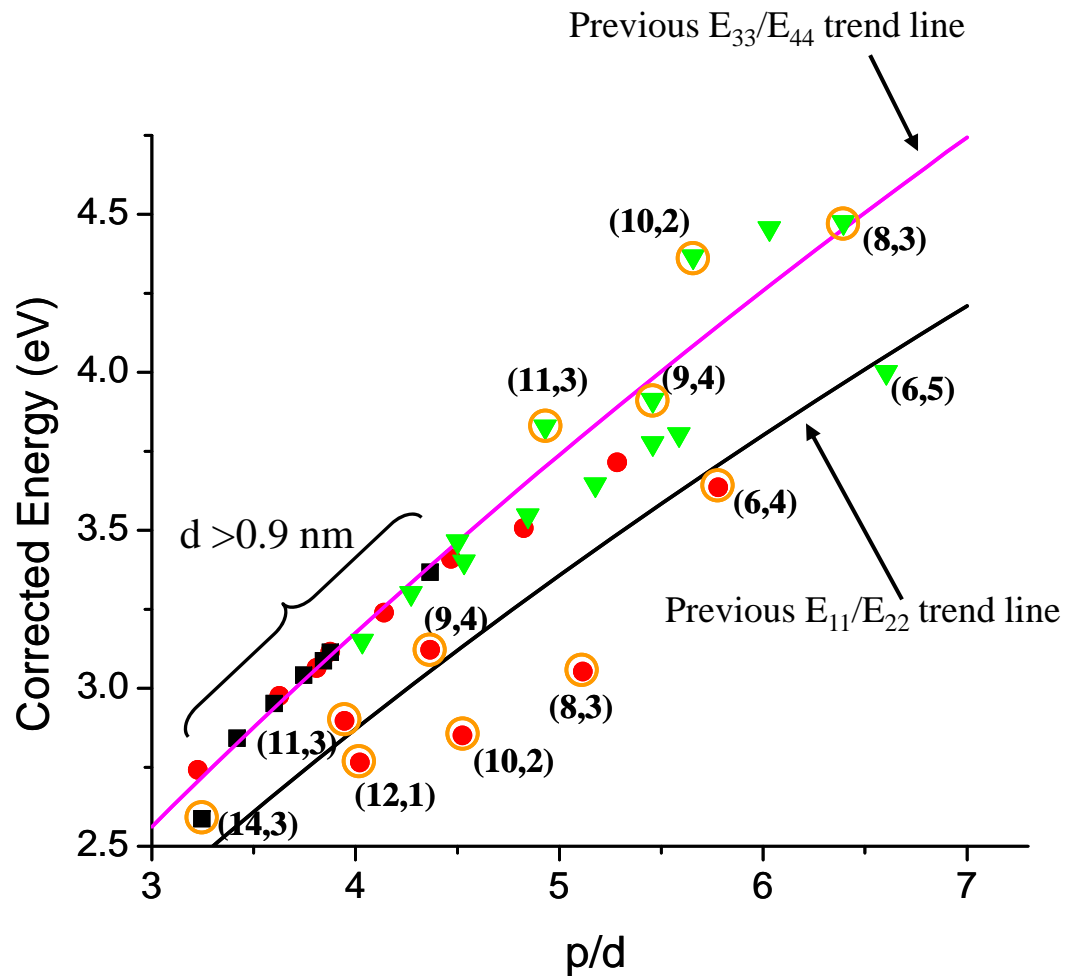
$E_{33}(\text{expt.}) > E_{33}(\text{ETB})$
(For PLE and Raman)

E_{44} trend less clear.

Haroz, et. al., *Phys. Rev. B*,
77, 125405 (2008).

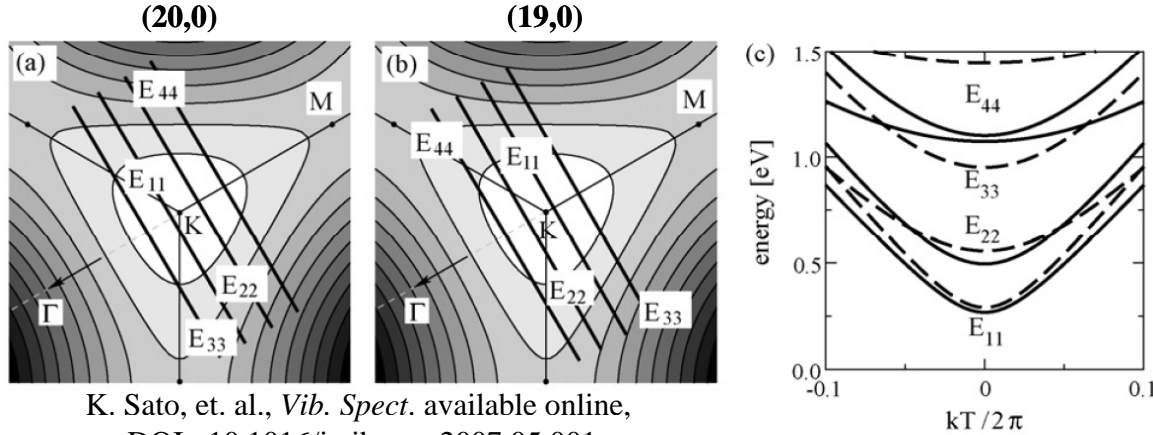
Scaling Law Analysis on Small Diameter Nanotubes

Haroz, et. al., *Phys. Rev. B*, **77**, 125405 (2008).



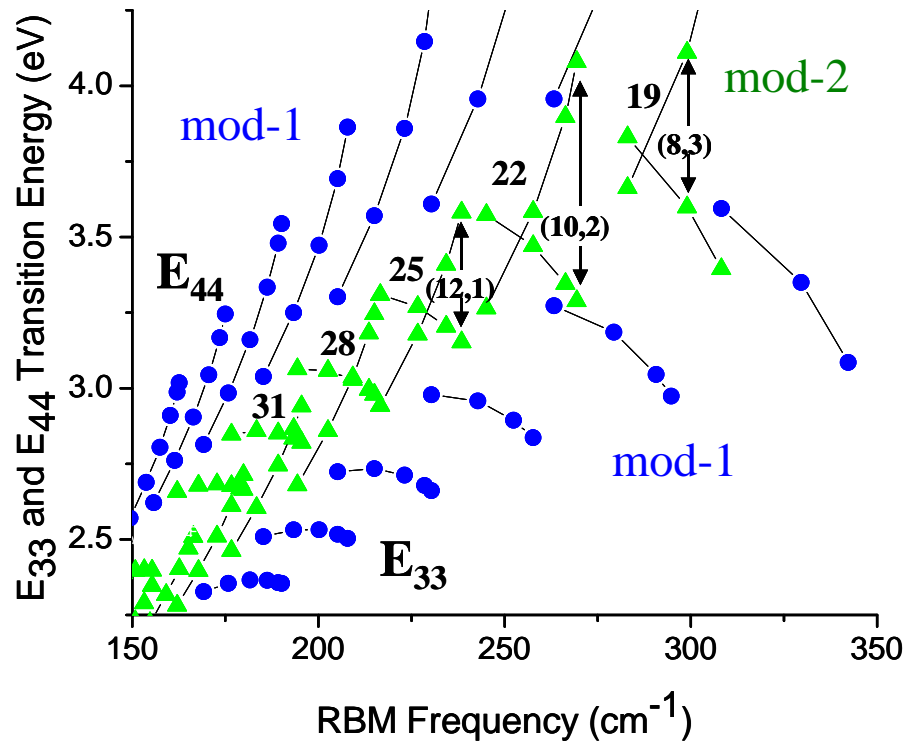
- Large d (>0.9 nm) nanotubes fit well to previous scaling trend.
- Small diameters appear to deviate.
- Significant scatter for selected chiralities.

Band Crossing in mod-2 Chiralities



K. Sato, et. al., *Vib. Spect.* available online,
DOI: 10.1016/j.vibspec.2007.05.001

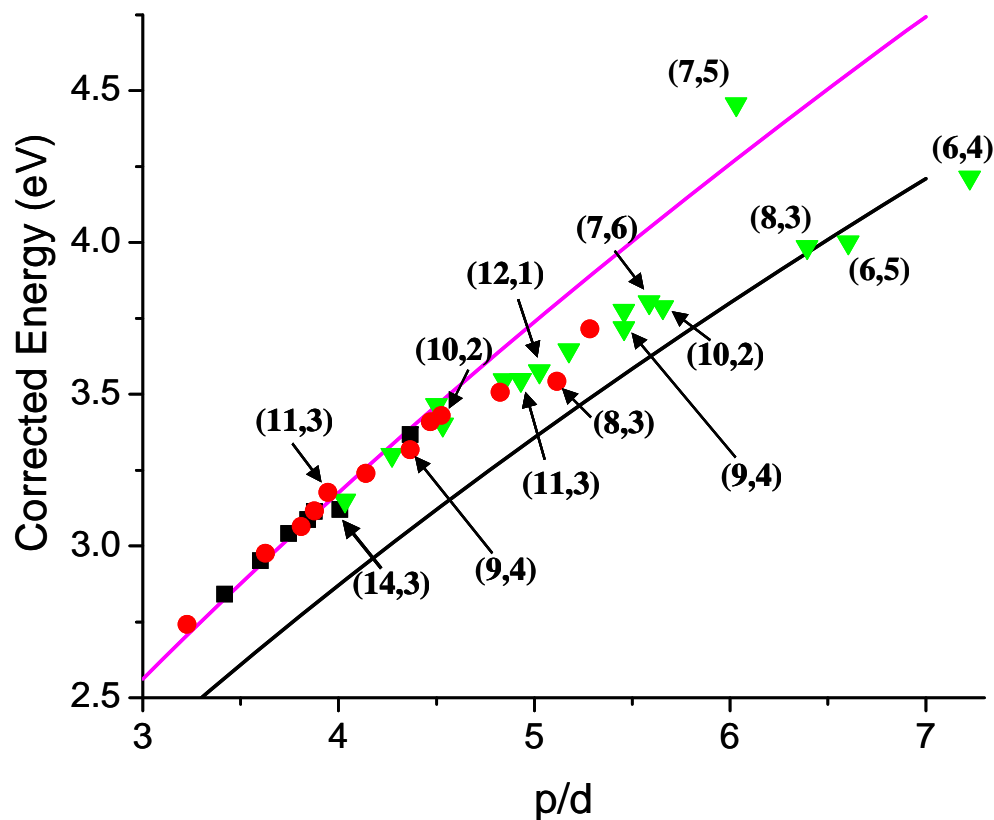
Mod-1 bends away.
Mod-2 bends toward.



