

One-Dimensional Electron and Thermal Transport in Suspended Carbon Nanotubes

Steve Cronin, I-Kai Hsu, Adam Bushmaker

University of Southern California

Vikram Deshpande, Scott Hsieh, Marc Bockrath

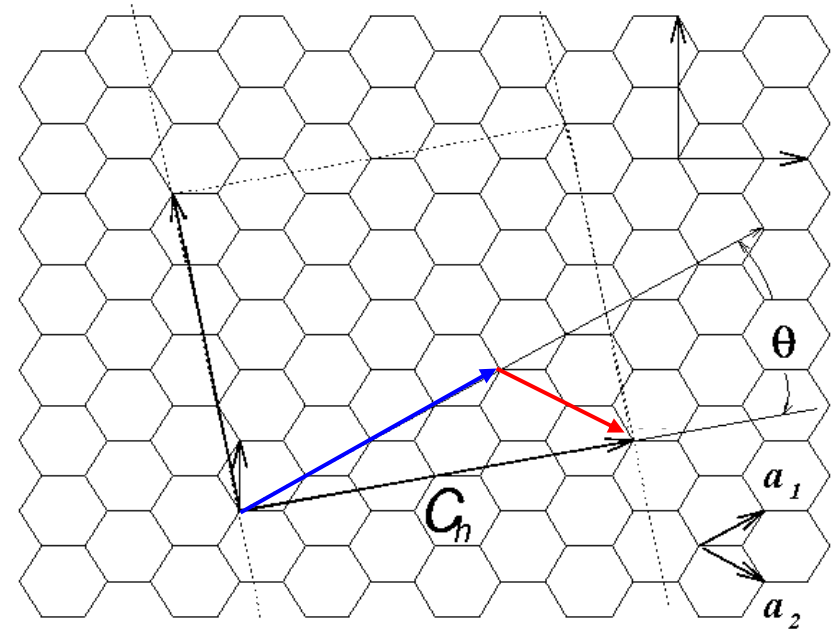
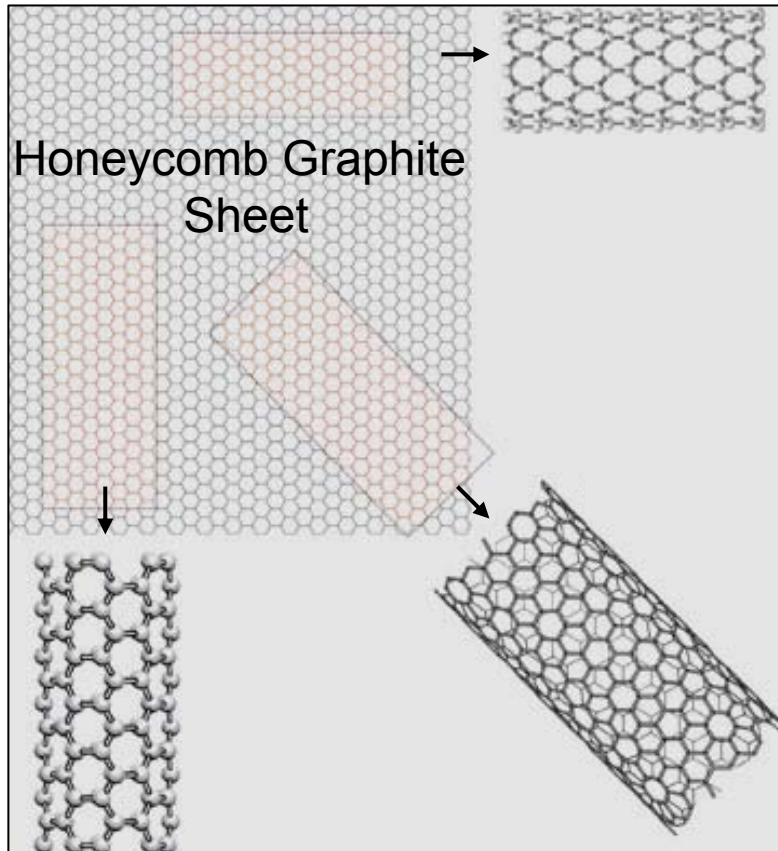
California Institute of Technology

Michael Pettes and Li Shi

University of Texas, Austin

What is a Carbon Nanotube?

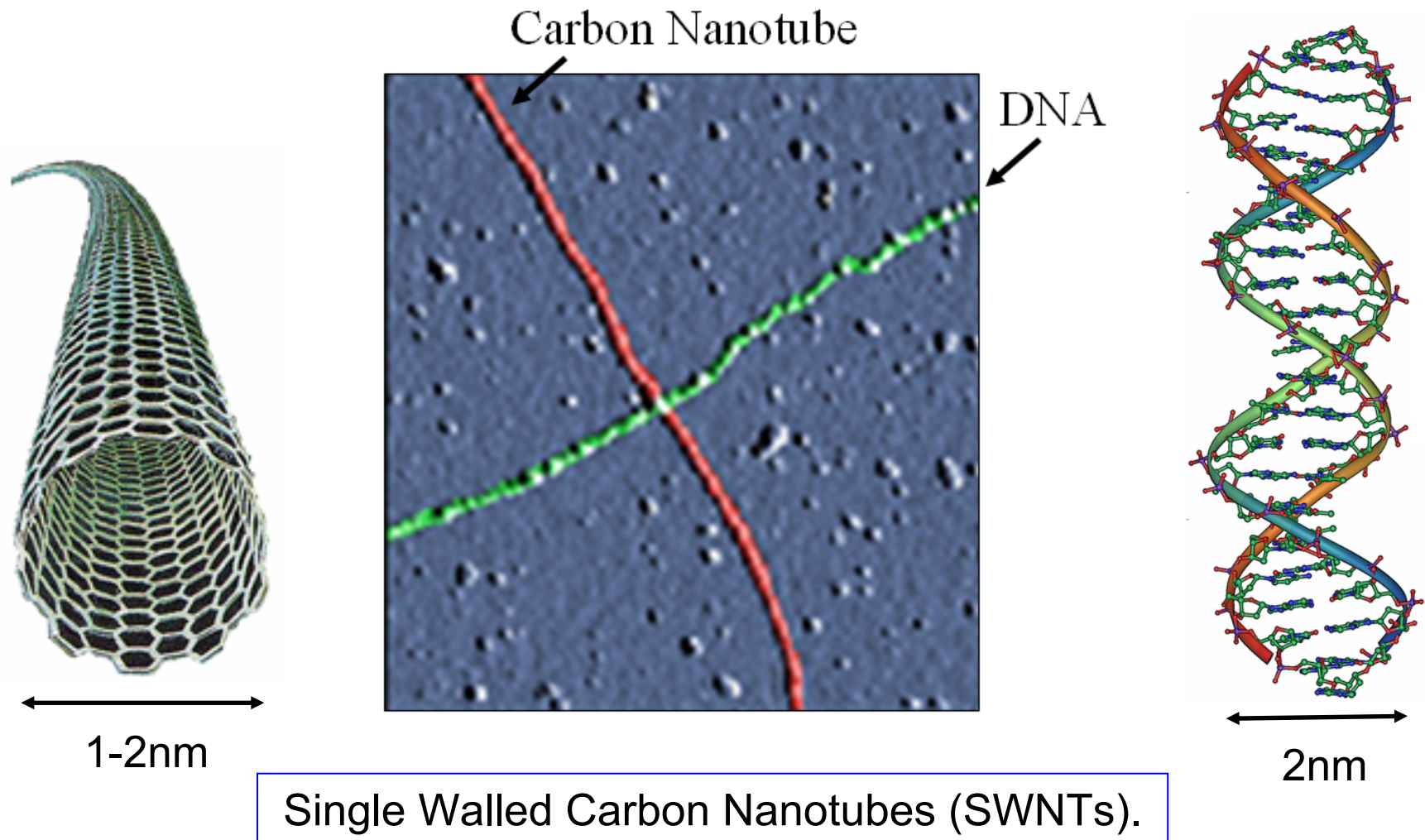
Imagine rolling a sheet of graphite into a seamless cylindrical tube.



Chirality (n,m) :

$$\mathbf{C}_h = 4\mathbf{a}_1 + 2\mathbf{a}_2 = (4,2)$$

Atomic Force Microscope (AFM)



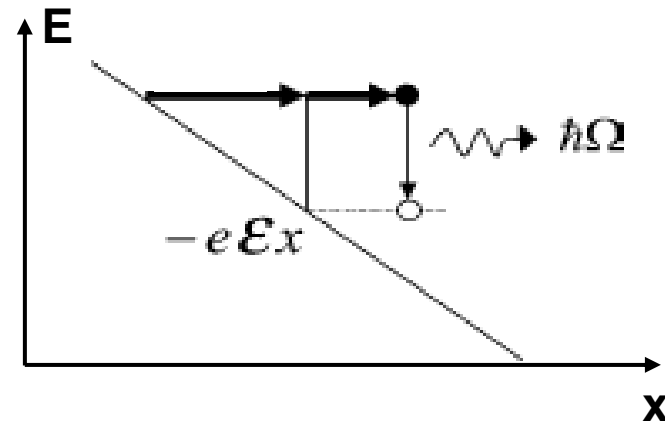
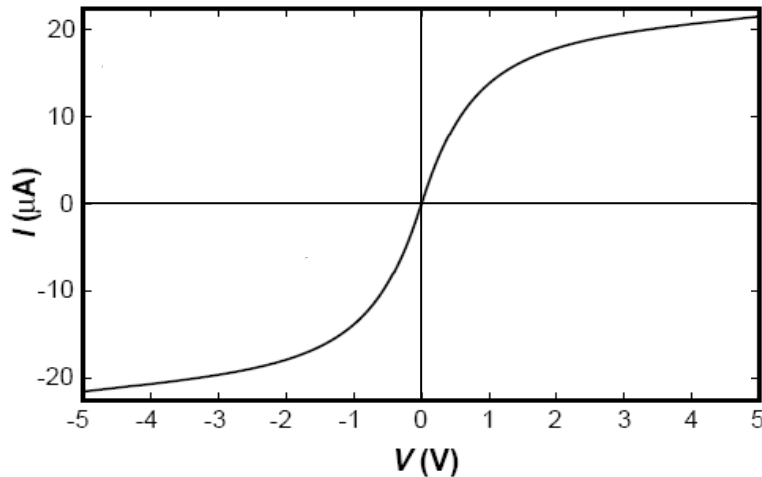
Why Study Carbon Nanotubes?

- 1nm in diameter, 1cm in length, aspect ratio $\sim 10^7$
- 100% surface-to-volume ratio
- 1 defect in 10^{12} C atoms \Rightarrow ballistic conduction
- High melting point $\sim 3800^\circ\text{C}$
- High Young's modulus 1TPa ($>$ diamond)
- High electronic current carrying capacity ($10^9\text{A}/\text{cm}^2$)
 $\sim 10^3$ times higher than that of the noble metals
- Thermal conductivity 6600W/mK at room temperature is twice the maximum known bulk thermal conductor, isotropically pure diamond = 3320W/mK

Despite 46,000 publications, no large scale commercial applications of nanotubes

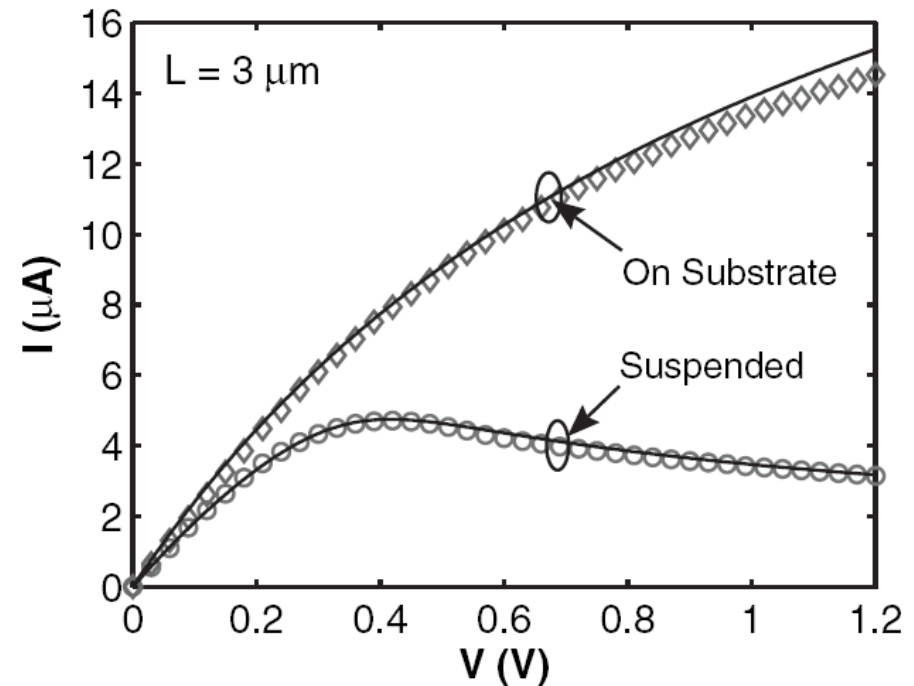
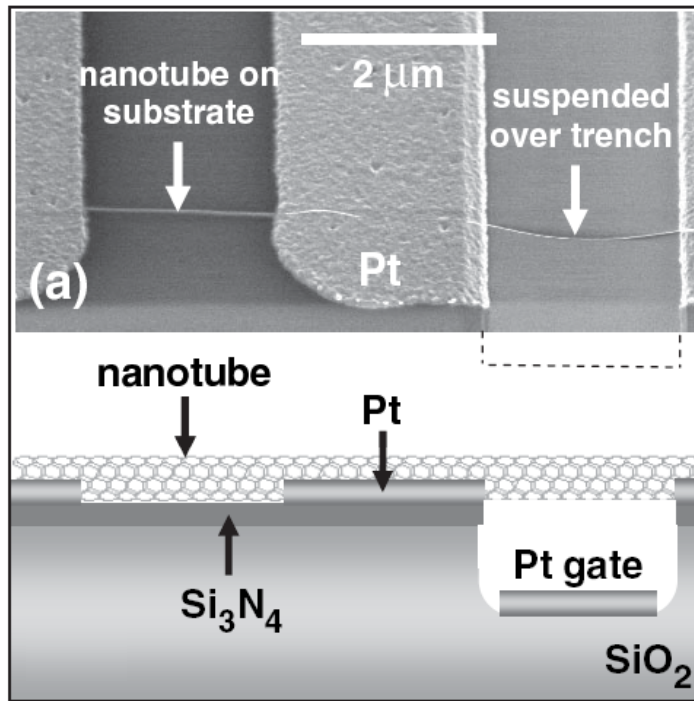
Review: High Field Transport

- Carbon nanotubes have extremely high current density threshold $\sim 10^9 \text{A/cm}^2$.



- Yao *et al.* proposed G band optical phonon emission.