--mod-2 E_{33} transitions
cross over E_{44} .

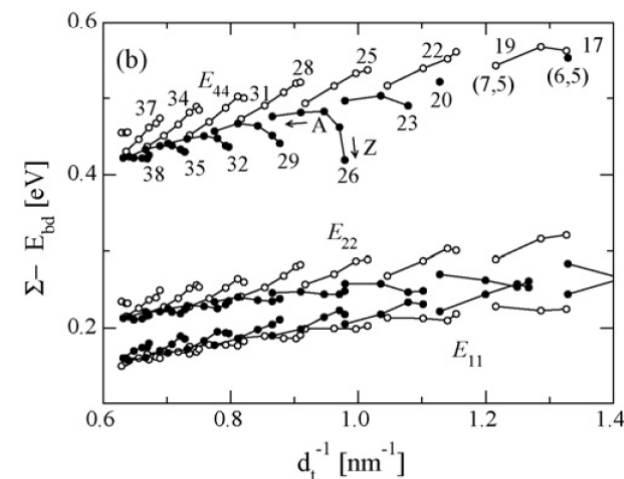
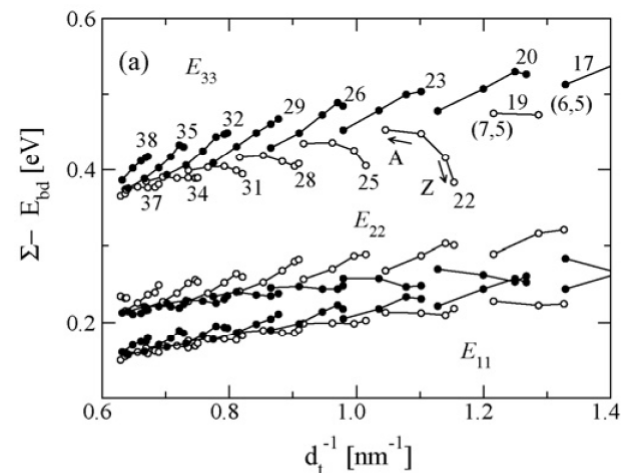
--Transition assignment
based on energy ordering
does not work.

Chirality Dependence of Many-Body Effects



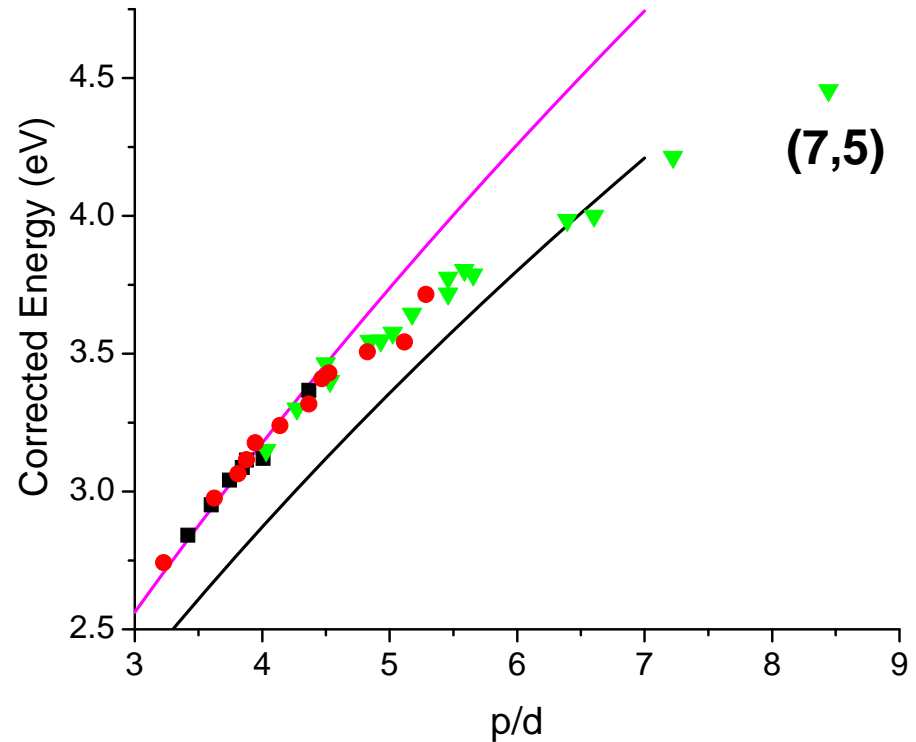
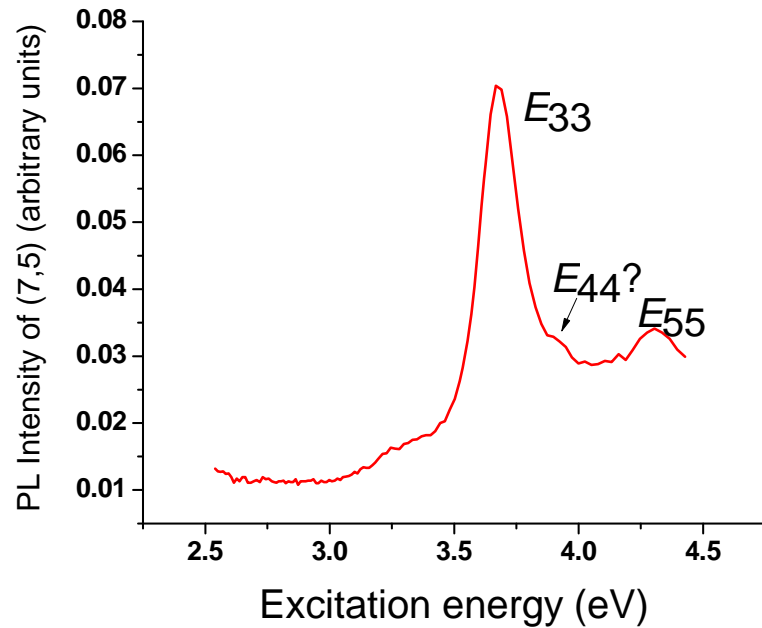
--Downward trend at small d is due to chiral angle dependence of the many-body effects.

--Many-body energy spread in opposite direction from that induced by trigonal warping effects.



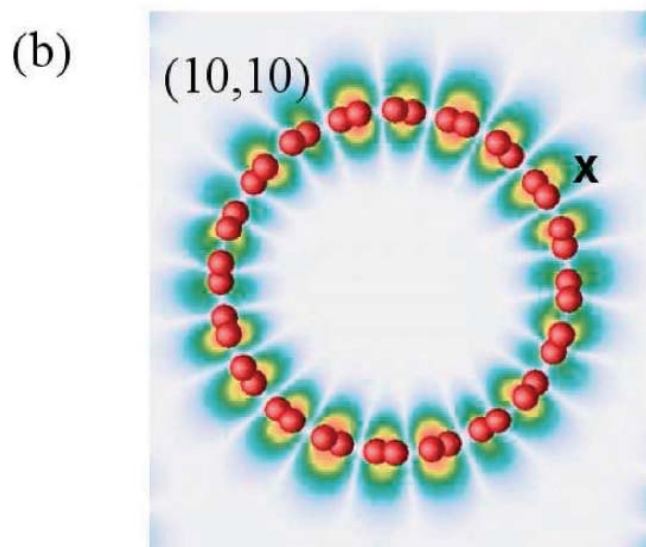
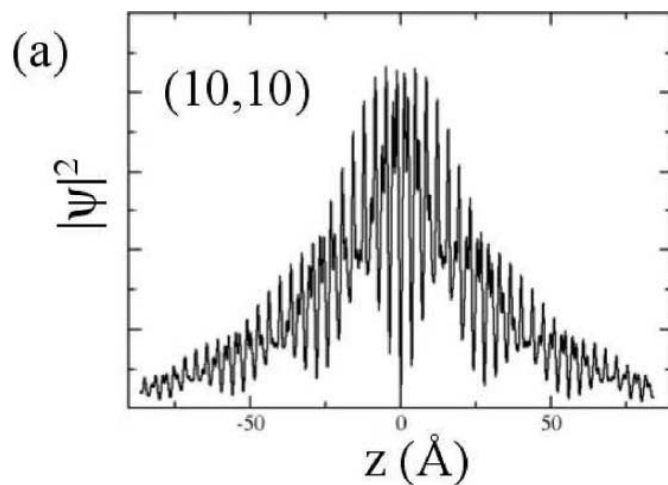
K. Sato, et. al., *Vib. Spect.* **45**, 89 (2007)

The (7,5) Assignment



If we consider the (7,5) point as an E_{55} transition energy, it matches the trend with the other chiralities.

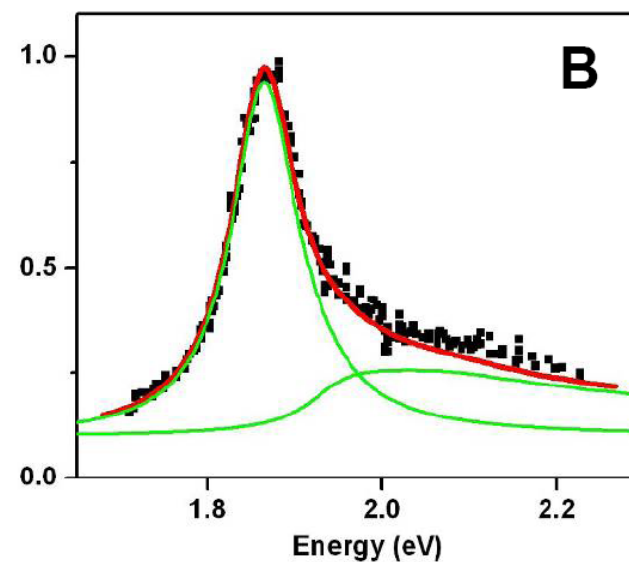
Excitons in One-Dimensional Metals



J. Deslippe et. al., *Nano Lett.*,
7, 1626 (2007).

Strong screening prohibits exciton formation in 2-D and 3-D metals.

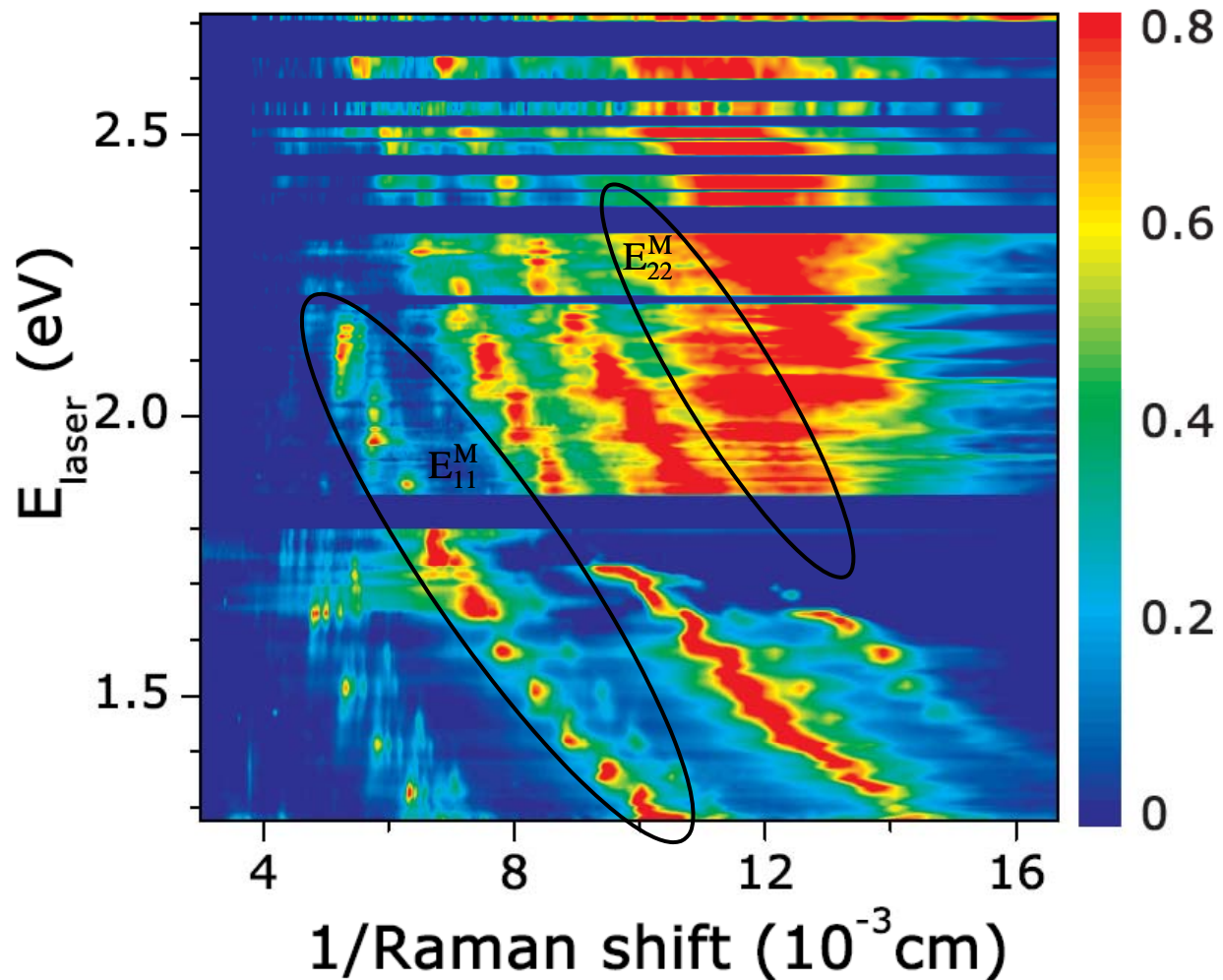
Long-range interactions overcome screening in 1-D metallic SWNTs.



F. Wang et. al., *Phys. Rev. Lett.*,
99, 227401 (2007).

Raman Analysis of Metallic Nanotubes

S.K. Doorn, et. al., *Phys. Rev. B*, **78**, 165408 (2008)



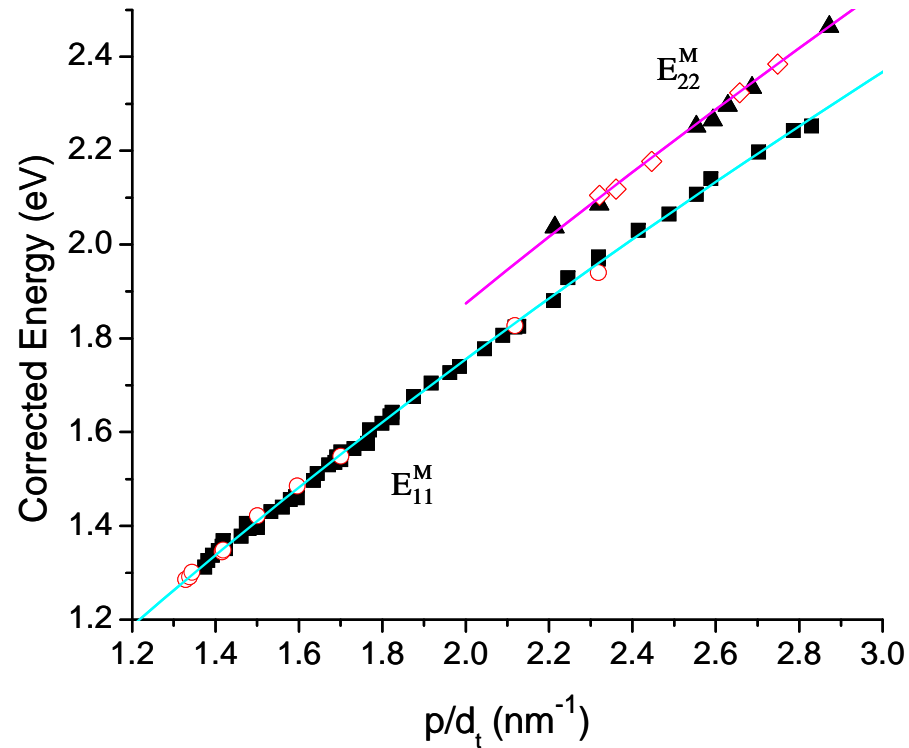
Both E_{11} and E_{22} metallic transitions observed.

Upper branches observed at large diameter.

Scaling Analysis of Metallic Energies

S.K. Doorn, et. al., *Phys. Rev. B*, **78**, 165408 (2008)

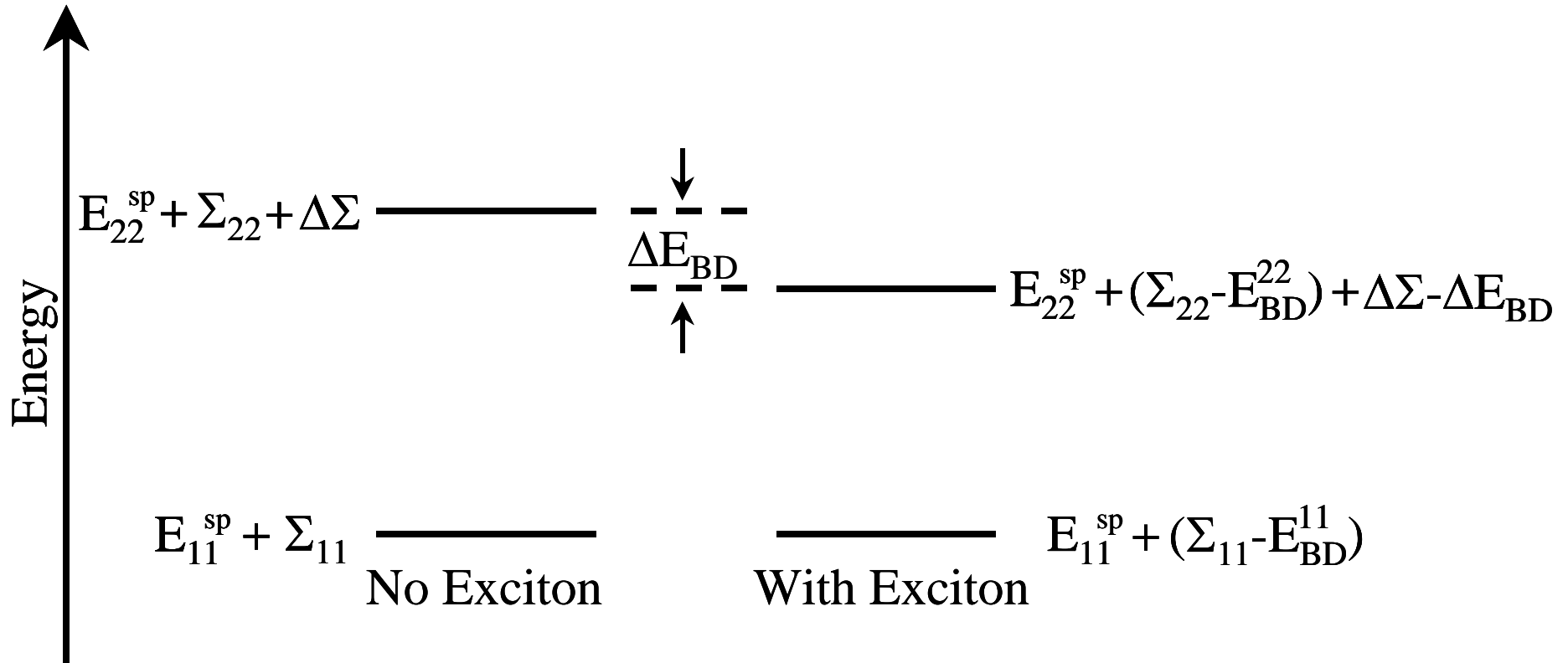
Semiconducting scaling lines
compared to metallic energy data.