Review: High Field Transport



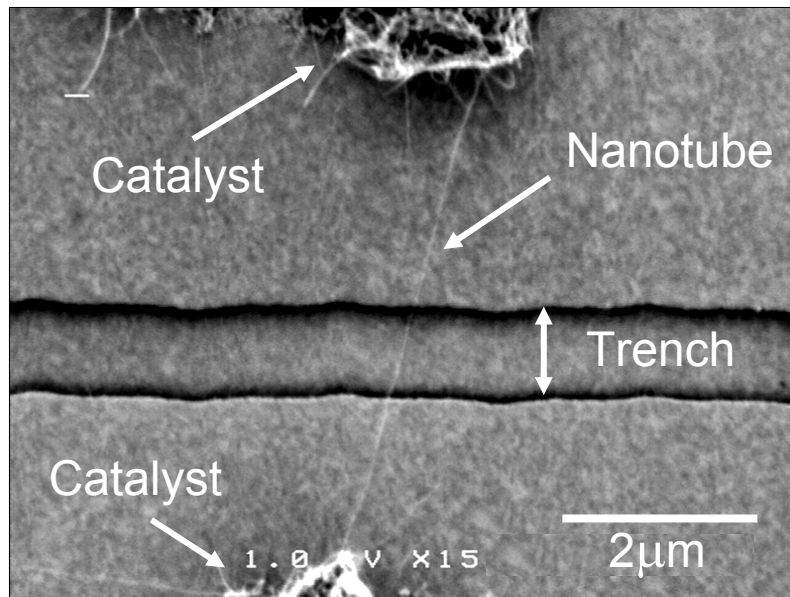
- Suspended nanotubes with low contact resistance show negative differential conductance.
 - Landauer Model

Using Raman Spectroscopy...

What is the temperature of the nanotube at each point on the I-V curve?

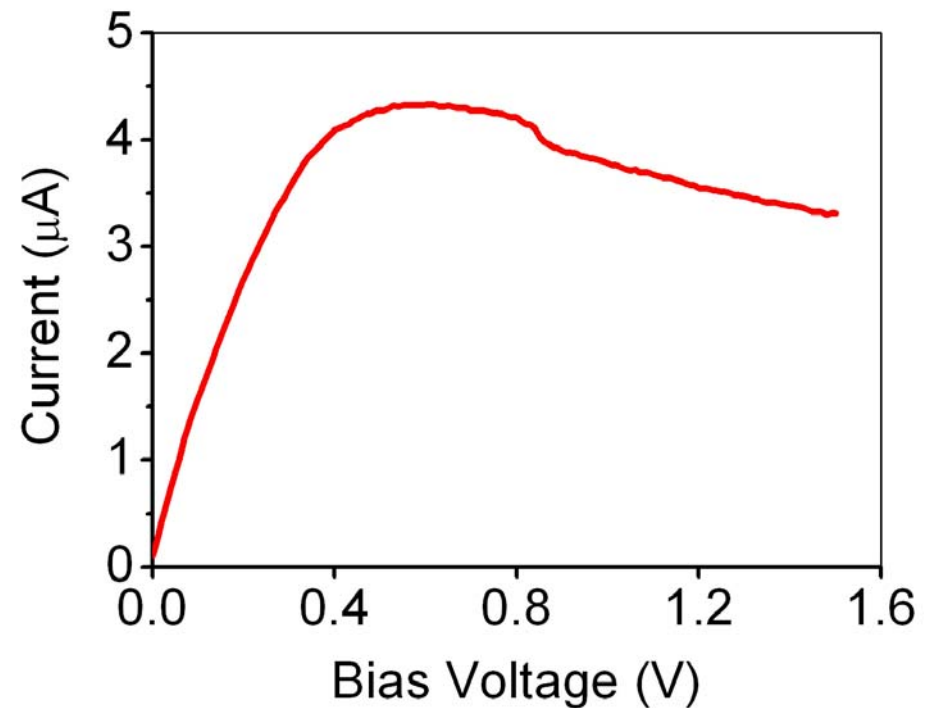
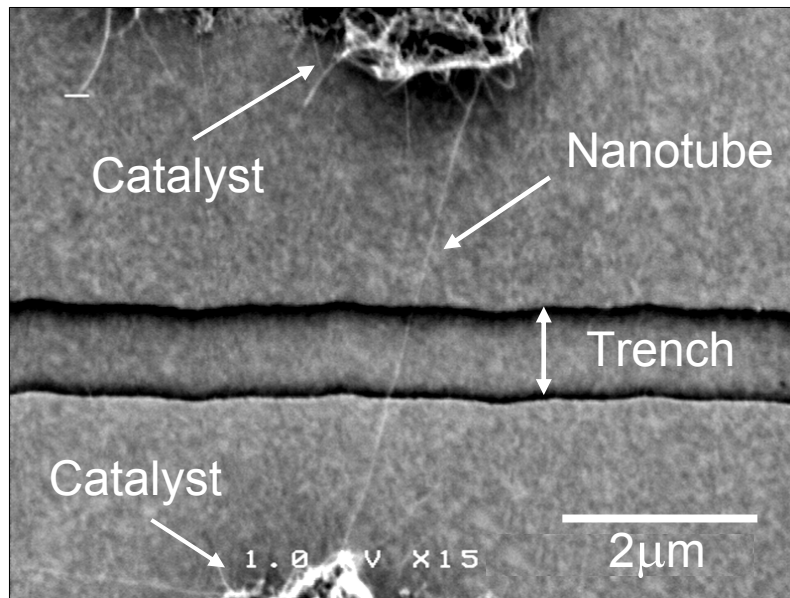
Can we observe optical phonon emission?

Sample Fabrication



- Pt electrodes deposited on top of trenches etched in Si/SiO₂.
- Catalyst patterned lithographically
- Nanotubes are growth in a mixture of methane and hydrogen over the wafer at 800°C.
- Measure the nanotubes, as-grown.

Negative Differential Conductance

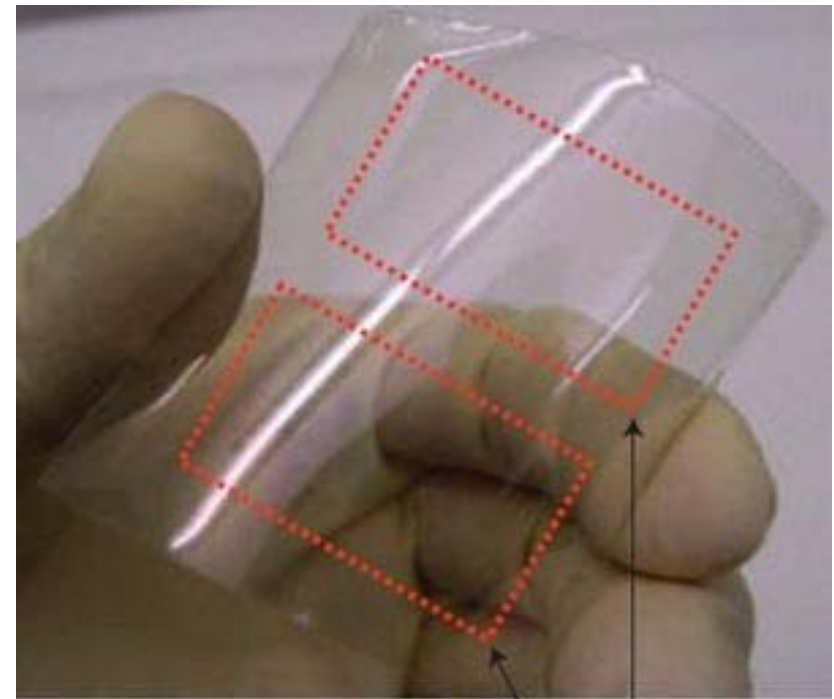


- Selected individual metallic nanotubes exhibiting NDC
- Pristine, highly crystalline, as-grown nanotubes
- Measure the intrinsic properties of nanotube

Carbon Nanotube Applications



Space Elevator

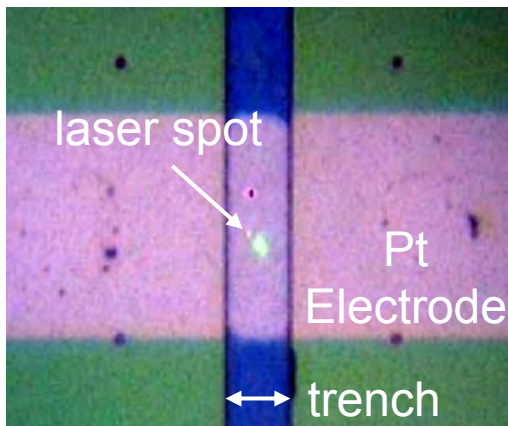


Transistor array regions

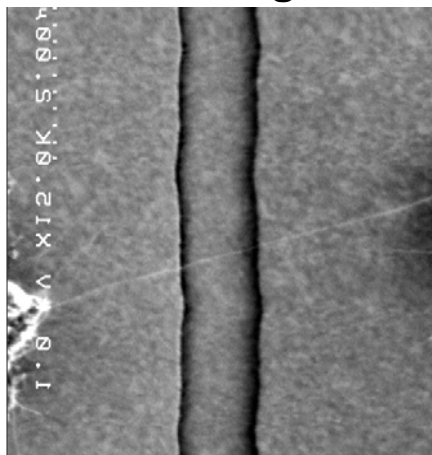
Flexible Electronics

Experimental Setup

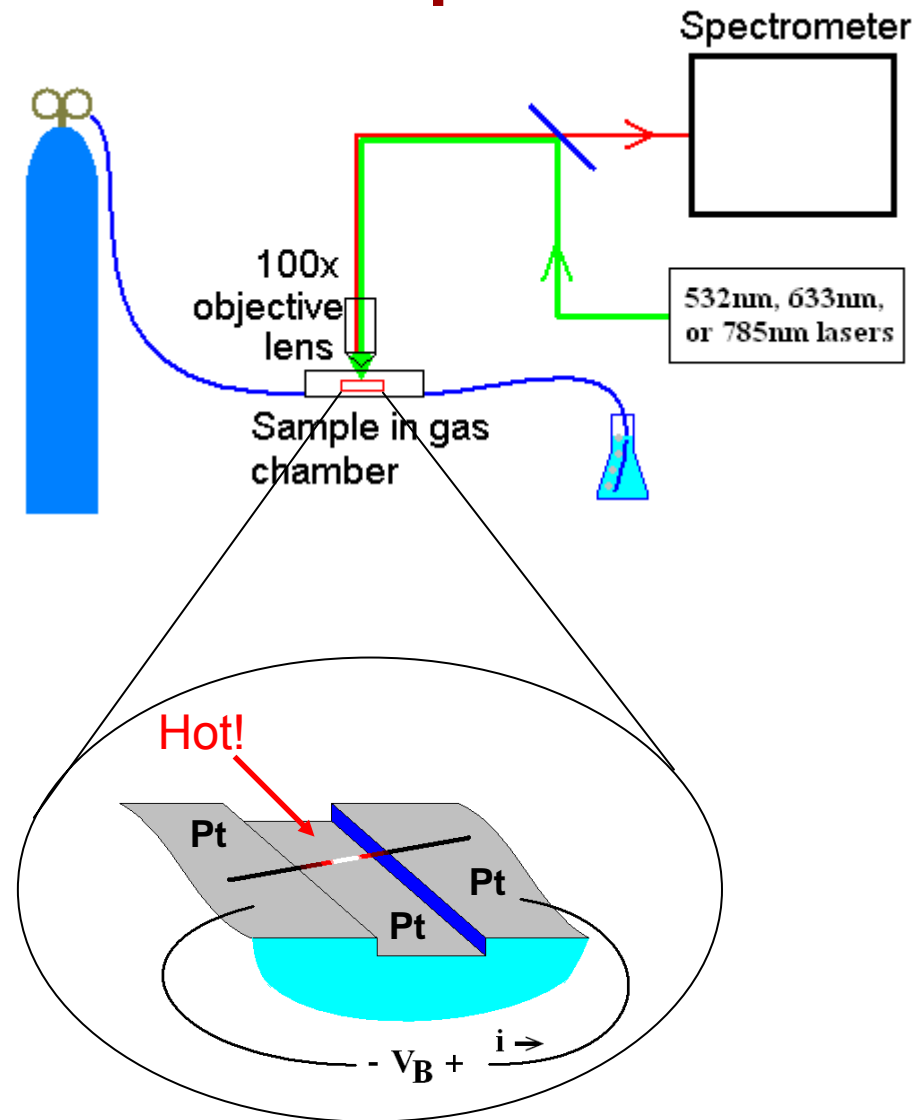
Optical Microscope Image:



SEM Image:



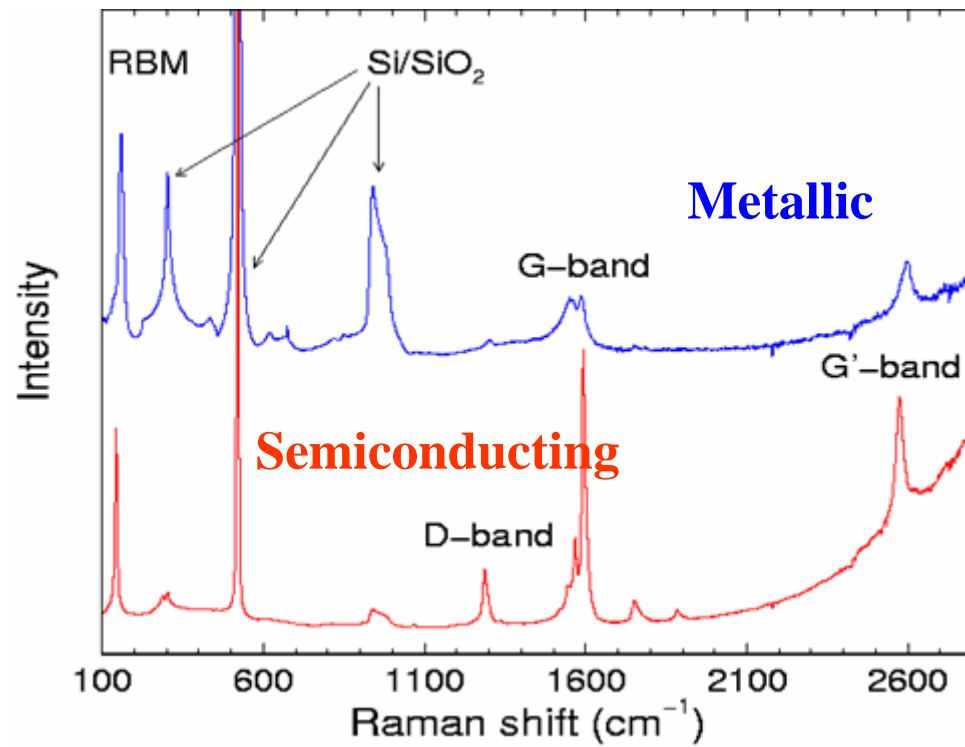
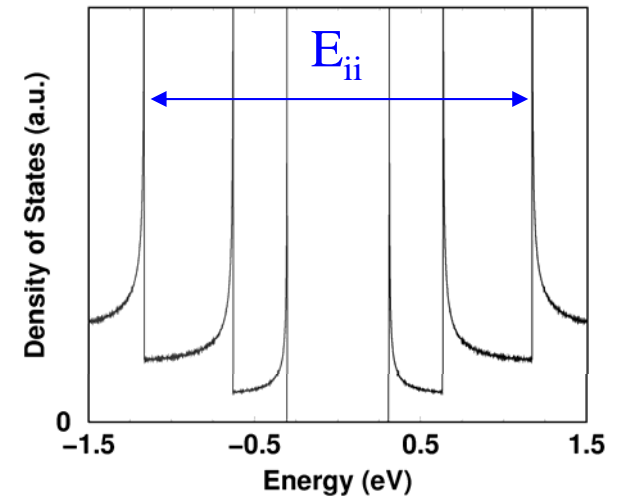
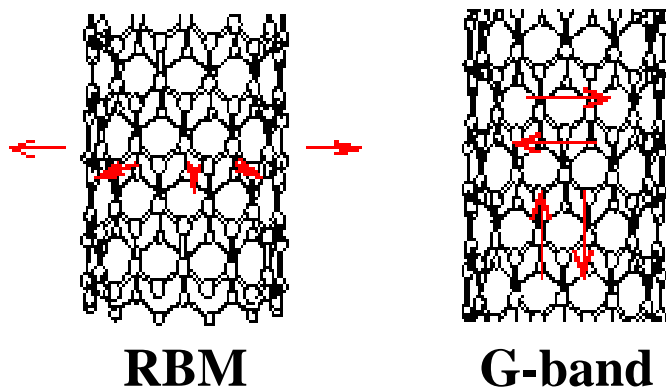
Our device (example)