Overlap of metallic scaling with semicond.
consequence of similar sampling of BZ.

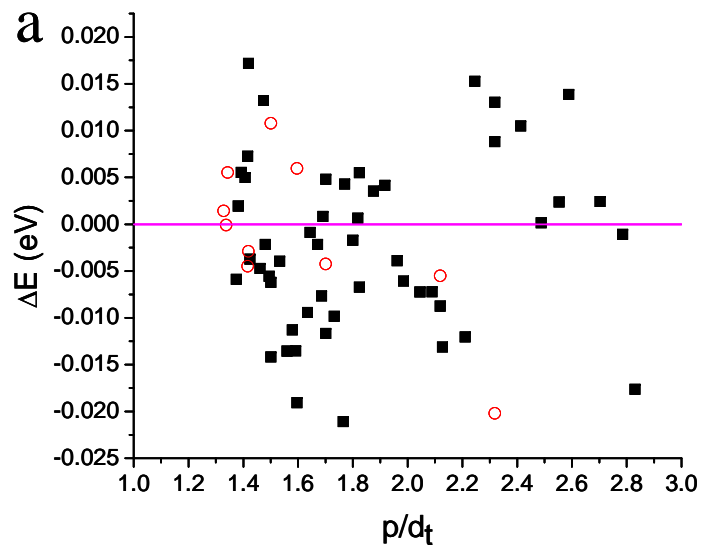
Scaling behavior provides indirect
evidence of exciton formation.

Schematic Analysis of Relative Scaling Behaviors

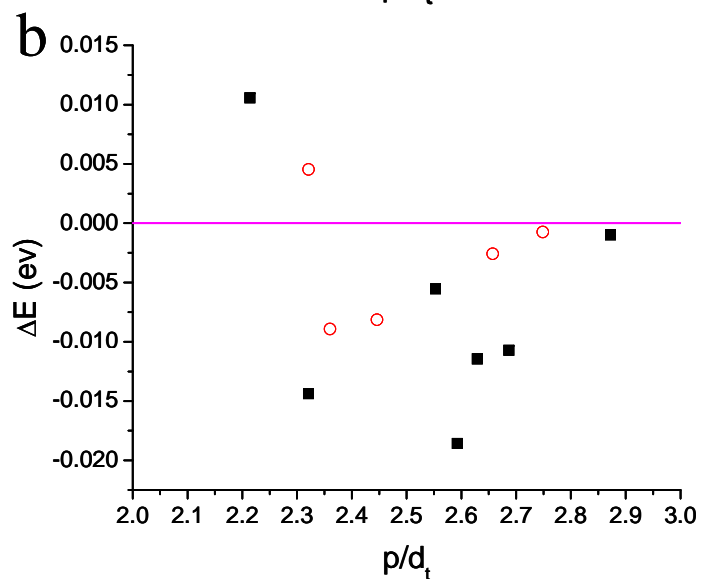


Lack of exciton formation should place E_{22} metallic data above the semiconductor scaling line.

Metallic and Semiconductor Scaling Law Comparisons

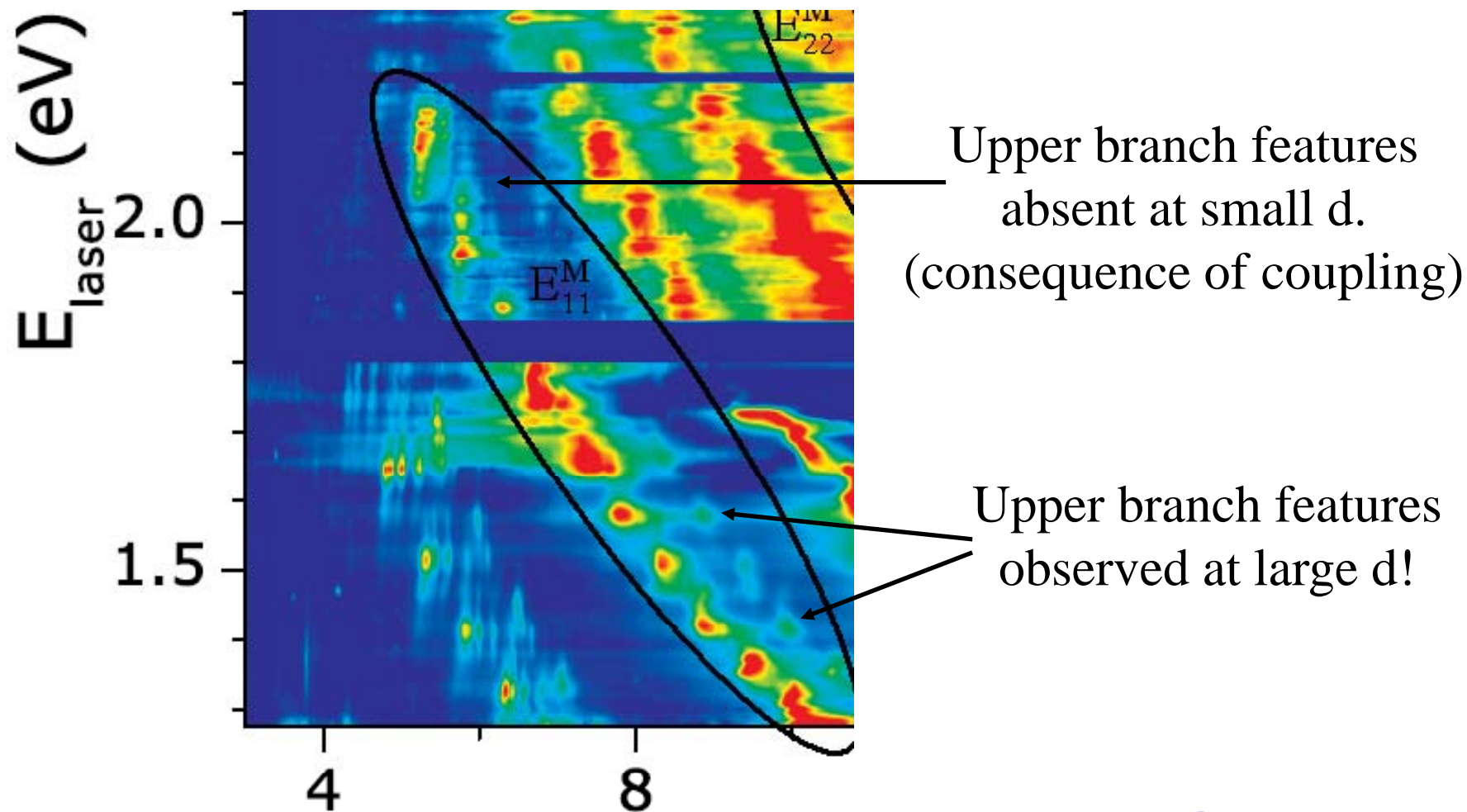


E_{11} metallic data evenly
distributed about trend line.

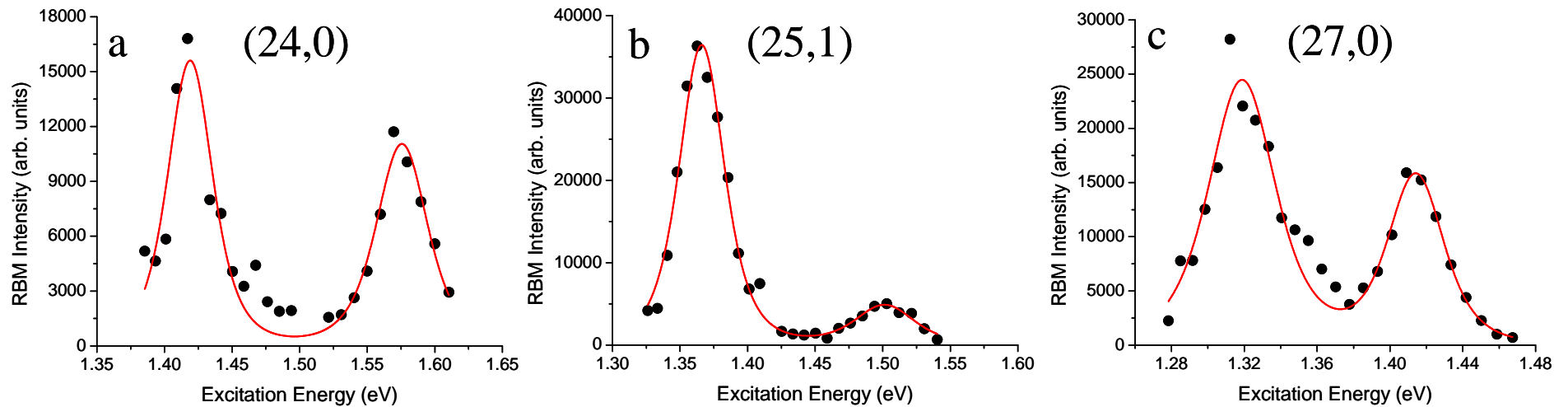


E_{22} metallic data
weighted
below the trend line.

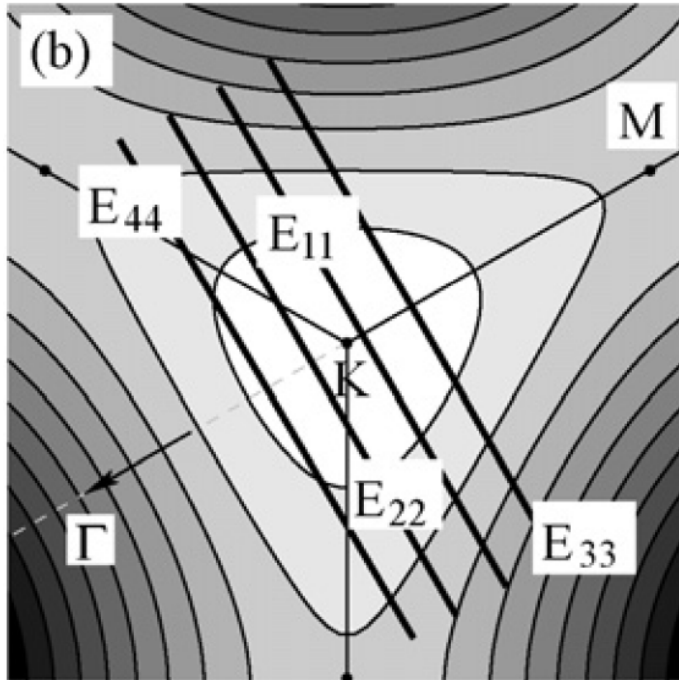
Exciton-Phonon Coupling In Metallic Nanotubes



Comparable Upper and Lower Branch Intensities

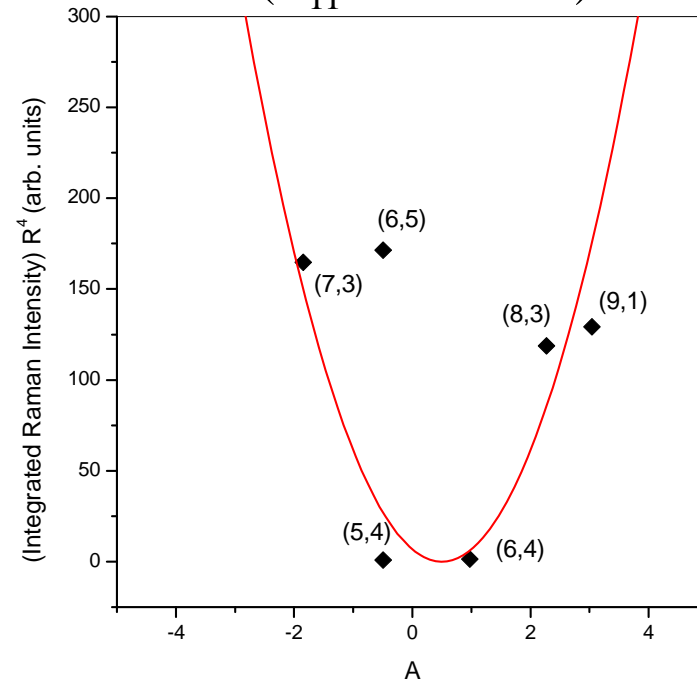


K- Γ vs. K-M Valley Coupling



Transitions originating in K-M valley expected to be stronger.

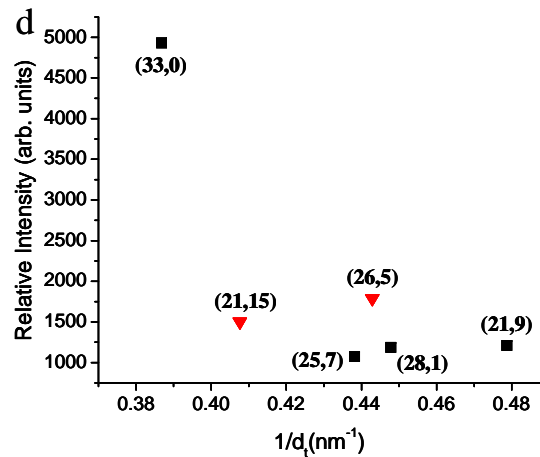
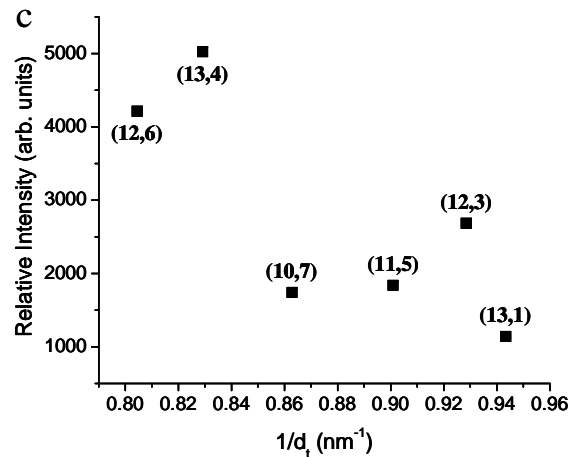
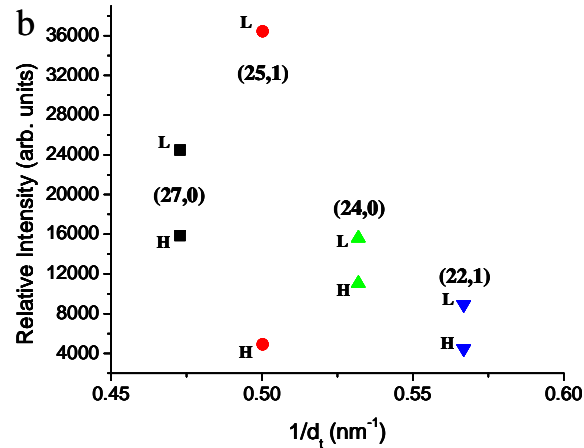
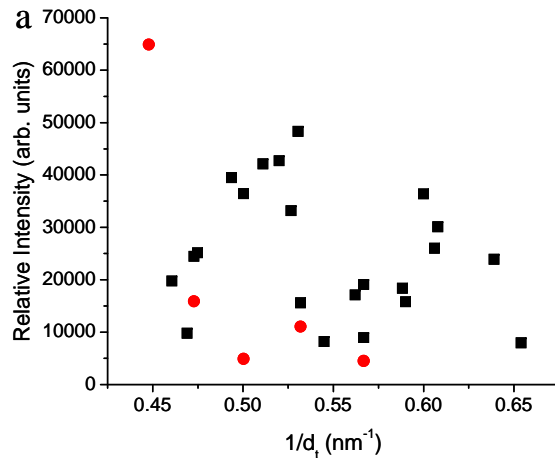
Nodal behavior of coupling.
(E_{11} Excitation)



Larger diameter tubes in the K- Γ valley can lie further from the low intensity node.

Goupalov, et. al., *Phys. Rev. B*, **73**, 115401 (2006).

General Metallic Intensity Behavior

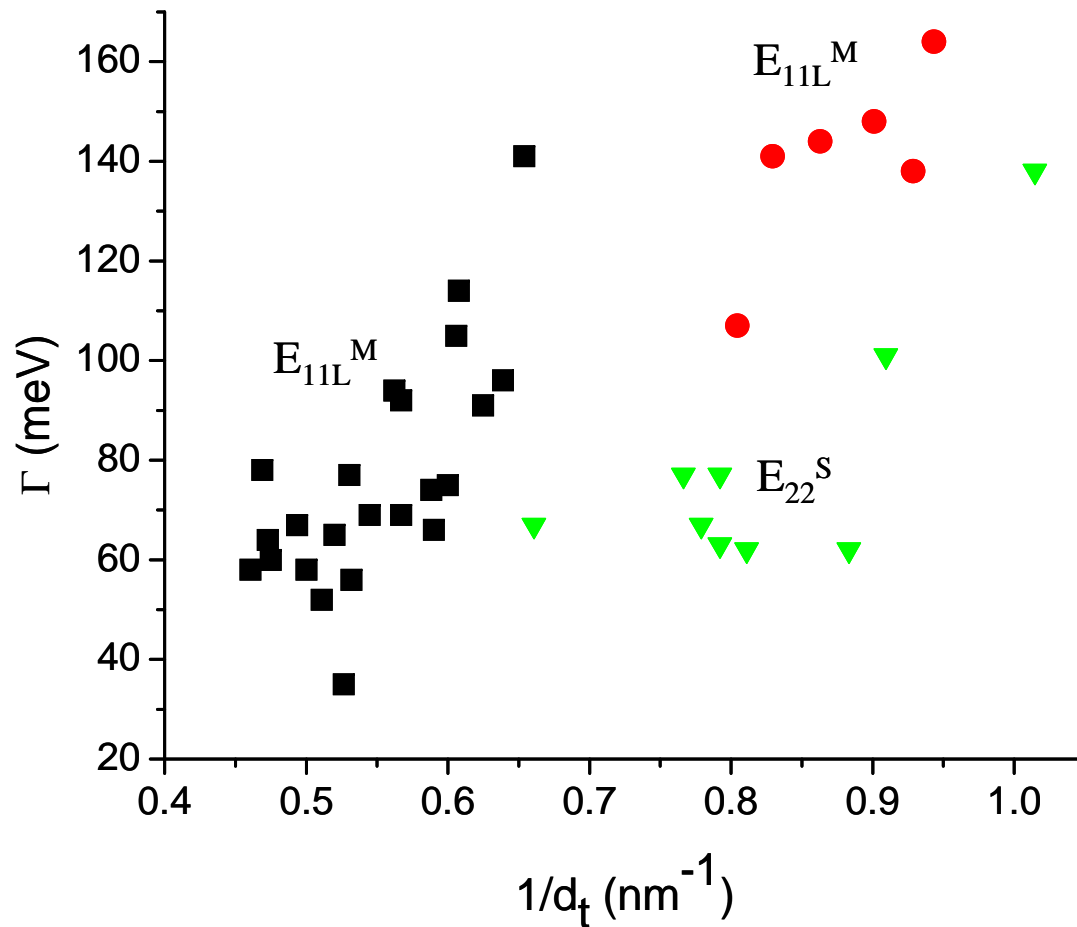


--Intensity decreases as diameter decreases.