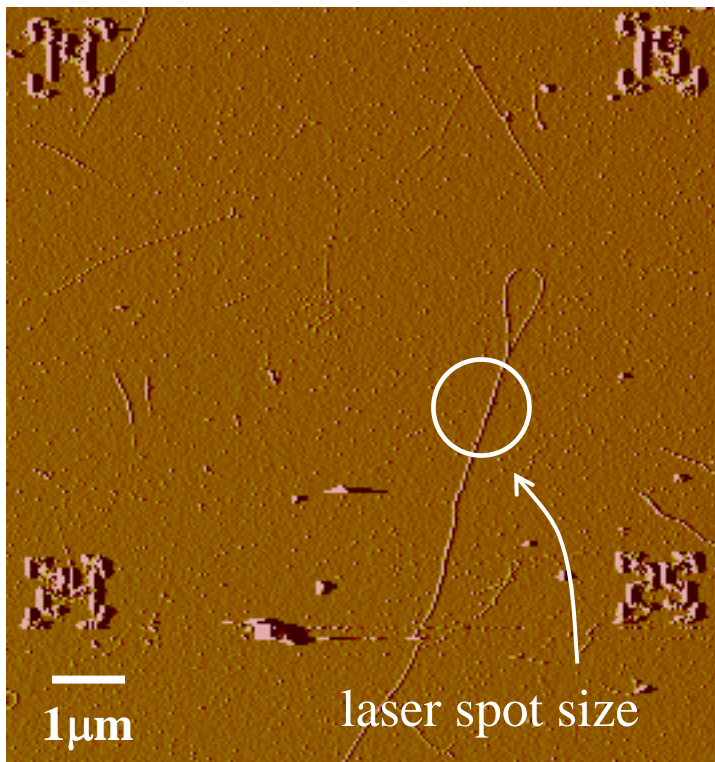
Single Nanotube Raman Spectroscopy

Despite the extremely small geometric cross-section the Raman signal from a single isolated nanotube can be observed.

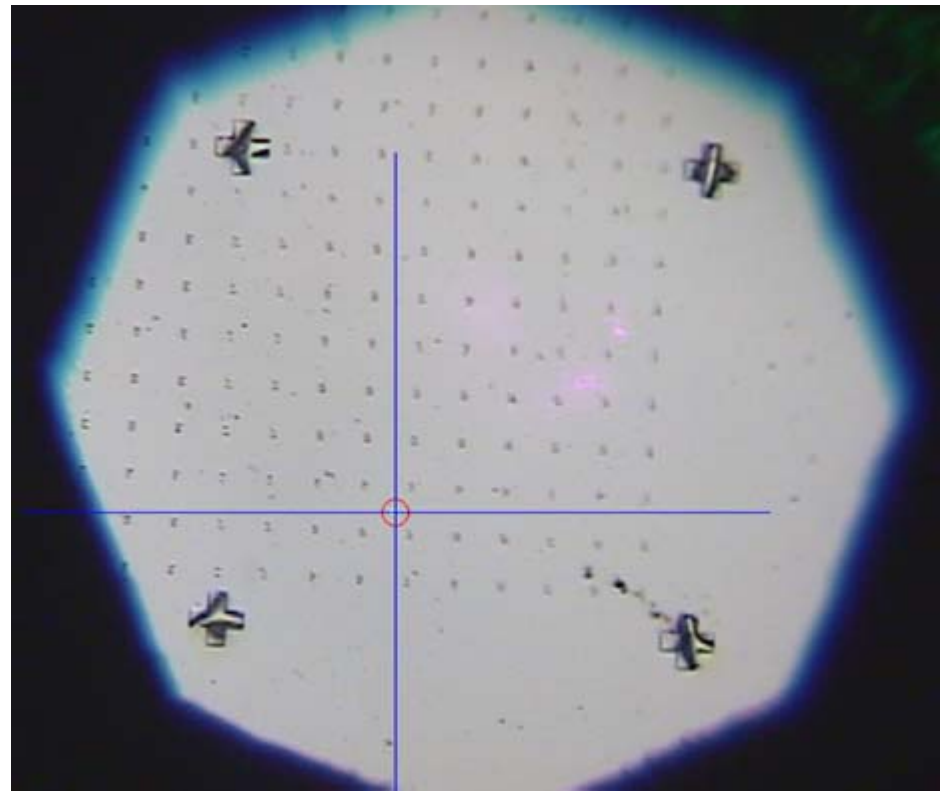
- 10^5 enhancement in scattering cross-section due to singularities in the DOS
- Resonance occurs when $E_{\text{laser}} = E_{\text{ii}}$
- Only observe nanotubes that are resonant with E_{laser}



Nanotubes on a Substrate with Grid Patterned by e-Beam Lithography

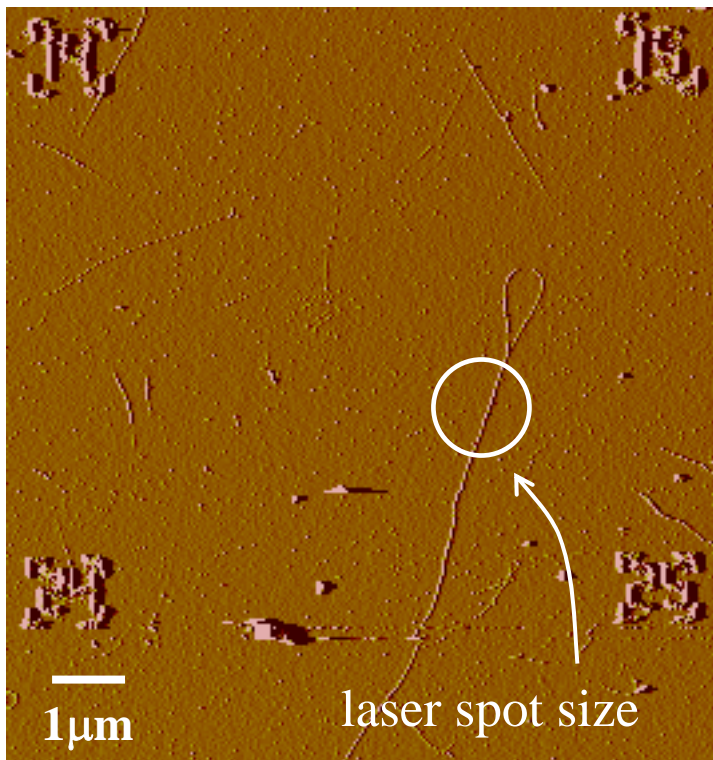


AFM Image

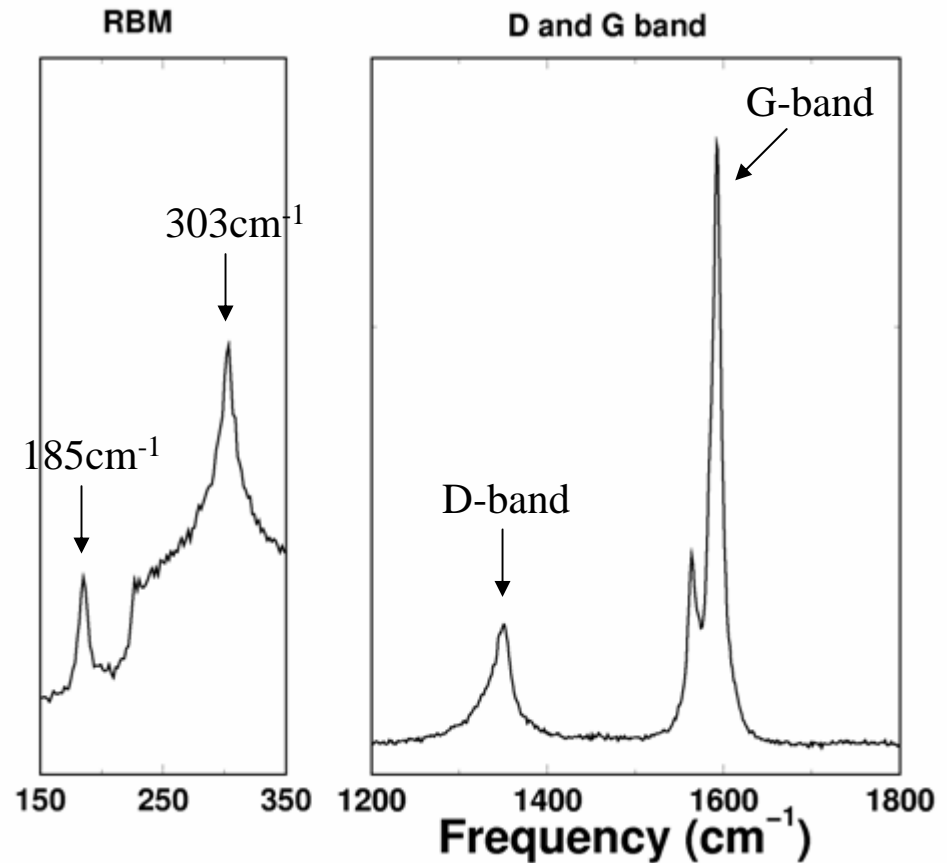


Optical Image

Raman Spectroscopy of Individual Nanotube on a Substrate with Grid



AFM Image



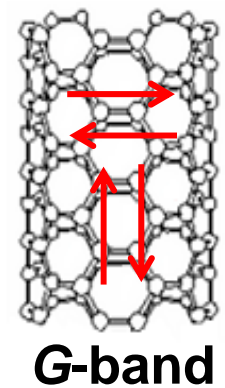
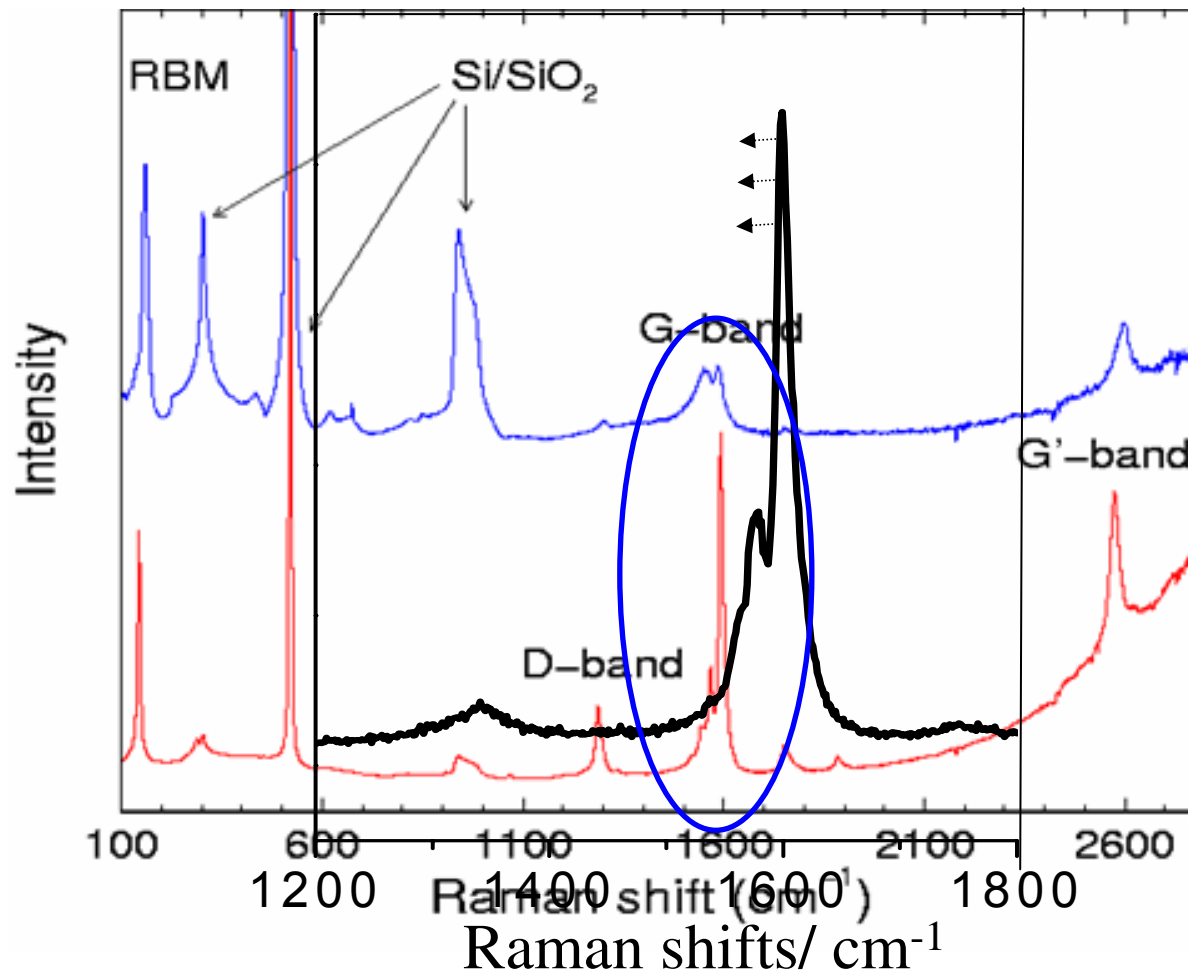
Raman Spectrum

Raman tells us:

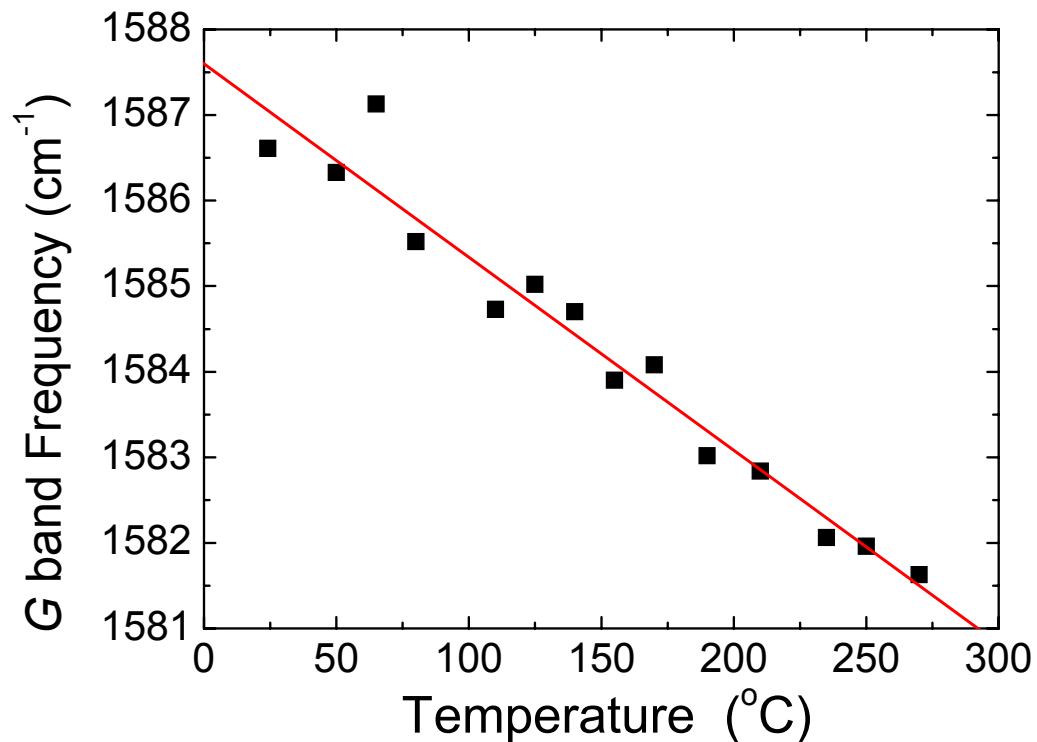
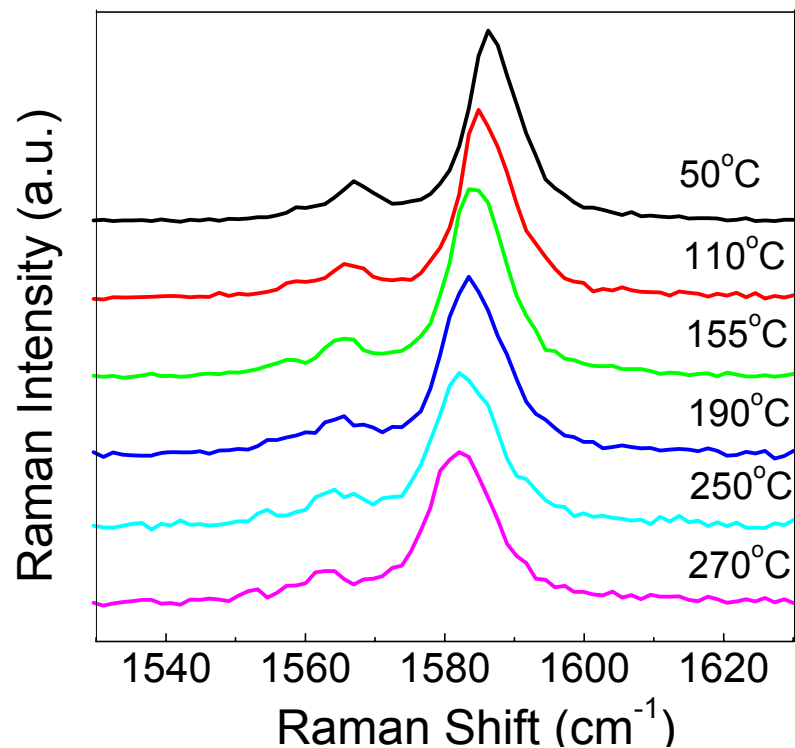
- Metal / Semiconducting nature \rightarrow G Band lineshape
- Diameter ($\pm 5\%$) \rightarrow Inverse to RBM
- Electronic transition energies (E_{ij}) \rightarrow Laser Energy
- Rough chirality $(n,m) \rightarrow E_{ij}$ and ω_{RBM}
- Orientation of a nanotube \rightarrow Polarized Raman Intensity
- Defect concentration \rightarrow D Band Intensity
- Strain / Temperature effects \rightarrow G band frequency shift
- Non-contact
- Non-destructive

Raman Spectra

- Downshift of the G-band is caused by thermal expansion when the laser heats the nanotube and lengthens the interatomic C-C length.



Temperature Coefficient of G band



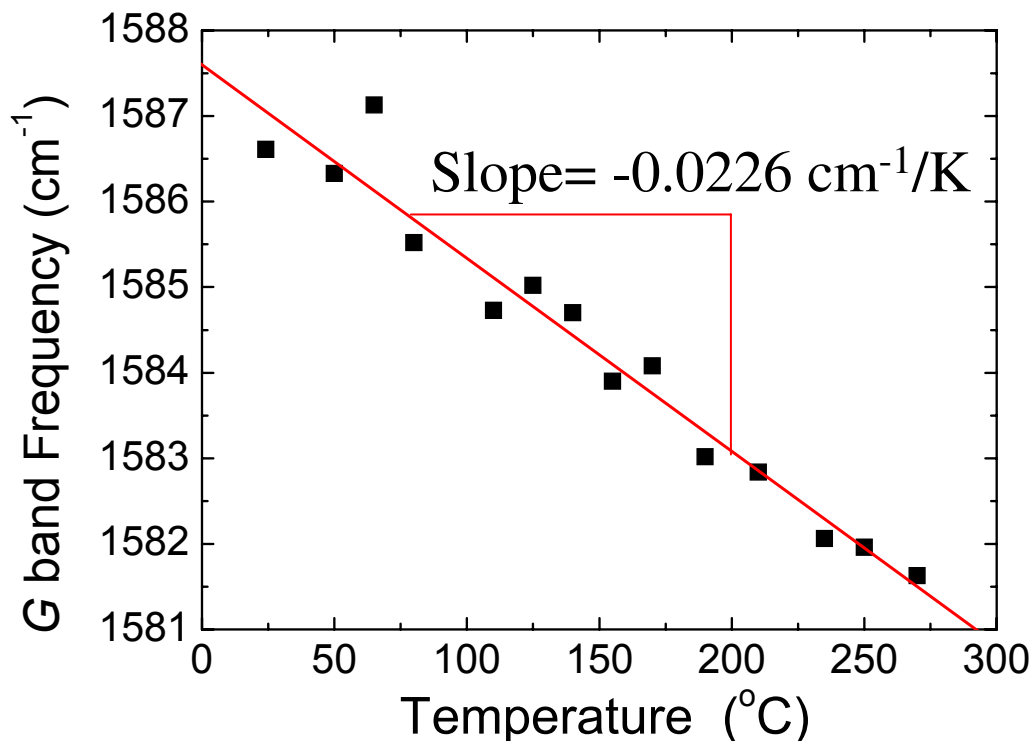
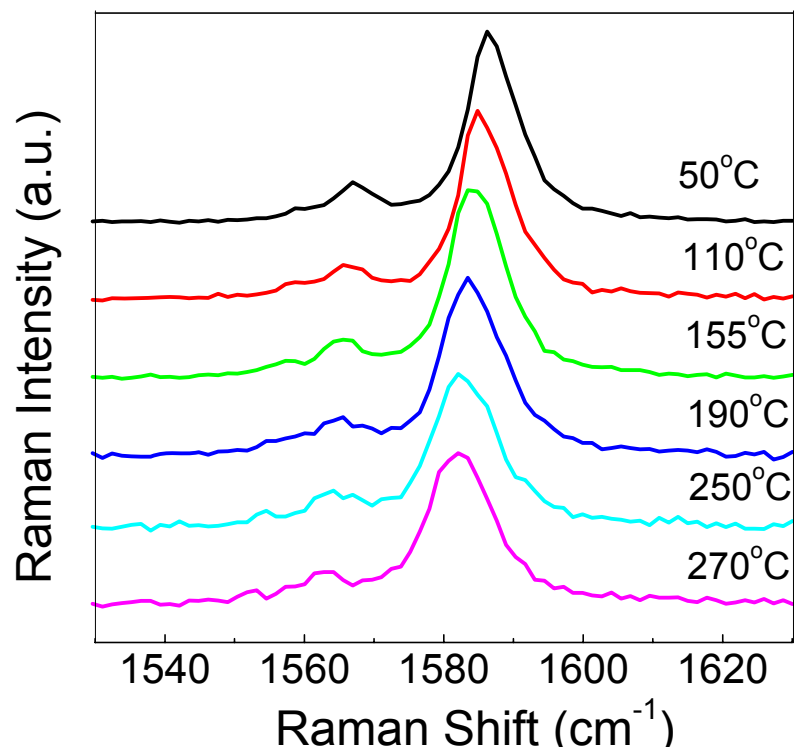
TOC = $-0.044 \text{cm}^{-1}/\text{K}$, Atashbar, *et. al.*, *App. Phys. Lett.* 86, 123112 (2005)

TOC = $-0.042 \text{cm}^{-1}/\text{K}$, H. D. Li, *et. al.*, *App. Phys.* 76, 2053 (2000)

TOC = $-0.028 \text{cm}^{-1}/\text{K}$, Huang, *et. al.*, *App. Phys.* 84, 4022 (1998)

I.K. Hsu, *et. al.*, *App. Phys. Lett.* 92, 063119 (2008)

Temperature Coefficient of G band



Temperature coefficient of G band shifts = $-0.0226 \text{ cm}^{-1}/\text{K}$

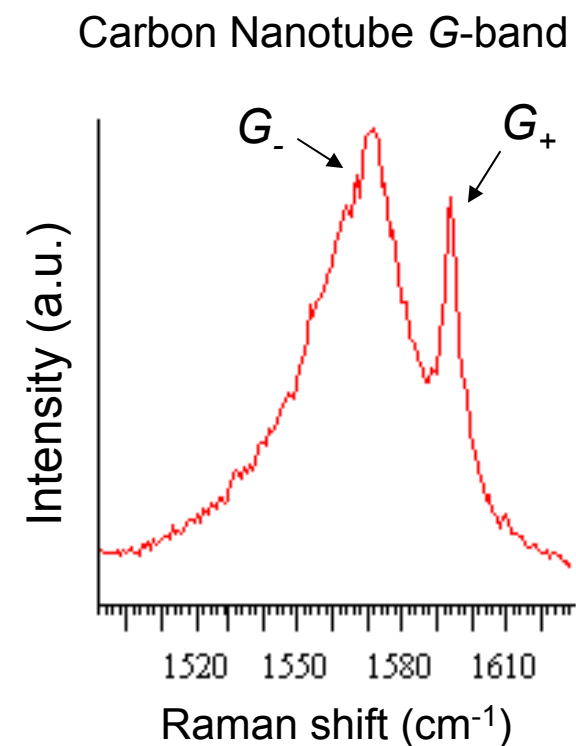
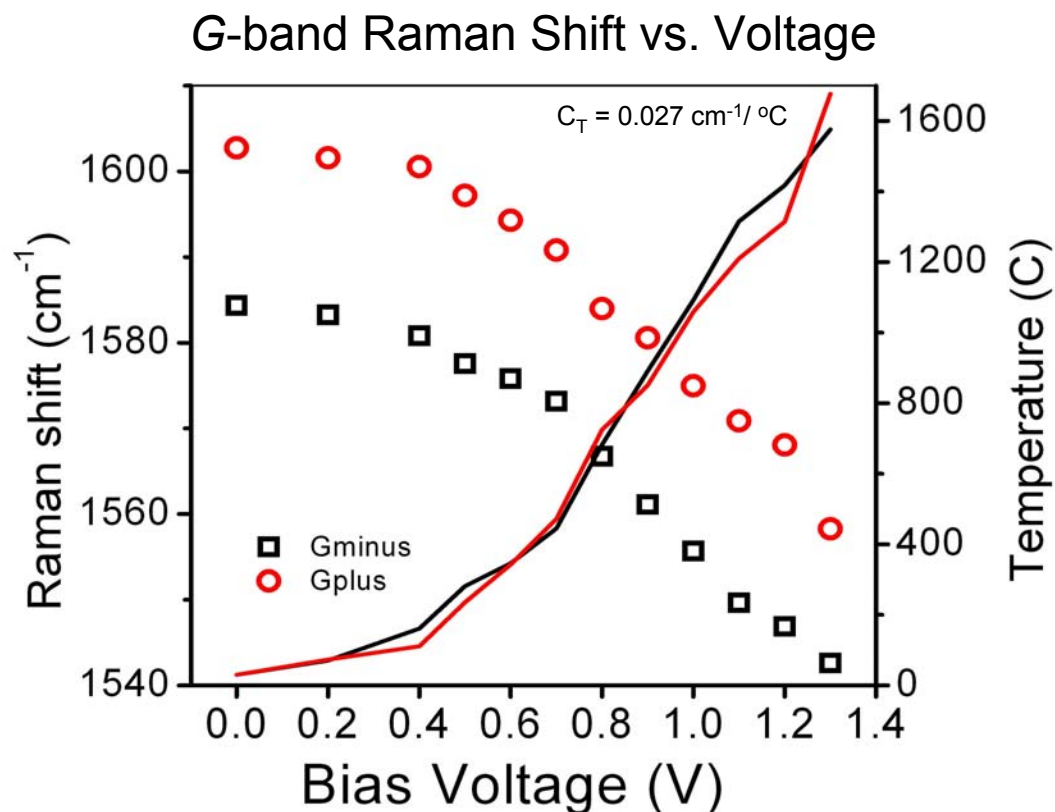
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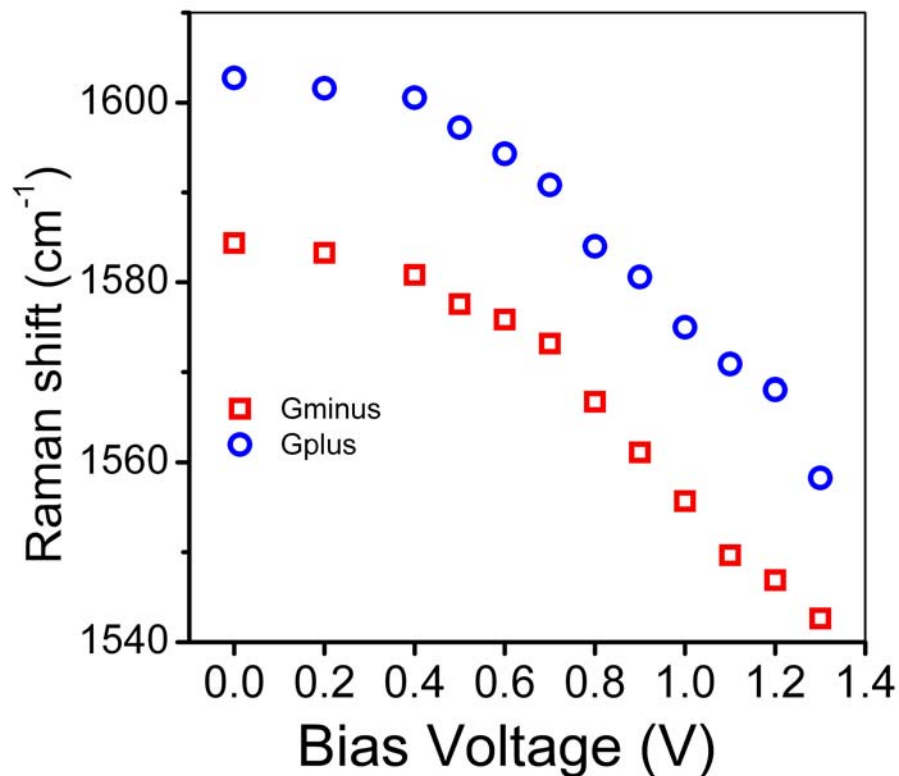
Electrical Heating



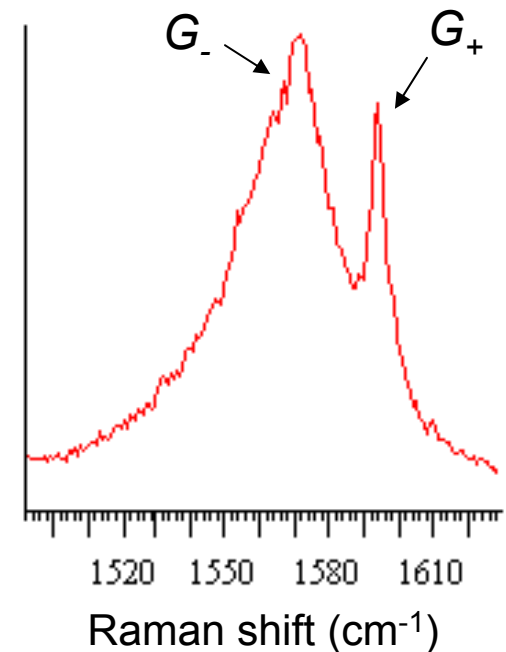
Nanotube reaches 1600°C at high voltage bias

Electrical Heating

G-band Raman Shift vs. Voltage

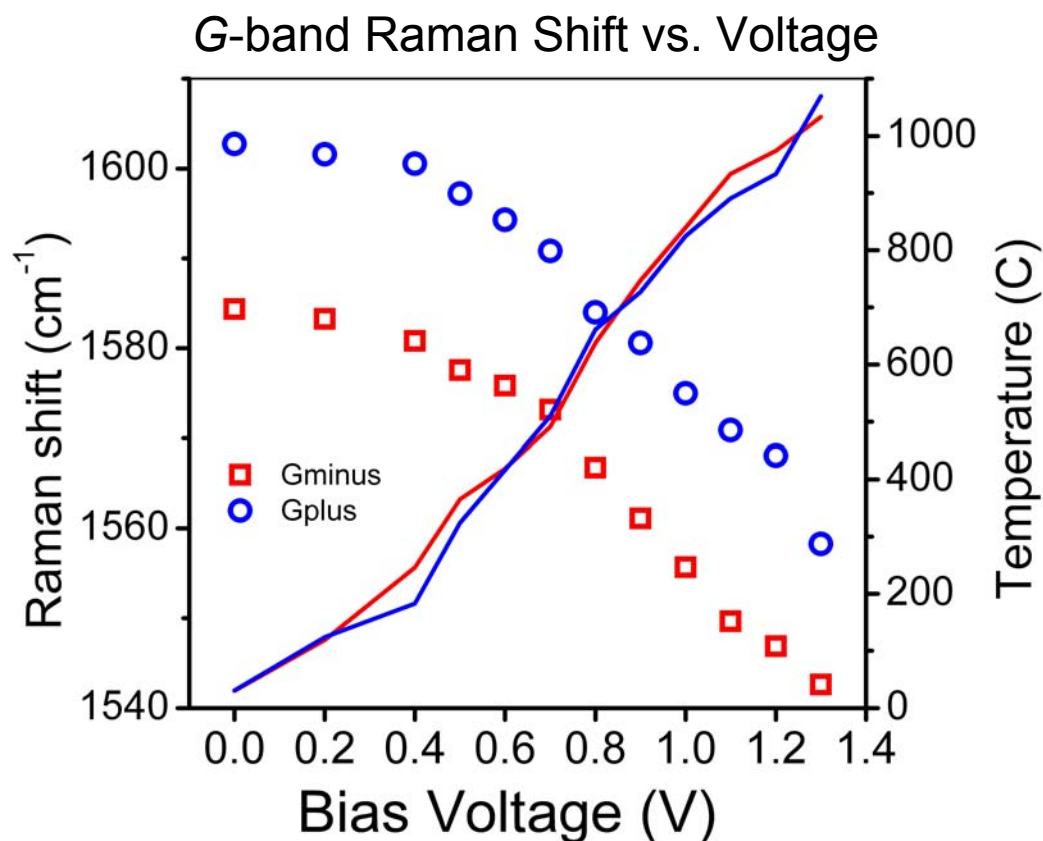


Carbon Nanotube G-band

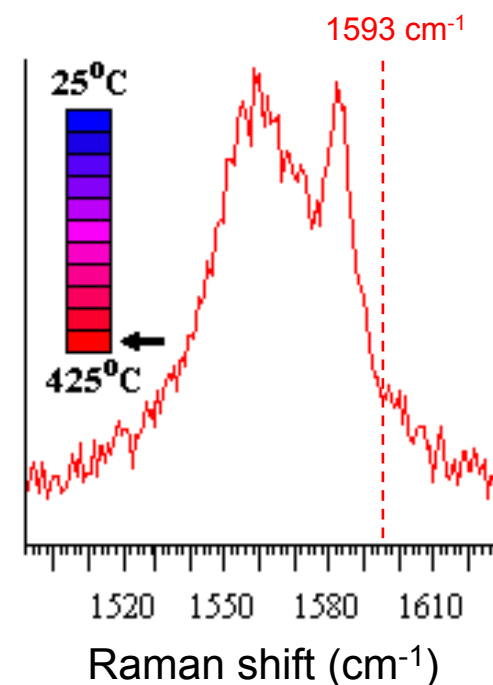


Nanotube reaches 1000°C at high voltage bias, with quadratic fit.

Electrical Heating

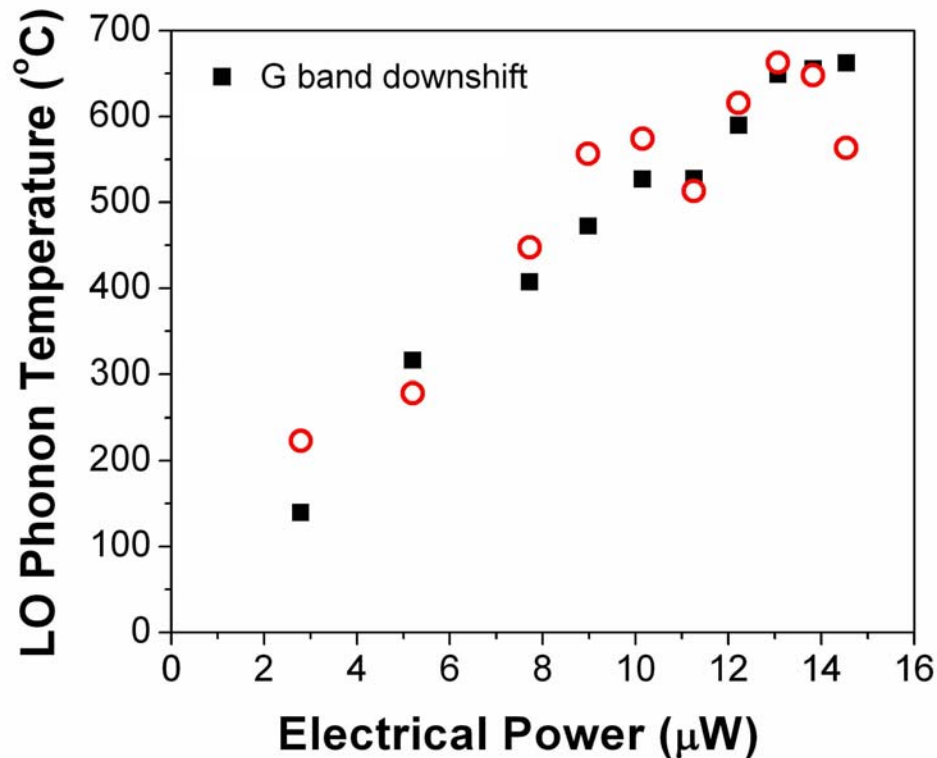


Carbon Nanotube G-band



Nanotube reaches 1000°C at high voltage bias, with quadratic fit.

Comparison of Raman Downshift vs. Stokes/AS Temperature Measurement

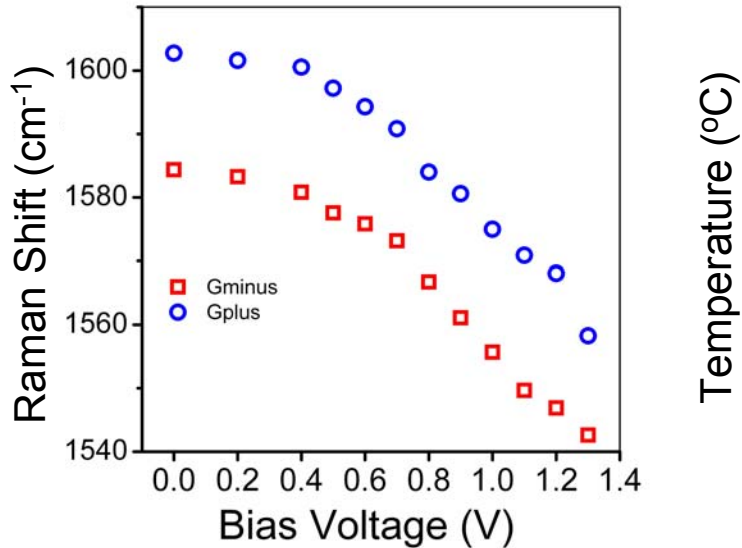


Anti-Stokes
thermal factor:

$$e^{-E_{ph}/k_B T}$$

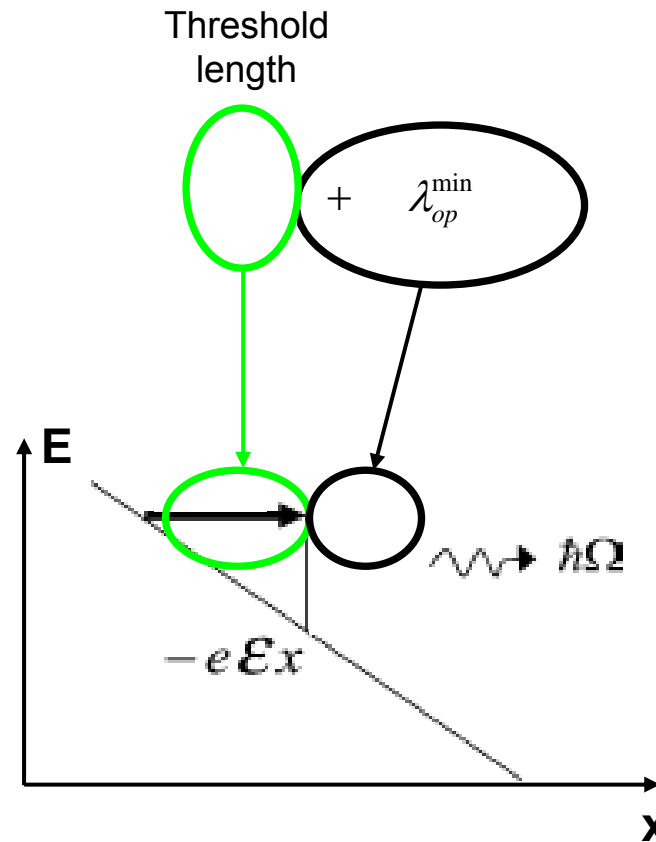
High temperatures were corroborated with AS spectroscopy

G_+ (TO) Preferential Heating



- Preferential shift indicates large non-equilibrium phonon populations
 - G_+ @ 1000°C
 - G_- @ 40°C
- An extreme state of thermal non-equilibrium

Optical Phonon Emission



- Electrons must exceed threshold phonon energy.
- And then, travel a little bit more before scattering.

Electron Transport Model

- The Landauer Model: (Pop, Mann, Park, Yao, and others)

$$R(V, T) = R_c + \frac{h}{4q^2} \left[\frac{L}{\lambda_{\text{eff}}(V, T)} + 1 \right]$$

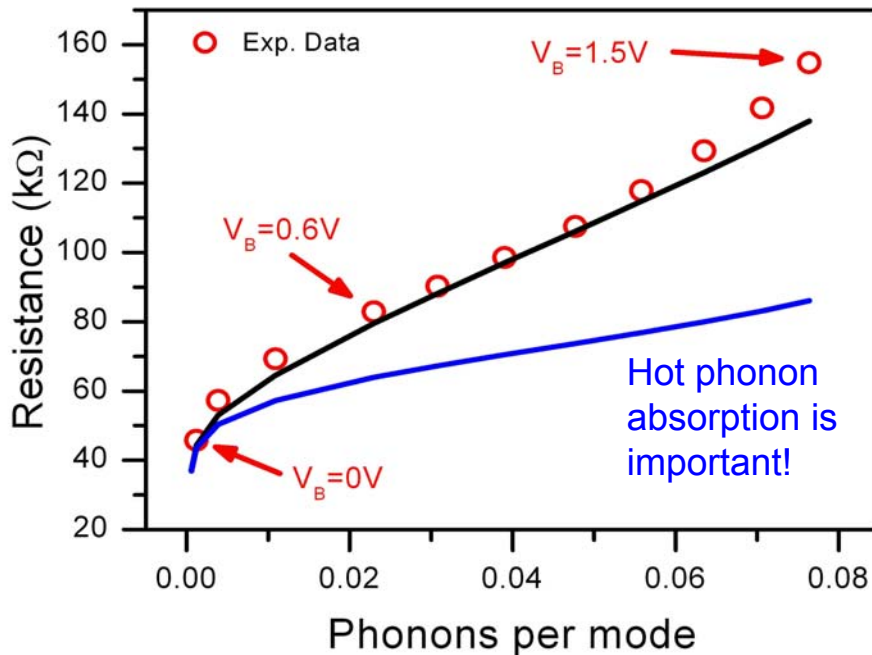
Phonon scattering contribution
Quantum Conductance

- Matthiessen's Rule:

$$\lambda_{\text{eff}}^{-1} = \lambda_{\text{ac}}^{-1} + \lambda_{\text{op,ems}}^{-1} + \lambda_{\text{op,abs}}^{-1}$$

$\lambda_{\text{op,abs}}$
Dominant contribution to high bias scattering and nanotube heating

Results of the Model



Sample	Length (μm)	M/SC	Downshift	d _t (nm)	λ_{op}^{\min} (nm)
1	0.5	M	G ₊	-	18
2	2	M	G ₋	1.70	26
3	1	M	G ₊	-	35
4	2	SC	G ₊	2.01	9
5	0.5	M	G ₊	-	28