--Opposite of expectation based on dependence of exciton-phonon coupling.

Diameter Dependence of Gamma

S.K. Doorn, et. al., *Phys. Rev. B*, **78**, 165408 (2008)



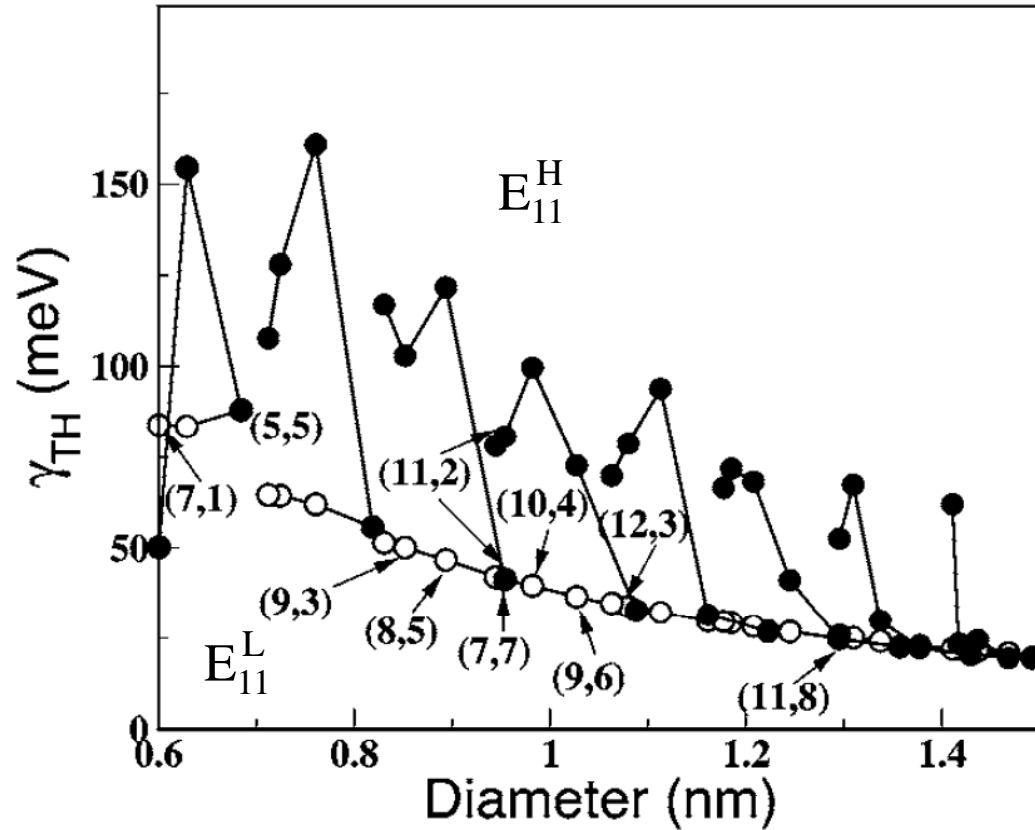
-- Γ increases as d decreases.

-- Γ influence on intensity overcomes coupling dependence.

-- Γ larger for metallics. Consequence of coupling to the plasmon.

Expectations From Theory

Park et. al., *Phys. Rev. B*, **74**, 165414 (2006).



--at large d , Γ increases nearly linearly as d decreases.

--upper and lower branch Γ similar at large d .

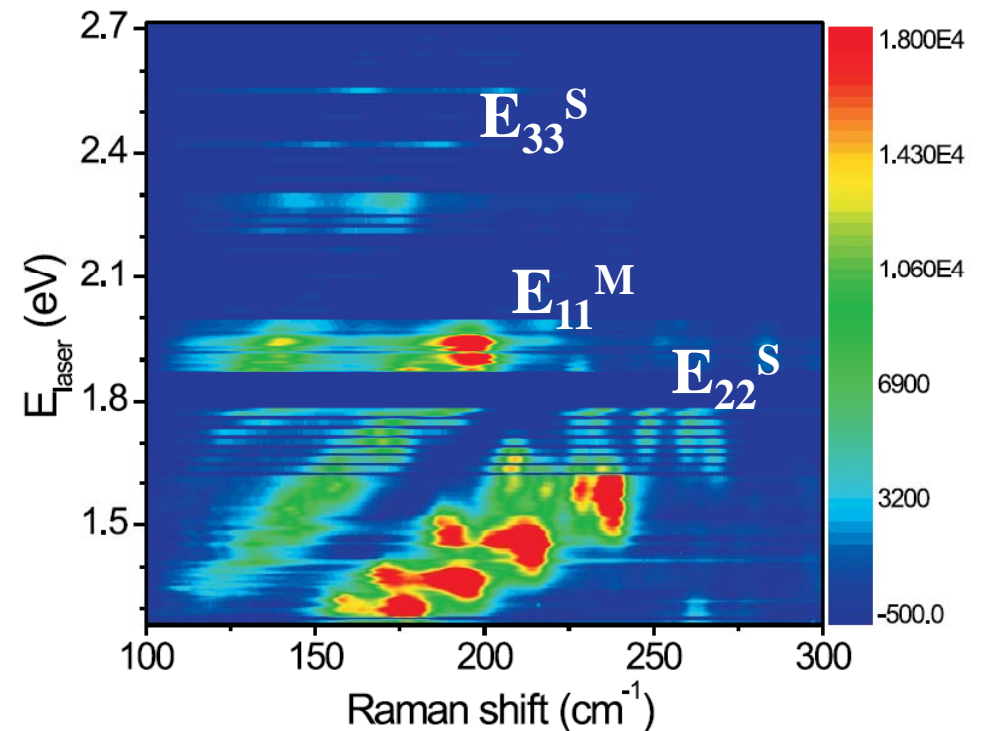
Parallels to Semiconductor Intensity Behavior

E11/E22

(n, m)	I_{11}	I_{22}
(6,4)	0.26	0.054
(8,3)	12.6	1.03
(9,1)	15.6	0.33
(5,4)	0.21	0.0037
(6,5)	20.7	0.0720
(7,3)	16.4	0.0014

B.C. Satishkumar, et. al., *Phys. Rev. B*, **74**, 155409 (2006).

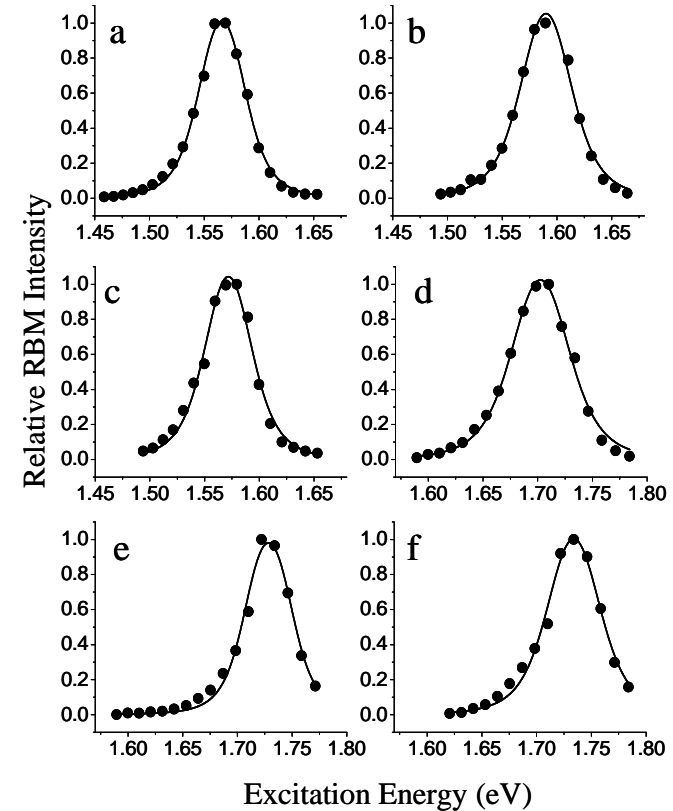
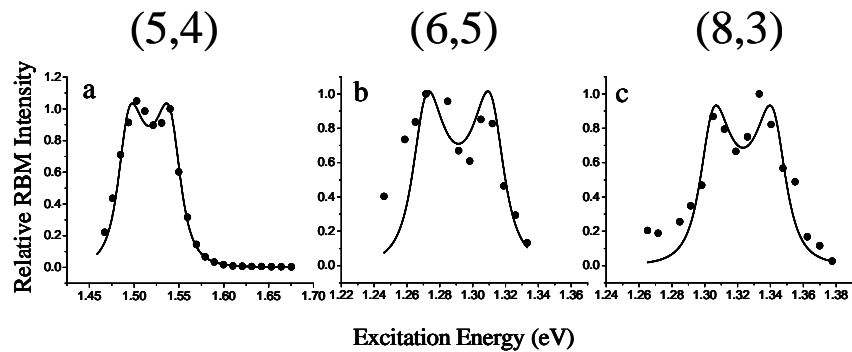
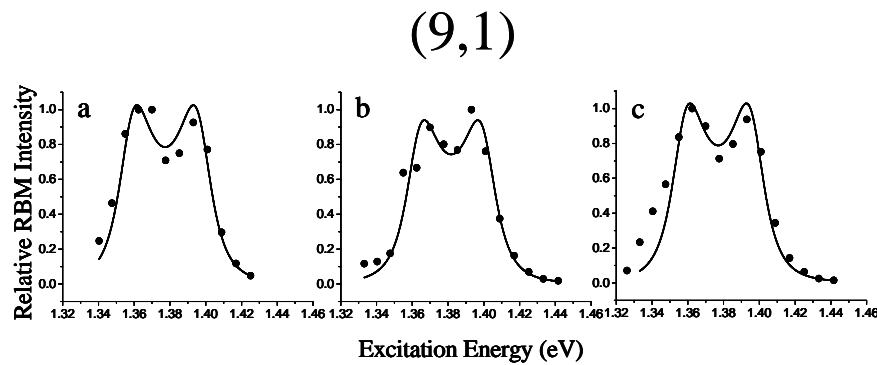
E33



P.T. Araujo, et. al., *Phys. Rev. Lett.*, **98**, 067401 (2007).