Confirms literature values for $\lambda_{op}^{\min} \sim 15 - 30 \text{ nm}$

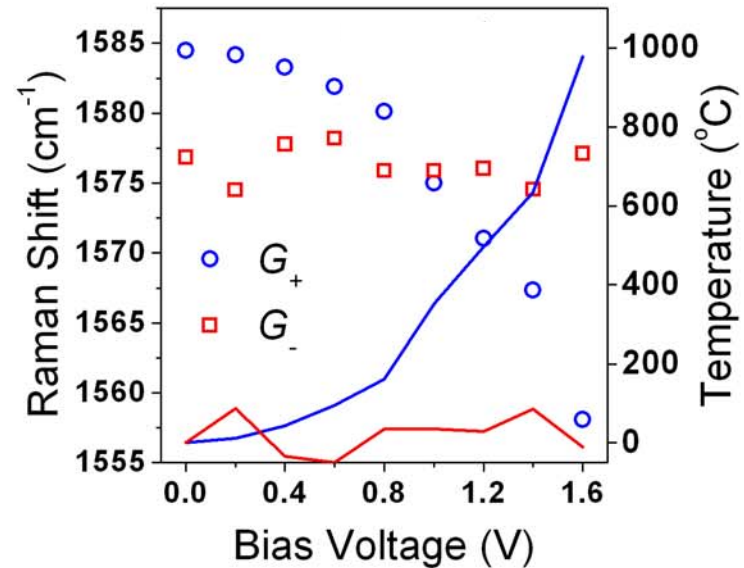
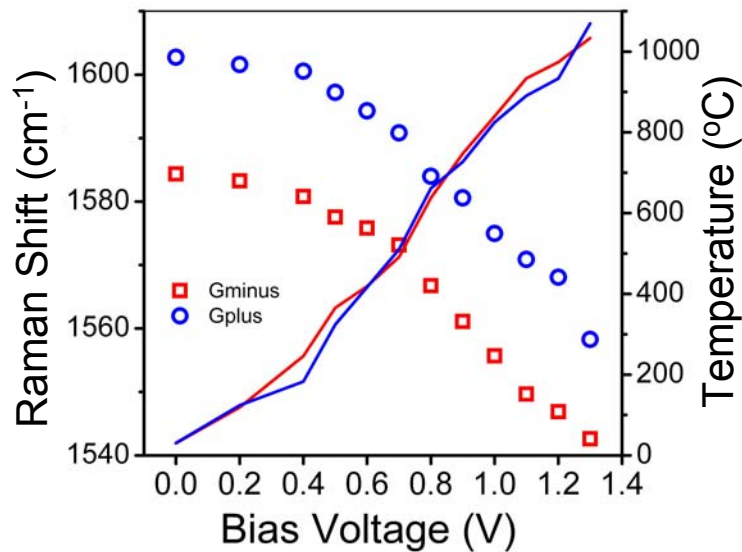
- Park *et al.* Nano Lett. **4** 517 (2004)
- Yao *et al.* PRL **84** 2941 (2000)
- Javey *et al.* PRL **92** 106804 (2004)
- Pop *et al.* PRL **95** 155505 (2005)

Apply the Landauer model:

R_C , λ_{op}^{\min} = fitting parameters

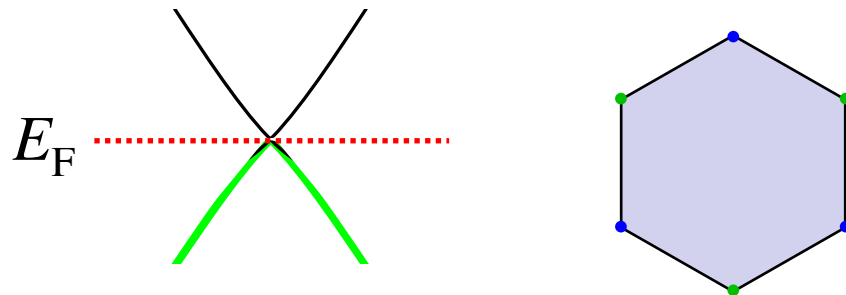
$$R(V, T) = R_c + \frac{h}{4q^2} \left[\frac{L}{\lambda_{eff}(V, T)} + 1 \right]$$

G_+ Preferential Heating



Why is there preferential heating of only one phonon mode?

The Kohn Anomaly

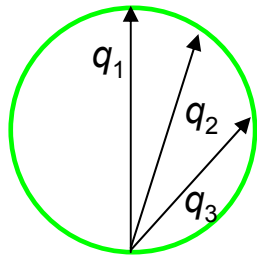


$$E_{\text{TOT}} = \int_{-\infty}^{E_F} E \cdot D(E) dE$$

- Strong coupling of electrons and phonons at the Fermi energy.
- Atomic displacements create dynamic bandgap, which lowers the energy of the phonons.

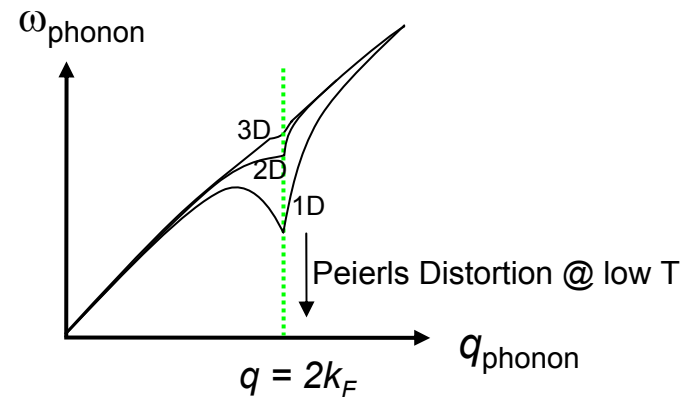
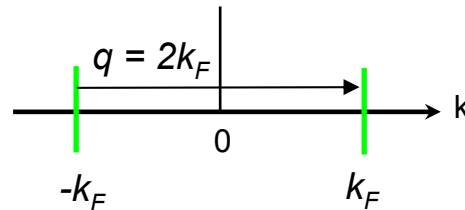
The Kohn Anomaly in 1D

Fermi surface in 2D



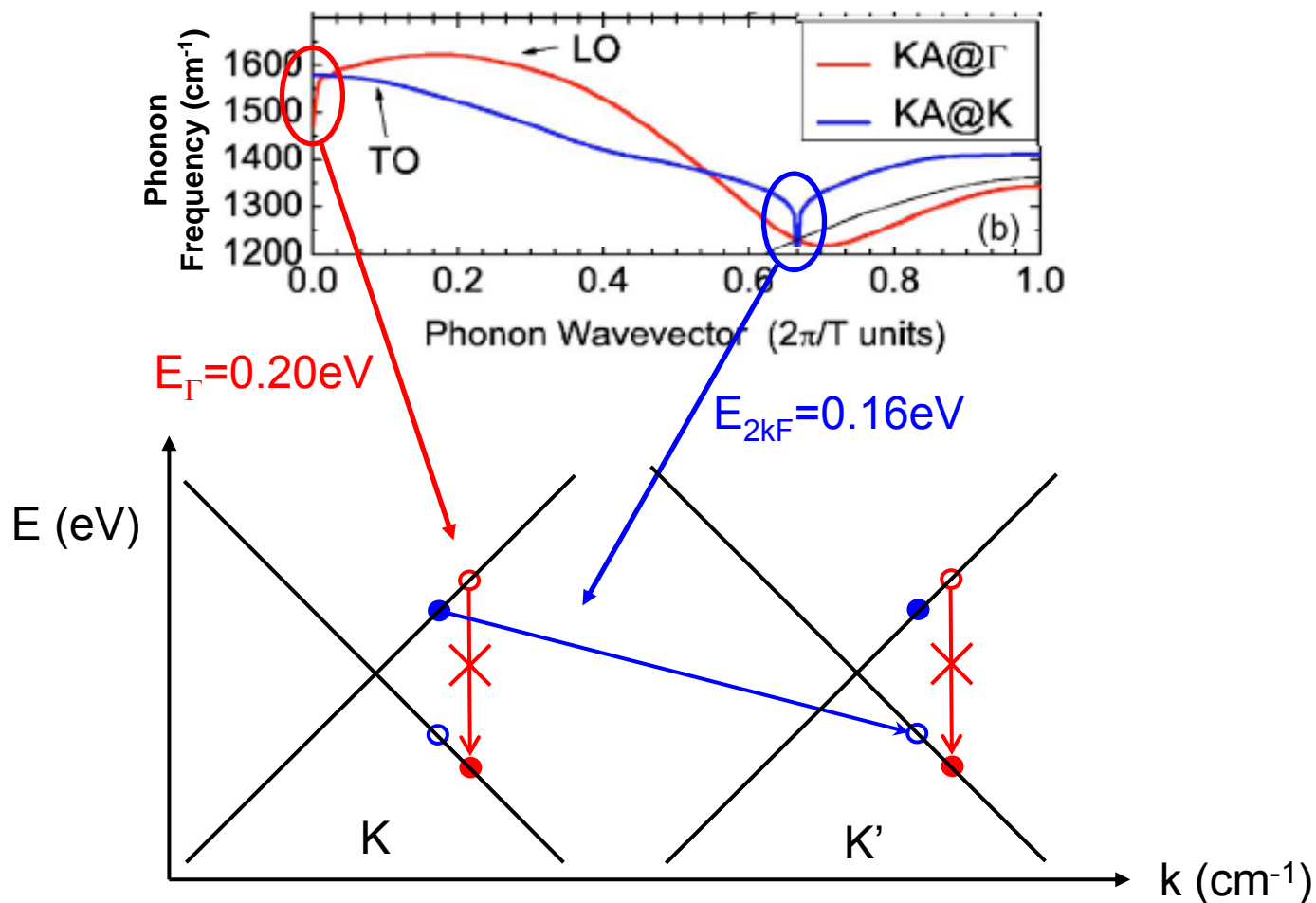
$$q_1 \neq q_2 \neq q_3 \leq 2k_F$$

Fermi surface in 1D



- The Kohn anomaly is exacerbated in 1D because of the limited number of phonon and electronic states.
- At low temperatures the phonon frequency reduced to zero resulting in a permanent lattice (Peierls) distortion.

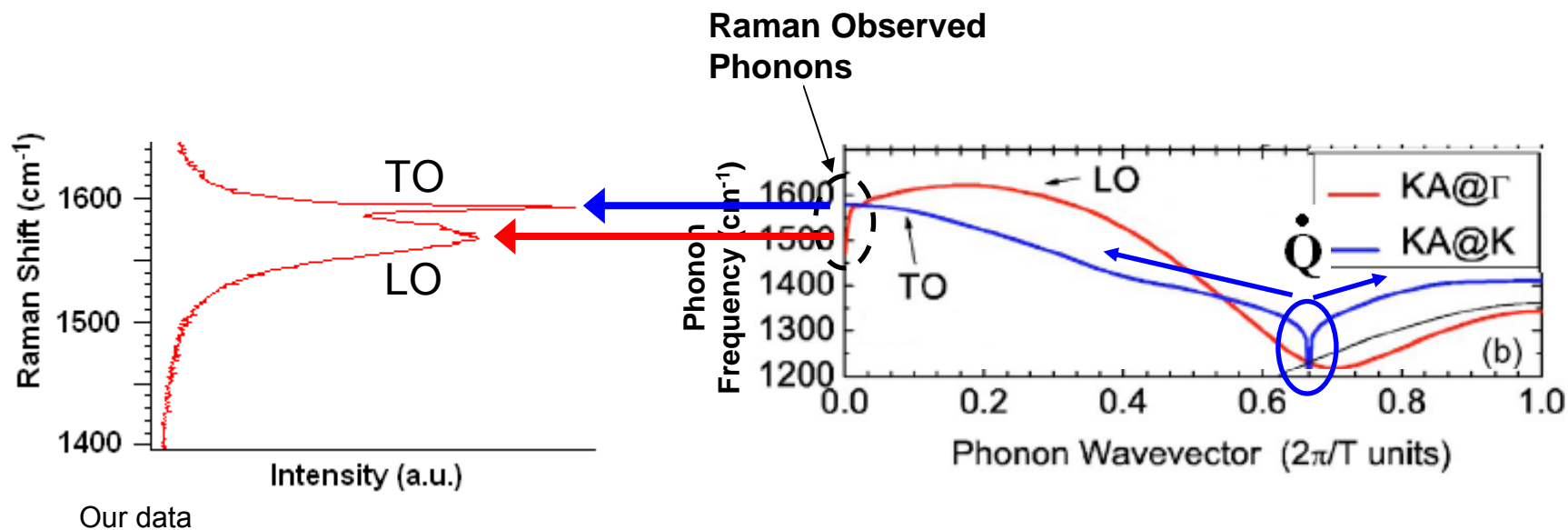
Preferential Electron-Phonon Coupling



$E_{2k_F} < E_{\Gamma} \implies$ Only TO phonons are emitted

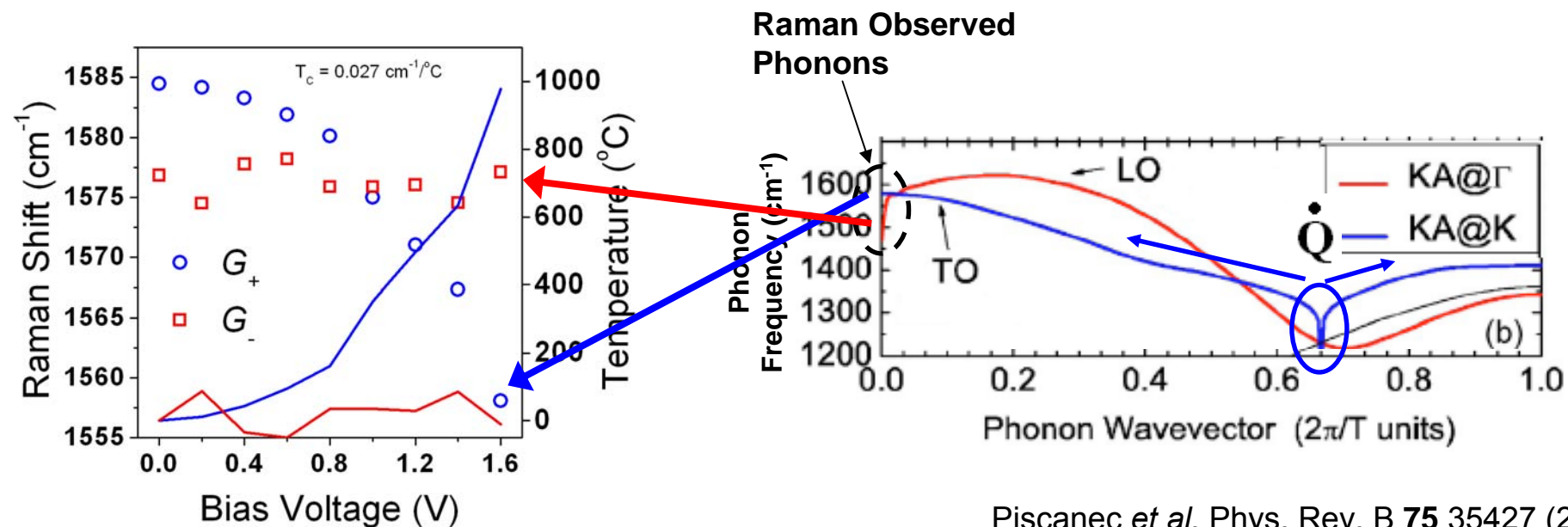
Observation of $2K_F$ Optical Phonon Preferential Coupling

- $2k_F$ phonons are heated
- Each phonon band is expected to be in thermal equilibrium
- We observe the temperature of the bands with Raman spectroscopy
- The TO band is preferentially heated

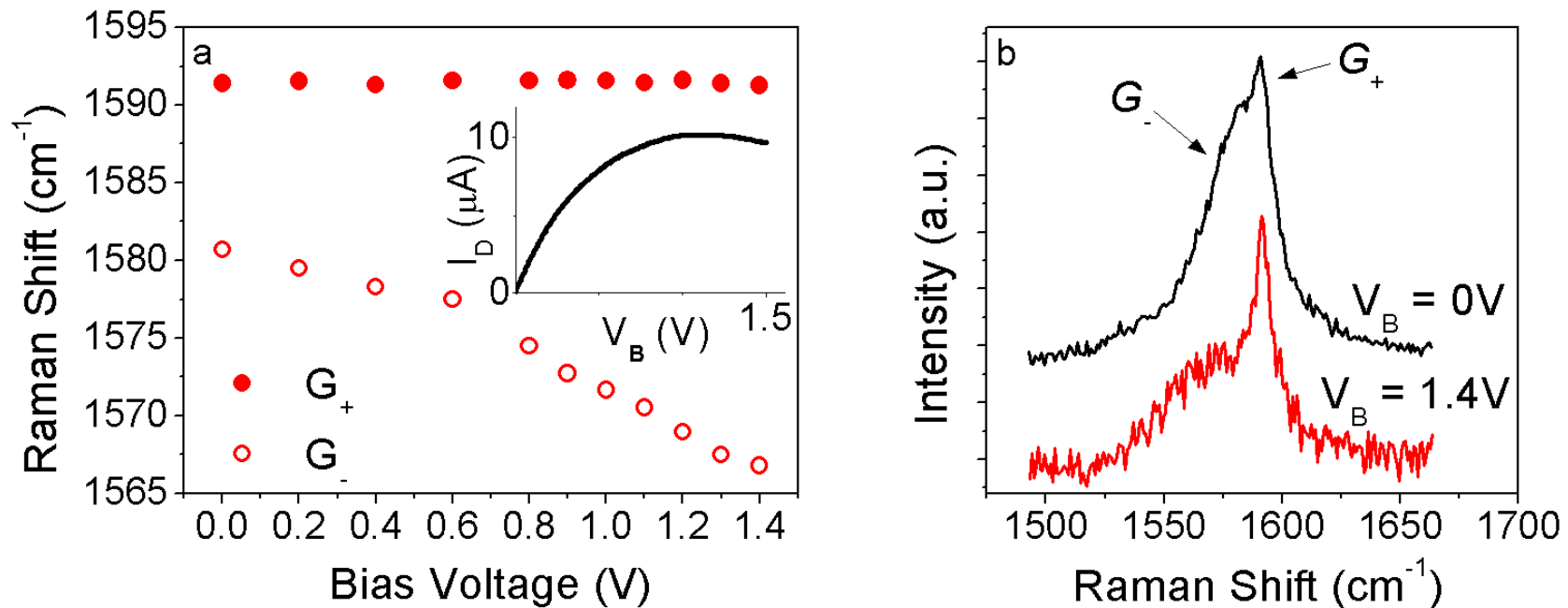


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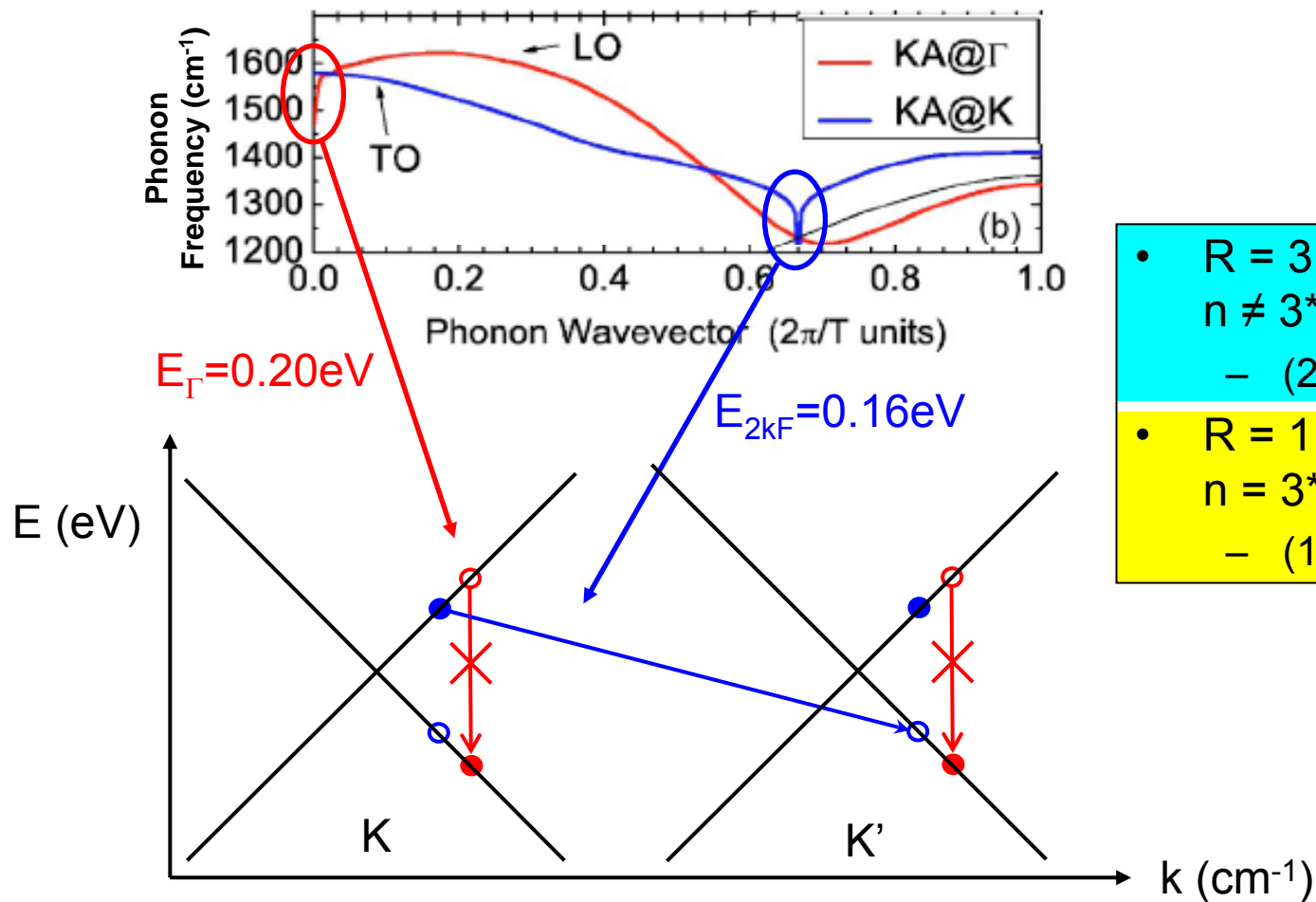


Sometimes, Only G_- (LO) is Heated



- Even less frequently, only G_- is heated (1 tube)
- Theory: rare chirality with non-Raman active TO phonon branch
 - Results in observation of heating in only LO branch
- Two types of metallic nanotubes: $R = 1,3$
 - LO heating is observed in $R = 1$ nanotubes

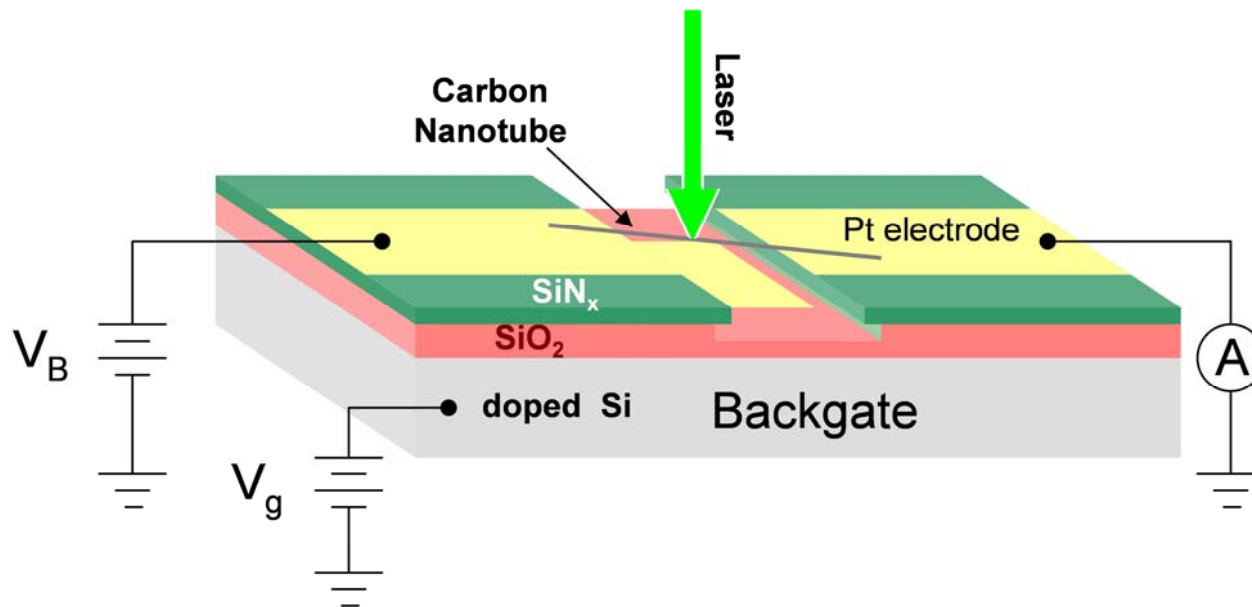
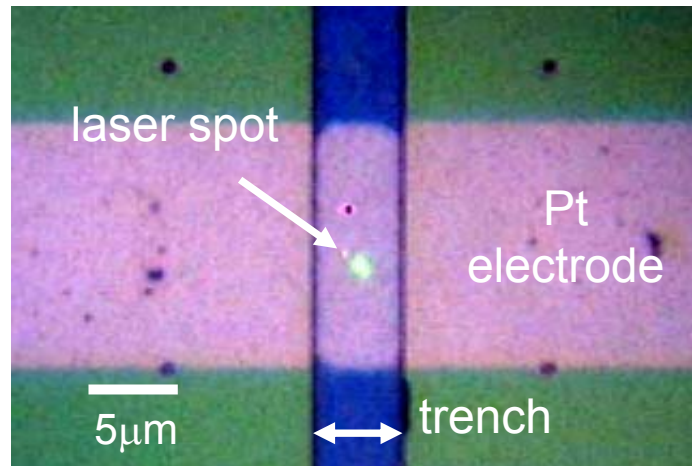
Preferential G_{-} (LO) Heating



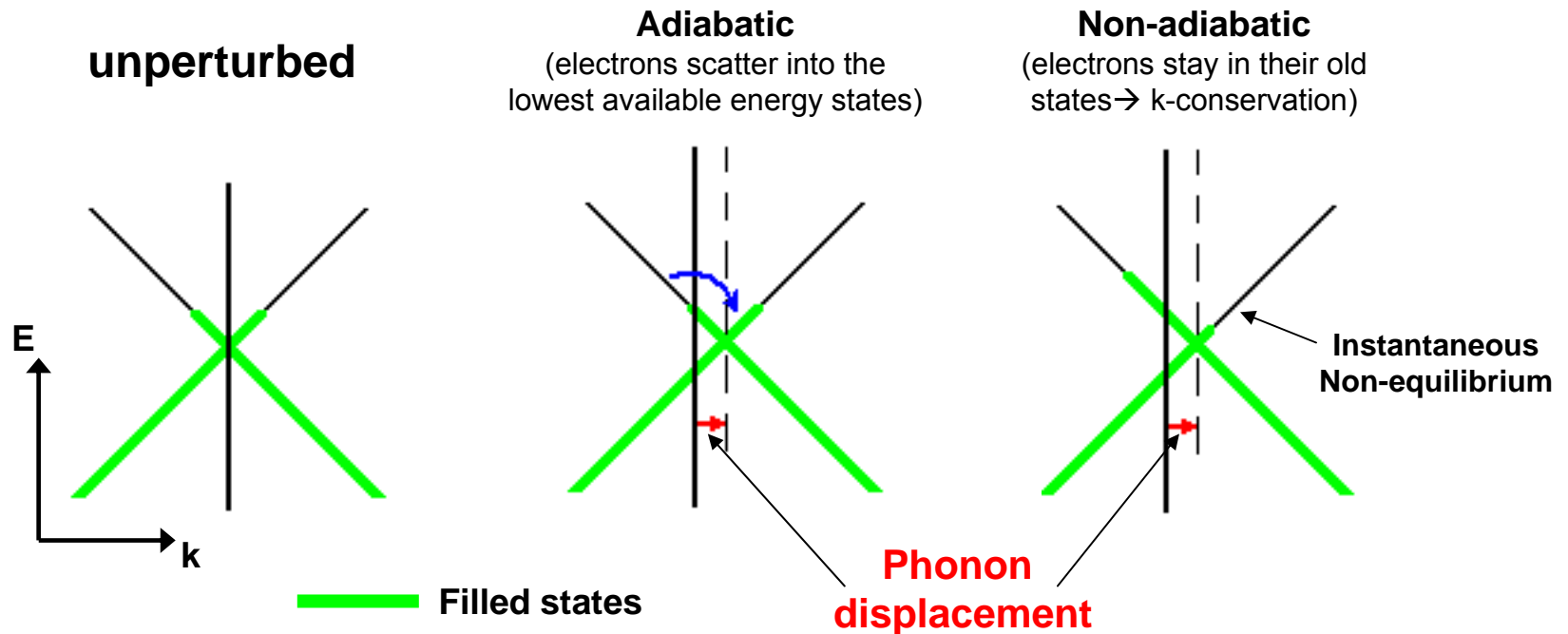
- R = 3 Armchair tubes **or** n ≠ 3*i* tubes
 - (2/3 of m-CNTs)
- R = 1 Non-armchair **and** n = 3*i* tubes
 - (1/3 of m-CNTs)

E_{Γ} only \Rightarrow Only LO phonons are emitted

Gate Voltage

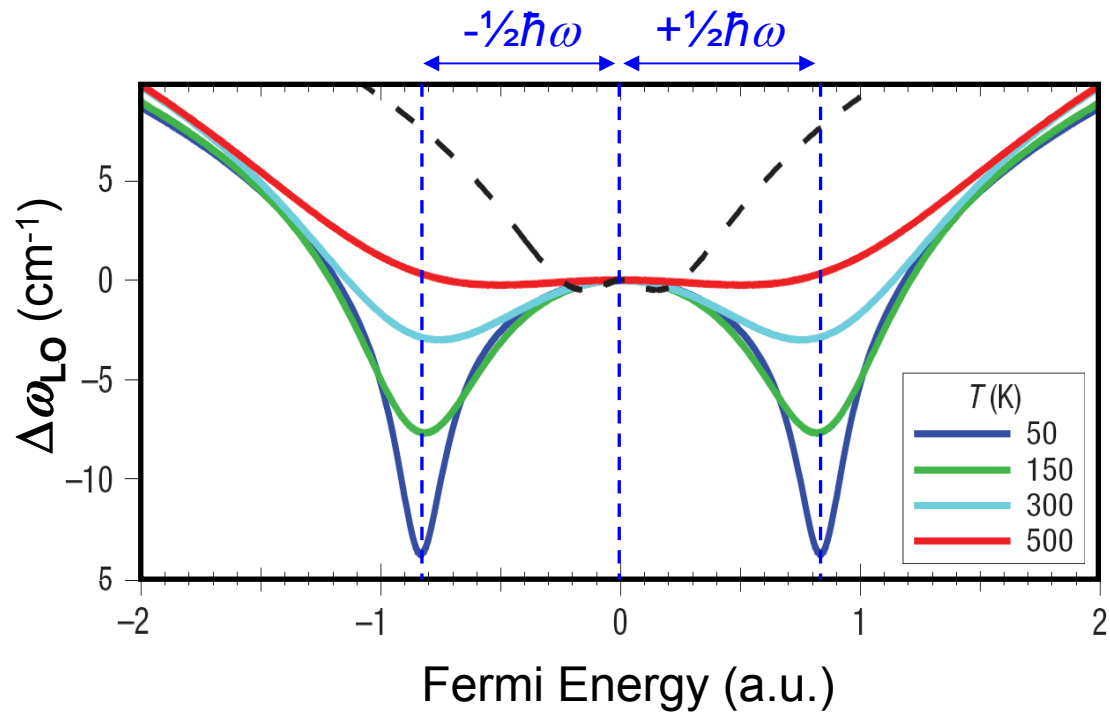
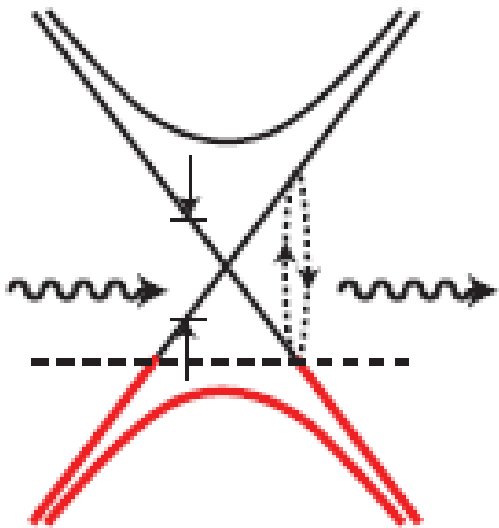