E11/E22 profiles and Gammas

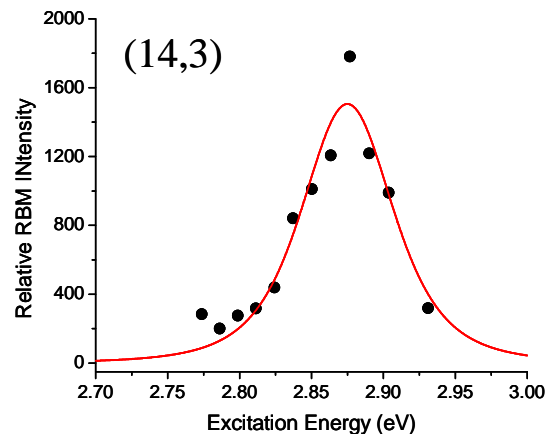
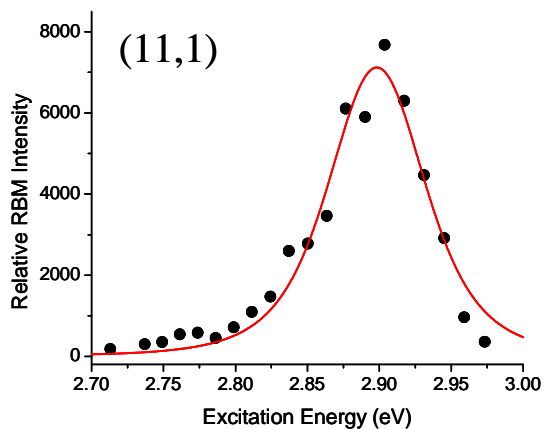
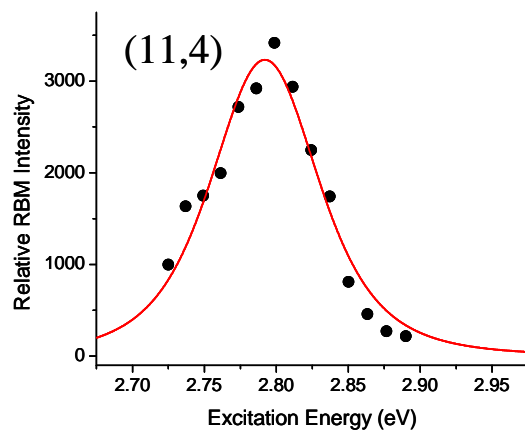
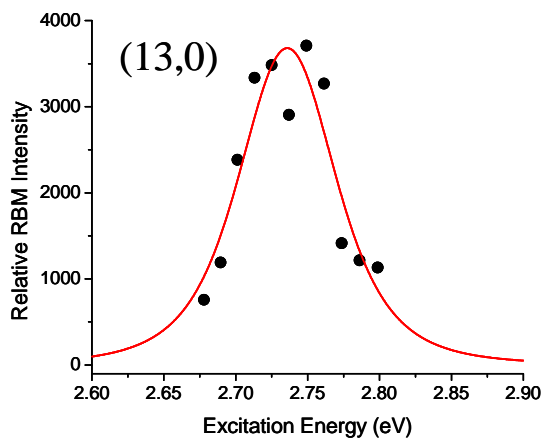
$$\alpha = \sum_{i,j} \frac{M_{ee}^{g,i} M_{ee}^{i,g} M_{ep}^{i,j}}{(E_{\text{laser}} - E - i\Gamma_r) (E_S - E - i\Gamma_r)}$$



$$\Gamma_{11} \sim 20\text{-}30 \text{ meV}$$

$$\Gamma_{22} \sim 60\text{-}70 \text{ meV}$$

E33 Excitation Profiles

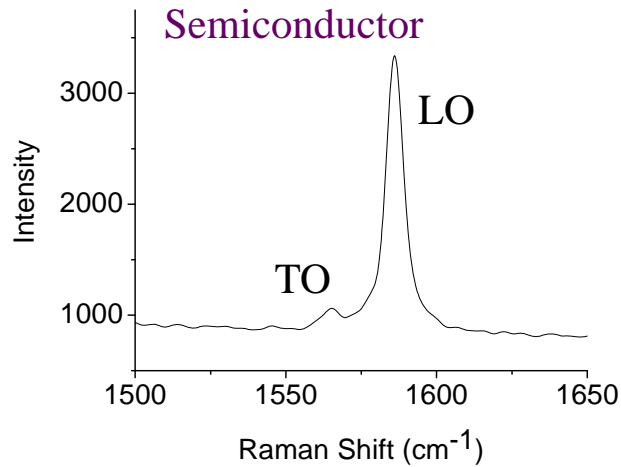


- Trend to larger Γ_{ii} continues.
- Expect D.O.S. and relaxation pathways to increase.
- Agrees with quantum chemical results—increased number of closely spaced states.

Haroz, et. al., *Phys. Rev. B*,
77, 125405 (2008).

$$\Gamma_{33} \sim 110-130 \text{ meV}$$

Phonon Coupling to Low-Energy Excitations: G-Band



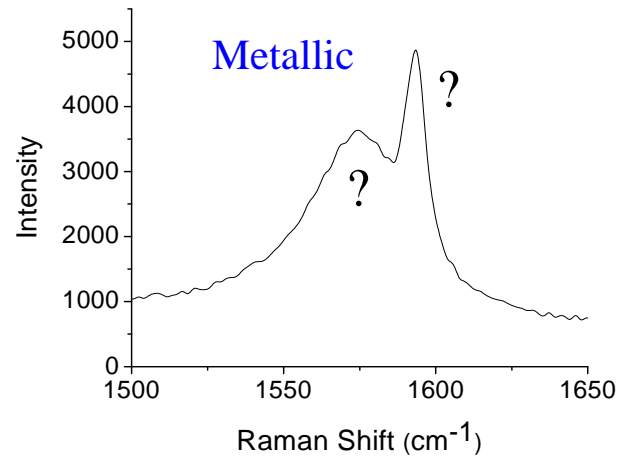
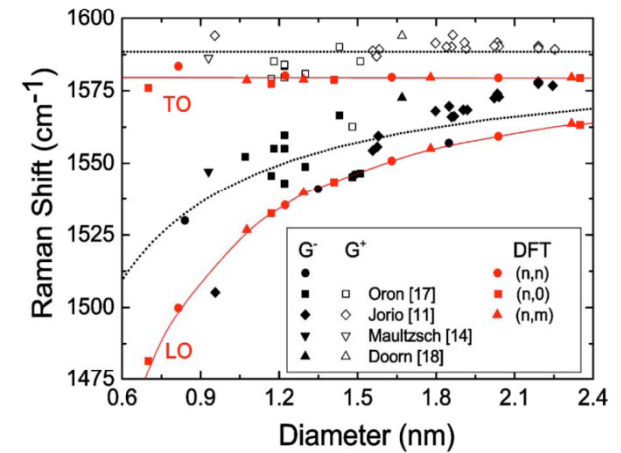
Kohn Anomaly



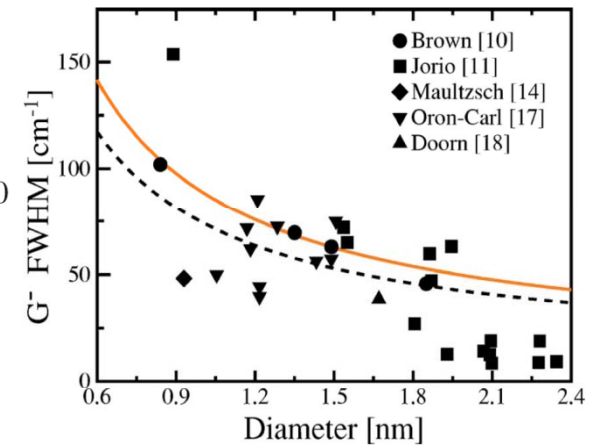
ω decrease for strongly coupled modes

LO: EPC Large
TO: EPC = 0

$$\omega_{LO}^2 = A^2 - B/d; \quad \omega_{TO}^2 = C$$



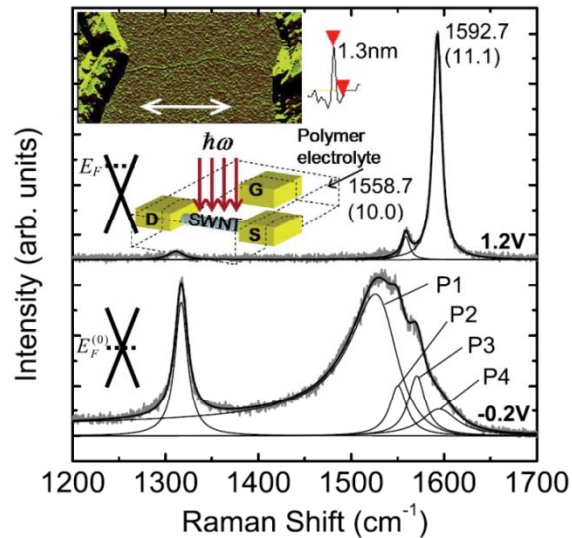
$$\gamma_{\Gamma-LO}^{EP} = \frac{2\sqrt{3}\hbar a_0^2 \langle D_{\Gamma}^2 \rangle_F}{\pi M \omega_{\Gamma} \beta d}; \quad \gamma_{\Gamma-TO}^{EP} = 0$$



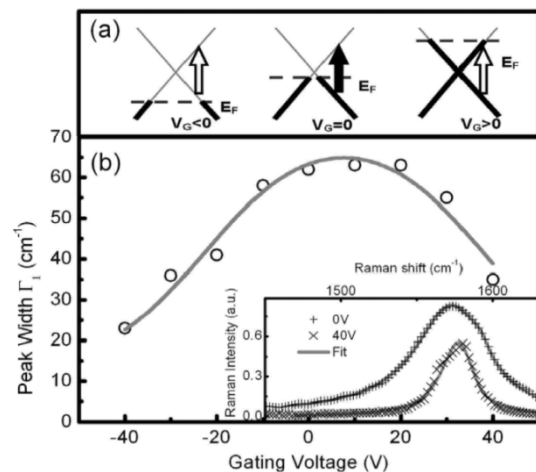
Plasmon Coupling?
-OR-
Excitations at Fermi Level?

Lazzeri, et. al., *Phys. Rev. B*, **73**, 155426 (2006).

G-Band Lineshape: Gating and Chirality Effects



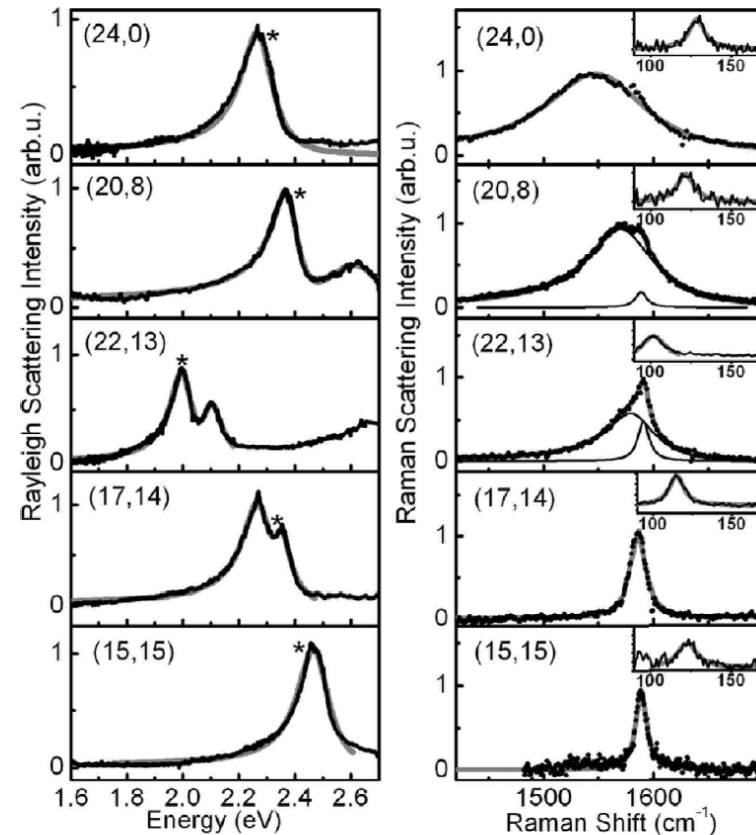
Intensity (arb. units)
Raman Shift (cm^{-1})
Nguyen, et. al., *Phys. Rev. Lett.*
98, 145504 (2007).



Peak Width Γ_1 (cm^{-1})
Gating Voltage (V)
Wu, et. al., *Phys. Rev. Lett.*
99, 027402 (2007).

Electrostatic Gating

- Shifts Fermi Level
- Excitations > 200 meV
- no G-band coupling
- LO Narrows



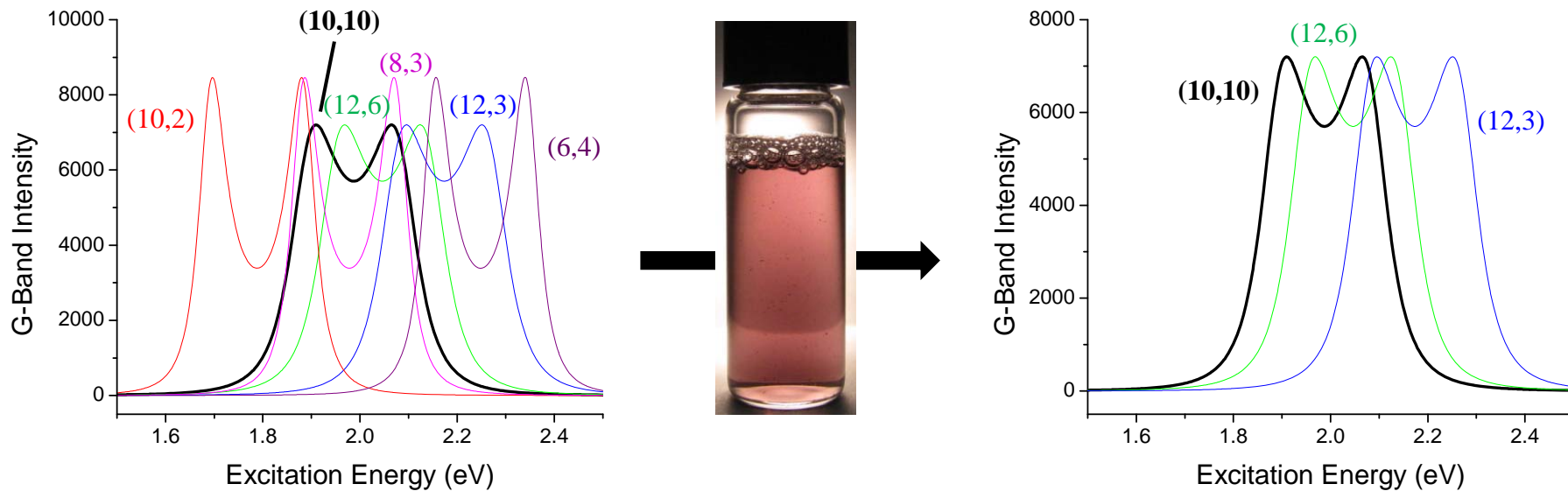
Rayleigh Scattering Intensity (arb.u.)
Raman Scattering Intensity (arb.u.)
Energy (eV)
Raman Shift (cm^{-1})
Wu, et. al., *Phys. Rev. Lett.*
99, 027402 (2007).

Lineshape inherent to single metallic tube.

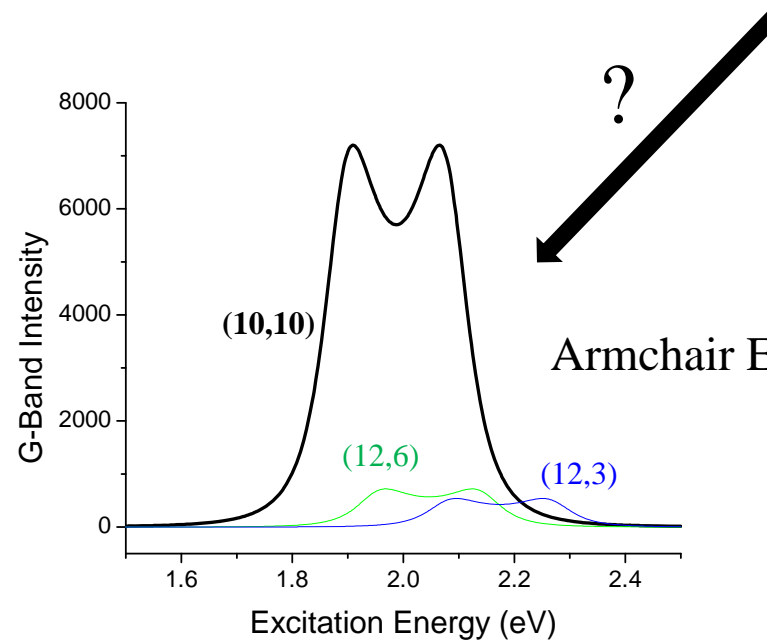
Strong chirality dependence.

**Is single sharp G^+ band a
general result for armchairs?**

Ensemble Spectra: Approaching the Pure Armchair Limit



G-Band not as Clean as RBM



DGU Separations:
“Eliminate” Semiconductors

Armchair Enrichment?

Density Gradient Separations

- Widely used in biochemistry and pharmaceutical industry.
- Separation based on very subtle differences in buoyant densities of components—diameter dependent.
- Ultracentrifugation within a preformed gradient that varies in density with height.
- Nanotubes separate into levels of gradient of matching density upon ultracentrifugation.
- Typically requires use of cosurfactant (various combinations of SDS, SDBS, cholates) suspension of nanotubes.

Summary

Raman is a sensitive probe of structure, phonon-coupling, and nature of excited states.

First evidence for transition cross-over effect and chirality dependence of many-body influences.

Support for the existence of excitons in 1-D metals.

First Raman evidence of metallic upper branches.

G+ (TO) dominance in armchair spectra is a general behavior.

Acknowledgements

Curvature Effects

On E33/E44

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R. Bruce Weisman (Rice)

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Kenji Hata (AIST)

Armchair Raman

Erik Haroz (Rice)

Jun Kono (Rice)