Born-Oppenheimer Approximation



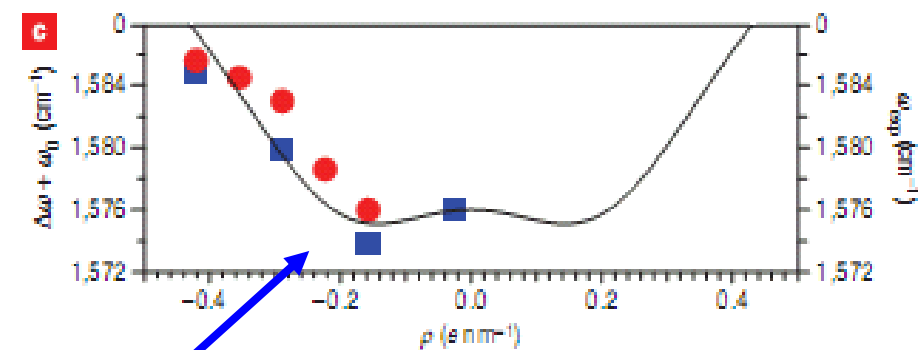
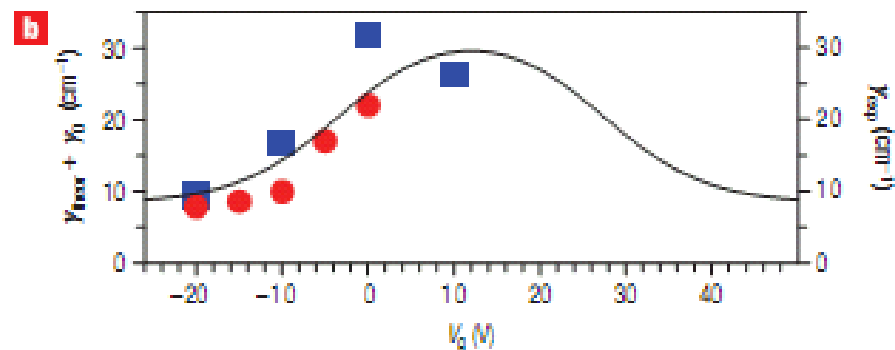
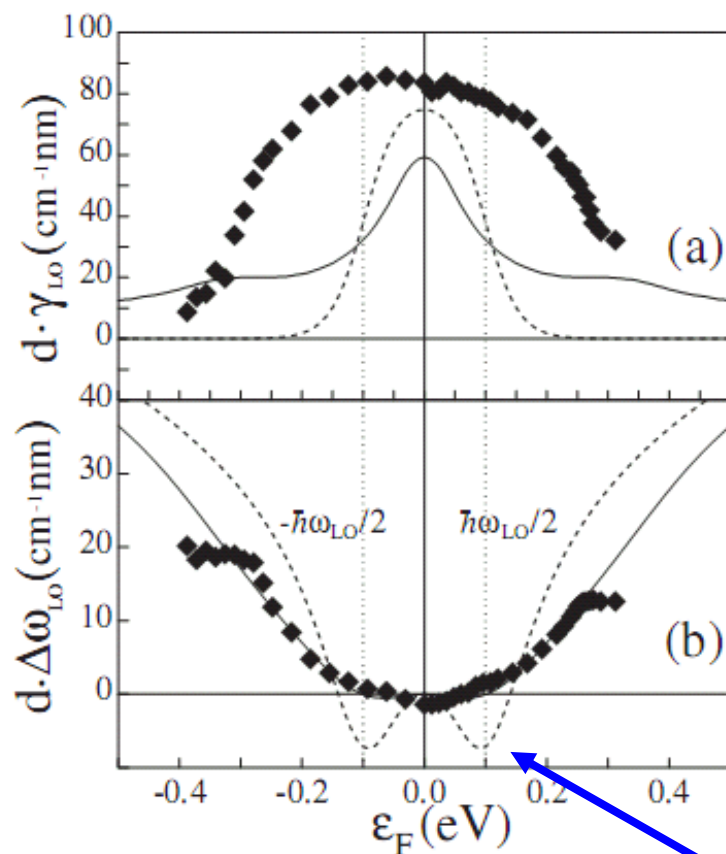
- The approximation: Electrons equilibrate to the instantaneous ground state faster than the atoms vibrate.
- Approximation breaks down in SWNTs
 - 1.) fast vibrational motion ($\tau_{OP} = 0.02\text{ps}$)
 - 2.) long electron lifetimes ($\tau_{\text{electron}} = 2\text{ps}$)

Non-Adiabatic Phonon Renormalization



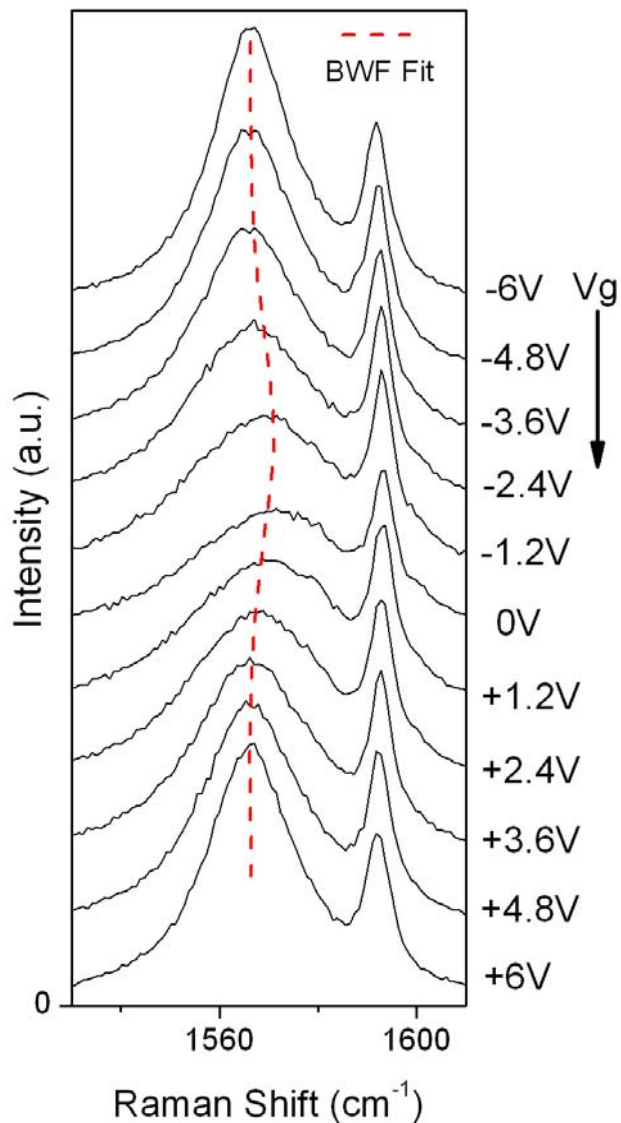
$$\omega_{\Gamma} = \sqrt{\frac{D_{\Gamma}^0}{M} + \frac{D_{\Gamma}^{KA}}{M}}$$

Comparison with Literature



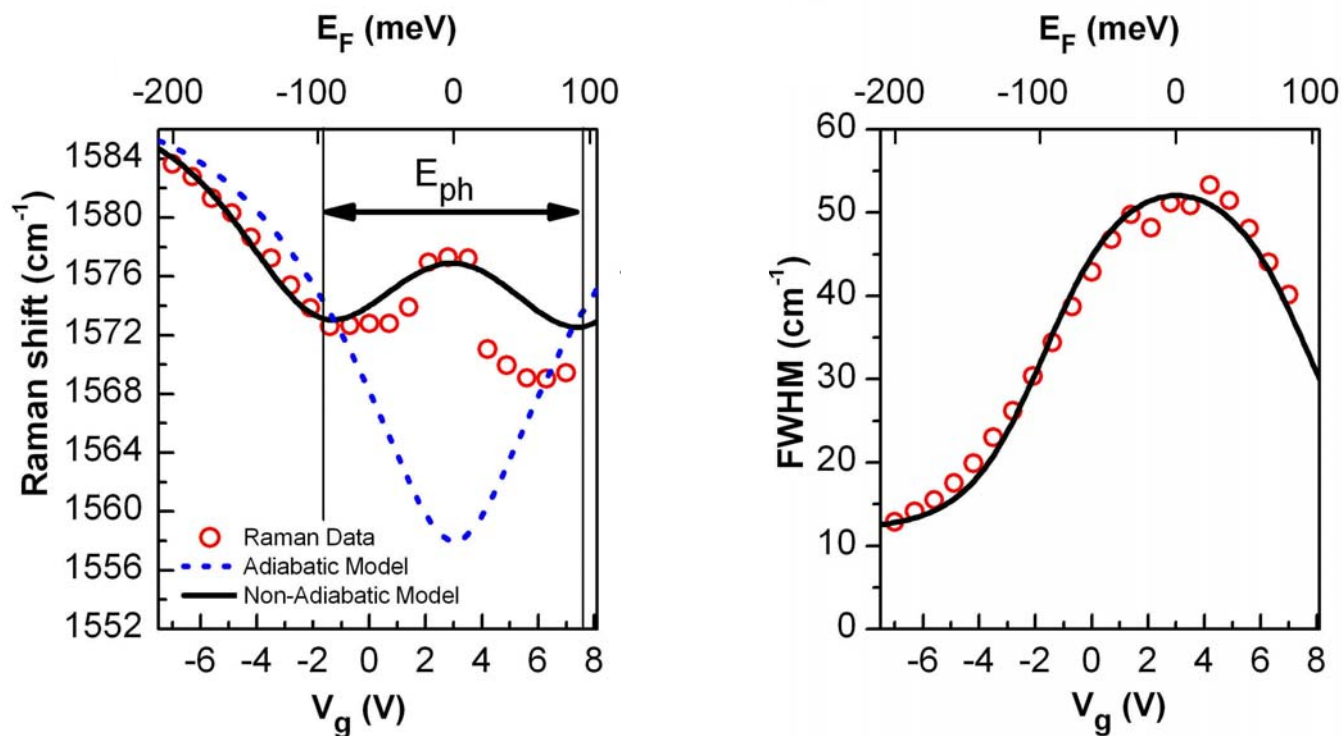
"W" shape predicted by theory but not observed in experiment!

Gate Voltage Dependence



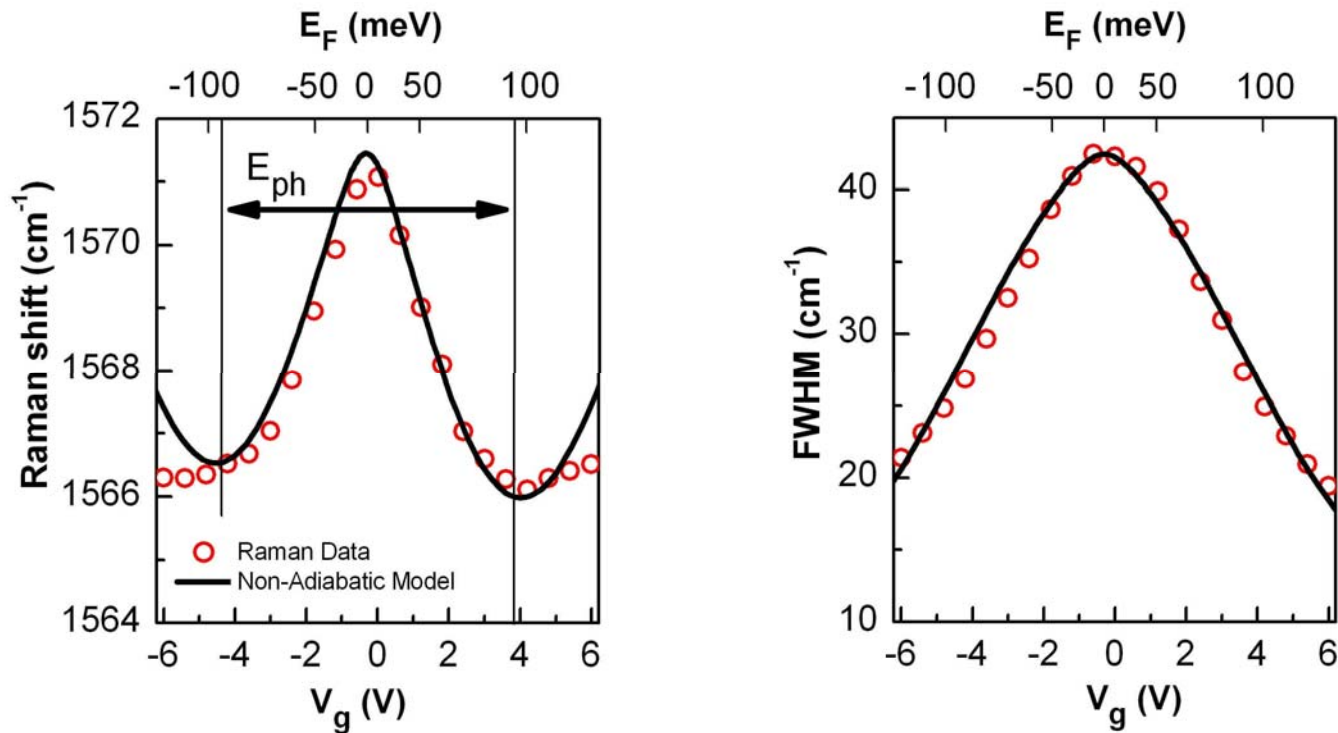
Our Devices:
Few defects
Suspended
Long electron lifetimes

Gate Voltage Dependence



- “W” shape indicates breakdown of the BOA.
- Non-adiabatic model also predicts FWHM.

Gate Voltage Dependence



- “W” shape is more pronounced.
- Good agreement with FWHM.

Conclusion

- Carbon nanotubes provide an ideal system for studying 1D physics.
- Preferential heating, mode selective el-ph coupling, non-equilibrium phonon populations, negative differential conductance (NDC), caused by the Kohn anomalies.
- Breakdown of the Born-Oppenheimer approximation, due to long electron lifetimes.

Thank You

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NSF Graduate Research Fellowship Program